

A GEOMETRICAL STRUCTURE IN THE λ -PARAMETER FAMILY OF GENERALIZED METRIC SPACES *

By Panayiotis STAVRINOS and Satoshi IKEDA.

Introduction. In the context of the time dependent geometric subjects of the Finslerian approach, the authors will provide a different standpoint from the other one that has already been developed for the rheonomic systems. In that case, the time is taken as an arbitrary (external) parameter ([1]¹), [6]).

In the present work, we study the geometry of the Finslerian space-time structure in relation with a λ -parameter family of a Finsler space-time $(F^{(4)}, \lambda)_{\lambda \in I}$ and its fiber metric structure which was proposed by one of the authors [5].

From a physical viewpoint, this consideration of λ -parameter as an *internal coordinate*, gives rise to an "evolution" of the Finslerian space-time $F_0^{(4)}$. For this reason, it is useful to consider a *second order deformed bundle* $T^{(2)}(DF)$ as the total space of the deformation of $F_0^{(4)}$ constructed locally by $X^A = (x^k, y^i, y^0 = \lambda)$, where the internal variable λ is considered independent of x^k and y^i , ($k, i = 1, 2, 3, 4$).

§ 1. **Connection.** The metrical structure in $T^{(2)}(DF)$ is given by the form

$$G = g_{\lambda k} dx^k dx^\lambda + g_{ij} \delta y^i \delta y^j + g_{00} \delta \lambda \delta \lambda,$$

with $\delta \lambda = d\lambda + M_\lambda^0 dx^\lambda + L_i^0 dy^i$, $\delta y^i = dy^i + N_\lambda^i dx^\lambda$, $\lambda = y^0$.

The coefficients $M_\lambda^0(x, y, \lambda)$, $L_i^0(x, y, \lambda)$ transform like the components of a co-vector field on M :

$$\tilde{M}_\lambda^0 = \frac{\partial x^k}{\partial \tilde{x}^\lambda} M_k^0, \quad \tilde{L}_j^0 = \frac{\partial x^j}{\partial \tilde{x}^j} L_i^0, \quad \tilde{N}_j^k \frac{\partial \tilde{x}^j}{\partial x^i} = \frac{\partial \tilde{x}^k}{\partial x^j} N_i^j - \frac{\partial^2 \tilde{x}^k}{\partial x^i \partial x^j} y^j.$$

In an adapted frame on $T^{(2)}(DF) = \left(\frac{\delta}{\delta x^i}, \frac{\delta}{\delta y^i}, \frac{\partial}{\partial \lambda} \right)$, the form of local bases will be given, as follows:

$$\frac{\delta}{\delta x^\mu} = \frac{\partial}{\partial x^\mu} - N_\mu^i \frac{\partial}{\partial y^i} - M_\mu^0 \frac{\partial}{\partial \lambda}, \quad \frac{\delta}{\delta y^i} = \frac{\partial}{\partial y^i} - L_i^0 \frac{\partial}{\partial \lambda}, \quad \frac{\partial}{\partial \lambda} = \frac{\partial}{\partial y^0}$$

Received November 5, 1994 (EH paper).

* The work was presented in the 3rd International Conference of Tensor Society on Differential Geometry and Its Applications held at Athenes, Greece, Aug. 15-20, 1994.

1) Numbers in brackets refer to the references at the end of the paper.

The transformation of local bases yields the relations:

$$\frac{\delta}{\delta x^\lambda} = \frac{\partial \bar{x}^k}{\partial x^\lambda} \frac{\delta}{\delta \bar{x}^k}, \quad \frac{\delta}{\delta y^i} = \frac{\partial \bar{x}^j}{\partial x^i} \frac{\delta}{\delta \bar{y}^j}$$

$$d\bar{x}^k = \frac{\partial \bar{x}^k}{\partial x^\lambda} dx^\lambda, \quad \delta \bar{y}^i = \frac{\partial \bar{x}^i}{\partial x^j} \delta y^j, \quad \lambda = \bar{\lambda}.$$

Considering that in a deformed bundle $T^{(2)}(DF)$, the D -connections $D\Gamma = \{L_{\lambda\mu}^\nu, L_{i\mu}^j, L_{0\mu}^0, C_{ij}^m, C_{\mu j}^\nu, C_{0j}^0, \Lambda_{\mu 0}^\nu, \Lambda_{i0}^j, \Lambda_{00}^0\}$ become d -connections, then we get

$$D_{\delta/\delta x^\mu} \frac{\delta}{\delta x^\lambda} = L_{\lambda\mu}^\nu \frac{\delta}{\delta x^\nu}, \quad D_{\delta/\delta x^\mu} \frac{\delta}{\delta y^i} = L_{i\mu}^j \frac{\delta}{\delta y^j}, \quad D_{\delta/\delta x^\mu} \frac{\partial}{\partial \lambda} = L_{0\mu}^0 \frac{\partial}{\partial \lambda},$$

$$D_{\delta/\delta y^j} \frac{\delta}{\delta x^\mu} = C_{\mu j}^\nu \frac{\delta}{\delta x^\nu}, \quad D_{\delta/\delta y^j} \frac{\delta}{\delta y^i} = C_{ij}^m \frac{\delta}{\delta y^m}, \quad D_{\delta/\delta y^j} \frac{\partial}{\partial \lambda} = C_{0j}^0 \frac{\partial}{\partial \lambda},$$

$$D_{\partial/\partial \lambda} \frac{\delta}{\delta x^\mu} = \Lambda_{\mu 0}^\nu \frac{\delta}{\delta x^\nu}, \quad D_{\partial/\partial \lambda} \frac{\delta}{\delta y^i} = \Lambda_{i0}^j \frac{\delta}{\delta y^j}, \quad D_{\partial/\partial \lambda} \frac{\partial}{\partial \lambda} = \Lambda_{00}^0 \frac{\partial}{\partial \lambda}.$$

In the framework of our consideration, the differential of fundamental function $F(x, y, \lambda)$ of the family of spaces is given, for example, by the relation,

$$dF(x, y, \lambda) = \frac{\partial F}{\partial x^k} dx^k + \frac{\partial F}{\partial y^i} dy^i + \frac{\partial F}{\partial \lambda} d\lambda = \frac{\delta F}{\delta x^k} \delta x^k + \frac{\delta F}{\delta y^i} \delta y^i + \frac{\partial F}{\partial \lambda} d\lambda.$$

The absolute differential of an arbitrary vector X^i has the form:

$$DX^i = dX^i + L_{j\mu}^i(x, y, \lambda) X^j \delta x^\mu + C_{jk}^i(x, y, \lambda) X^j \delta y^k + \Lambda_{j0}^i X^j d\lambda,$$

where

$$L_{j\mu}^i = \frac{1}{2} g^{ir} \left(\frac{\delta g_{rj}}{\delta x^\mu} + \frac{\delta g_{r\mu}}{\delta x^j} - \frac{\delta g_{j\mu}}{\delta x^r} \right),$$

$$C_{jk}^i = \frac{1}{2} g^{ir} \left(\frac{\delta g_{rj}}{\delta y^k} + \frac{\delta g_{rk}}{\delta y^j} - \frac{\delta g_{jk}}{\delta y^r} \right),$$

$$\Lambda_{j0}^i = \frac{1}{2} g^{il} \left(\frac{\delta g_{\ell 0}}{\delta y^j} + \frac{\partial g_{\ell j}}{\partial \lambda} - \frac{\delta g_{j0}}{\delta y^\ell} \right).$$

For the physical viewpoint, it is useful to impose the relation $g_{\ell 0} = g_{j0} = 0$. Also, the form of a curve $\tilde{C} : [0, 1] \rightarrow T^{(2)}(DF)$ is given locally by $X^A(t) = (x^k(t), y^i(t), \lambda(t))$. A curve \tilde{C} in $T^{(2)}(DF)$ is a geodesic with respect to the global affine connection $\tilde{\Gamma}_{BC}^A$, if and only if it satisfies the relation

$$\frac{d^2 X^A}{ds^2} + \tilde{\Gamma}_{BC}^A \frac{dX^B}{ds} \frac{dX^C}{ds} = 0.$$

§ 2. **Torsions and curvatures.** We shall calculate the torsions and the curvatures in a $T^{(2)}(DF)$, taking into account [2], [7] in our case. The torsion tensor field T of a d -connection D is given by $T(X, Y) = D_X Y - D_Y X - [X, Y]$, $\forall X, Y$. For the adapted frame $\left(\frac{\delta}{\delta x^\lambda}, \frac{\delta}{\delta y^i}, \frac{\partial}{\partial \lambda}\right)$, we have the relations:

$$(2.1) \quad \begin{aligned} a) \quad T\left(\frac{\delta}{\delta x^\mu}, \frac{\delta}{\delta x^\nu}\right) &= T_{\nu\mu}^\lambda \frac{\delta}{\delta x^\lambda} + T_{\nu\mu}^i \frac{\delta}{\delta y^i} + T_{\nu\mu}^0 \frac{\partial}{\partial \lambda}, \\ b) \quad T\left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta x^\mu}\right) &= T_{\nu\mu}^\lambda \frac{\delta}{\delta x^\lambda} + T_{\nu\mu}^i \frac{\delta}{\delta y^i} + T_{\nu\mu}^0 \frac{\partial}{\partial \lambda}, \\ c) \quad T\left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta y^j}\right) &= T_{ji}^\lambda \frac{\delta}{\delta x^\lambda} + T_{ji}^e \frac{\delta}{\delta y^e} + T_{ji}^0 \frac{\partial}{\partial \lambda}, \\ d) \quad T\left(\frac{\delta}{\delta x^\mu}, \frac{\partial}{\partial \lambda}\right) &= T_{\mu 0}^\lambda \frac{\delta}{\delta x^\lambda} + T_{\mu 0}^i \frac{\delta}{\delta y^i} + T_{\mu 0}^0 \frac{\partial}{\partial \lambda}, \\ e) \quad T\left(\frac{\delta}{\delta y^i}, \frac{\partial}{\partial \lambda}\right) &= T_{i0}^\lambda \frac{\delta}{\delta x^\lambda} + T_{i0}^j \frac{\delta}{\delta y^j} + T_{i0}^0 \frac{\partial}{\partial \lambda}, \\ f) \quad T\left(\frac{\partial}{\partial \lambda}, \frac{\partial}{\partial \lambda}\right) &= T_{00}^\lambda \frac{\delta}{\delta x^\lambda} + T_{00}^j \frac{\delta}{\delta y^j} + T_{00}^0 \frac{\partial}{\partial \lambda}. \end{aligned}$$

The relation (2.1)a) is written as

$$(2.2) \quad \begin{aligned} D_{\delta/\delta x^\mu} \frac{\delta}{\delta x^\nu} - D_{\delta/\delta x^\nu} \frac{\delta}{\delta x^\mu} - \left[\frac{\delta}{\delta x^\mu}, \frac{\delta}{\delta x^\nu} \right] \\ = L_{\nu\mu}^\lambda \frac{\delta}{\delta x^\lambda} - L_{\mu\nu}^\lambda \frac{\delta}{\delta x^\lambda} - R_{\mu\nu}^i \frac{\delta}{\delta y^i} - \tilde{U}_{\mu\nu}^0 \frac{\partial}{\partial \lambda}, \end{aligned}$$

which implies via the relations

$$\left[\frac{\delta}{\delta x^\mu}, \frac{\delta}{\delta x^\nu} \right] = \left(\frac{\delta N_\mu^i}{\delta x^\nu} - \frac{\delta N_\nu^i}{\delta x^\mu} \right) \frac{\partial}{\partial y^i} + \left(\frac{\delta M_\mu^0}{\delta x^\nu} - \frac{\delta M_\nu^0}{\delta x^\mu} \right) \frac{\partial}{\partial \lambda} = R_{\mu\nu}^i \frac{\partial}{\partial y^i} + U_{\mu\nu}^0 \frac{\partial}{\partial \lambda}$$

and

$$\tilde{U}_{\mu\nu}^0 = U_{\mu\nu}^0 + R_{\mu\nu}^i L_i^0.$$

Also comparing the relations (2.1)a) and (2.2), we obtain for the torsions

$$T_{\nu\mu}^\lambda = L_{\nu\mu}^\lambda - L_{\mu\nu}^\lambda, \quad T_{\nu\mu}^i = -R_{\mu\nu}^i, \quad T_{\nu\mu}^0 = -\tilde{U}_{\mu\nu}^0.$$

Similarly, from the relation (2.1)b), we have

$$T\left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta x^\mu}\right) = D_{\delta/\delta y^i} \frac{\delta}{\delta x^\mu} - D_{\delta/\delta x^\mu} \frac{\delta}{\delta y^i} - \left[\frac{\delta}{\delta y^i}, \frac{\delta}{\delta x^\mu} \right]$$

$$= C_{i\mu}^\lambda \frac{\delta}{\delta x^\lambda} - L_{i\mu}^j \frac{\delta}{\delta y^j} + \frac{\delta N_\mu^j}{\delta y^i} \frac{\delta}{\delta y^j} + \tilde{Y}_{\mu i}^0 \frac{\partial}{\partial \lambda},$$

$$T_{i\mu}^\lambda = C_{i\mu}^\lambda, \quad T_{i\mu}^j = \frac{\delta N_\mu^j}{\delta y^i} - L_{i\mu}^j, \quad T_{i\mu}^0 = -\tilde{Y}_{\mu i}^0,$$

where

$$\left[\frac{\delta}{\delta y^i}, \frac{\delta}{\delta x^\mu} \right] = \frac{\delta N_\mu^j}{\delta y^i} \frac{\delta}{\delta y^j} + \tilde{Y}_{\mu i}^0 \frac{\partial}{\partial \lambda}.$$

From (2.1)c), we get

$$T \left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta y^j} \right) = D_{\delta/\delta y^i} \frac{\delta}{\delta y^j} - D_{\delta/\delta y^j} \frac{\delta}{\delta y^i} - \left[\frac{\delta}{\delta y^i}, \frac{\delta}{\delta y^j} \right] = C_{ji}^\ell \frac{\delta}{\delta y^\ell} - C_{ij}^\ell \frac{\delta}{\delta y^\ell} - R_{ij}^0 \frac{\partial}{\partial \lambda}.$$

So we have

$$T_{ji}^\lambda = 0, \quad T_{ji}^\ell = C_{ji}^\ell - C_{ij}^\ell, \quad T_{ji}^0 = -R_{ij}^0,$$

where

$$\left[\frac{\delta}{\delta y^i}, \frac{\delta}{\delta y^j} \right] = R_{ij}^0 \frac{\partial}{\partial \lambda}.$$

The relation (2.1)d) is written in the form

$$\begin{aligned} T \left(\frac{\delta}{\delta x^\mu}, \frac{\partial}{\partial \lambda} \right) &= T_{\mu 0}^\lambda \frac{\delta}{\delta x^\lambda} + T_{\mu 0}^i \frac{\delta}{\delta y^i} + T_{\mu 0}^0 \frac{\partial}{\partial \lambda} \\ &= D_{\delta/\delta x^\mu} \frac{\partial}{\partial \lambda} - D_{\partial/\partial \lambda} \frac{\delta}{\delta x^\mu} - \left[\frac{\delta}{\delta x^\mu}, \frac{\partial}{\partial \lambda} \right] \\ &= L_{0\mu}^0 \frac{\partial}{\partial \lambda} - \Lambda_{0\mu}^\lambda \frac{\delta}{\delta x^\lambda} - \frac{\partial N_\mu^i}{\partial \lambda} \frac{\delta}{\delta y^i} - Z_\mu^0 \frac{\partial}{\partial \lambda}. \end{aligned}$$

Consequently, it follows:

$$T_{\mu 0}^\lambda = -\Lambda_{0\mu}^\lambda, \quad T_{\mu 0}^i = -\frac{\partial N_\mu^i}{\partial \lambda}, \quad T_{\mu 0}^0 = L_{0\mu}^0 - Z_\mu^0,$$

where

$$\left[\frac{\delta}{\delta x^\mu}, \frac{\partial}{\partial \lambda} \right] = \frac{\partial N_\mu^i}{\partial \lambda} \frac{\partial}{\partial y^i} + Z_\mu^0 \frac{\partial}{\partial \lambda} \quad \text{and} \quad Z_\mu^0 = \frac{\partial M_\mu^0}{\partial \lambda} + \frac{\partial N_\mu^i}{\partial \lambda} L_i^0.$$

From (2.1)e), we have

$$\begin{aligned} T \left(\frac{\delta}{\delta y^i}, \frac{\partial}{\partial \lambda} \right) &= T_{i 0}^\lambda \frac{\delta}{\delta x^\lambda} + T_{i 0}^j \frac{\delta}{\delta y^j} + T_{i 0}^0 \frac{\partial}{\partial \lambda} \\ &= D_{\delta/\delta y^i} \frac{\partial}{\partial \lambda} - D_{\partial/\partial \lambda} \frac{\delta}{\delta y^i} - \left[\frac{\delta}{\delta y^i}, \frac{\partial}{\partial \lambda} \right] \\ &= C_{0i}^0 \frac{\partial}{\partial \lambda} - \Lambda_{0i}^j \frac{\delta}{\delta y^j} - \frac{\partial L_i^0}{\partial \lambda} \frac{\partial}{\partial \lambda}. \end{aligned}$$

$$T_{i0}^\lambda = 0, \quad T_{i0}^j = -\Lambda_{0i}^j, \quad T_{i0}^0 = C_{0i}^0 - \frac{\partial L_i^0}{\partial \lambda},$$

where

$$\left[\frac{\delta}{\delta y^i}, \frac{\partial}{\partial \lambda} \right] = \frac{\partial L_i^0}{\partial \lambda} \frac{\partial}{\partial \lambda}.$$

Finally, from (2.1f), we obtain $T_{00}^\lambda = T_{00}^i = T_{00}^0 = 0$.

The curvature tensors and the λ -curvature tensors $I_{\nu 0 \mu}^k, I_{i 0 \mu}^j, I_{0 0 \mu}^0, J_{\nu 0 i}^k, J_{j 0 i}^\ell, J_{0 0 i}^0$ in $T^{(2)}$ (DF) are given by the relations:

$$\begin{aligned} R\left(\frac{\delta}{\delta x^\mu}, \frac{\delta}{\delta x^\nu}\right) \frac{\delta}{\delta x^\lambda} &= R_{\lambda\nu\mu}^\kappa \frac{\delta}{\delta x^\kappa}, & R\left(\frac{\delta}{\delta x^\mu}, \frac{\delta}{\delta x^\nu}\right) \frac{\delta}{\delta y^i} &= R_{i\nu\mu}^j \frac{\delta}{\delta y^j}, \\ R\left(\frac{\delta}{\delta x^\mu}, \frac{\delta}{\delta x^\nu}\right) \frac{\partial}{\partial \lambda} &= R_{0\nu\mu}^0 \frac{\partial}{\partial \lambda}, & R\left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta x^\mu}\right) \frac{\delta}{\delta x^\nu} &= P_{\nu\mu i}^\kappa \frac{\delta}{\delta x^\kappa}, \\ R\left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta x^\mu}\right) \frac{\delta}{\delta y^j} &= P_{ji\mu}^\ell \frac{\delta}{\delta y^\ell}, & R\left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta x^\mu}\right) \frac{\partial}{\partial \lambda} &= P_{0i\mu}^0 \frac{\partial}{\partial \lambda}, \\ R\left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta y^j}\right) \frac{\delta}{\delta x^\mu} &= S_{\mu ji}^\kappa \frac{\delta}{\delta x^\kappa}, & R\left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta y^j}\right) \frac{\delta}{\delta y^\ell} &= S_{\ell ji}^m \frac{\delta}{\delta y^m}, \\ R\left(\frac{\delta}{\delta y^i}, \frac{\delta}{\delta y^j}\right) \frac{\partial}{\partial \lambda} &= S_{0ji}^0 \frac{\partial}{\partial \lambda}, & R\left(\frac{\delta}{\delta x^\mu}, \frac{\partial}{\partial \lambda}\right) \frac{\delta}{\delta x^\nu} &= I_{\nu 0 \mu}^k \frac{\delta}{\delta x^k}, \\ R\left(\frac{\delta}{\delta x^\mu}, \frac{\partial}{\partial \lambda}\right) \frac{\delta}{\delta y^i} &= I_{i 0 \mu}^j \frac{\delta}{\delta y^j}, & R\left(\frac{\delta}{\delta x^\mu}, \frac{\partial}{\partial \lambda}\right) \frac{\partial}{\partial \lambda} &= I_{0 0 \mu}^0 \frac{\partial}{\partial \lambda}, \\ R\left(\frac{\delta}{\delta y^i}, \frac{\partial}{\partial \lambda}\right) \frac{\delta}{\delta x^\nu} &= J_{\nu 0 i}^k \frac{\delta}{\delta x^k}, & R\left(\frac{\delta}{\delta y^i}, \frac{\partial}{\partial \lambda}\right) \frac{\delta}{\delta y^j} &= J_{j 0 i}^\ell \frac{\delta}{\delta y^\ell}, \\ R\left(\frac{\delta}{\delta y^i}, \frac{\partial}{\partial \lambda}\right) \frac{\partial}{\partial \lambda} &= J_{0 0 i}^0 \frac{\partial}{\partial \lambda}, & R\left(\frac{\partial}{\partial \lambda}, \frac{\partial}{\partial \lambda}\right) \frac{\partial}{\partial \lambda} &= 0 \\ R\left(\frac{\partial}{\partial \lambda}, \frac{\partial}{\partial \lambda}\right) \frac{\delta}{\delta x^\mu} &= K_{\mu 0 0}^k \frac{\delta}{\delta x^\mu} = 0, & R\left(\frac{\partial}{\partial \lambda}, \frac{\partial}{\partial \lambda}\right) \frac{\delta}{\delta y^i} &= K_{i 0 0}^j \frac{\delta}{\delta y^i} = 0. \end{aligned}$$

The curvature tensor field R of a d -connection D has the form: $R(X, Y)Z = [D_x, D_y]Z - D_{[x, y]}Z, \forall X, Y, Z$. We apply this form the adapted frame, to produce:

$$\begin{aligned} R_{\lambda\nu\mu}^k \frac{\delta}{\delta x^k} &= R\left(\frac{\delta}{\delta x^\mu}, \frac{\delta}{\delta x^\nu}\right) \frac{\delta}{\delta x^\lambda} \\ &= D_{\delta/\delta x^\mu} \left(D_{\delta/\delta x^\nu} \frac{\delta}{\delta x^\lambda} \right) - D_{\delta/\delta x^\nu} \left(D_{\delta/\delta x^\mu} \frac{\delta}{\delta x^\lambda} \right) - D_{[\delta/\delta x^\mu, \delta/\delta x^\nu]} \frac{\delta}{\delta x^\lambda} \\ &= D_{\delta/\delta x^\mu} \left(L_{\lambda\nu}^k \frac{\delta}{\delta x^k} \right) - D_{\delta/\delta x^\nu} \left(L_{\lambda\mu}^k \frac{\delta}{\delta x^k} \right) - D_{R_{\mu\nu}^i \delta/\delta y^i + \tilde{U}_{\mu\nu}^0 \partial/\partial \lambda} \frac{\delta}{\delta x^\lambda}, \\ \frac{\delta L_{\lambda\nu}^k}{\delta x^\mu} \frac{\delta}{\delta x^\mu} + L_{\lambda\nu}^k D_{\delta/\delta x^\mu} \frac{\delta}{\delta x^k} - \frac{\delta L_{\lambda\mu}^k}{\delta x^\nu} - L_{\lambda\mu}^k D_{\delta/\delta x^\nu} \frac{\delta}{\delta x^k} - R_{\mu\nu}^i D_{\delta/\delta y^i} \frac{\delta}{\delta x^\lambda} - \tilde{U}_{\mu\nu}^0 D_{\partial/\partial \lambda} \frac{\delta}{\delta x^\lambda} \end{aligned}$$

$$= \frac{\delta L_{\lambda\nu}^k}{\delta x^\mu} \frac{\delta}{\delta x^k} + L_{\lambda\nu}^\rho L_{\rho\mu}^k \frac{\delta}{\delta x^k} - \frac{\delta L_{\lambda\rho}^k}{\delta x^\nu} \frac{\delta}{\delta x^k} - L_{\lambda\mu}^\rho L_{\rho\nu}^k \frac{\delta}{\delta x^k} - R_{\mu\nu}^i C_{i\lambda}^k \frac{\delta}{\delta x^k} - \tilde{U}_{\mu\nu}^0 \Lambda_{0\lambda}^k \frac{\delta}{\delta x^k}.$$

So, coefficients of the curvature of a d -connection D take the following form:

$$(2.3) \quad R_{\lambda\nu\mu}^k = \frac{\delta L_{\lambda\nu}^k}{\delta x^\mu} - \frac{\delta L_{\lambda\mu}^k}{\delta x^\nu} + L_{\lambda\nu}^\rho L_{\rho\mu}^k - L_{\lambda\mu}^\rho L_{\rho\nu}^k - R_{\mu\nu}^i C_{i\lambda}^k - \tilde{U}_{\mu\nu}^0 \Lambda_{0\lambda}^k,$$

where $\tilde{U}_{\mu\nu}^0$ is given by the relation (2.2).

With a similar method, we can calculate the other types of curvatures.

$$(2.4) \quad R_{i\nu\mu}^j = \frac{\delta L_{i\nu}^j}{\delta x^\mu} - \frac{\delta L_{i\mu}^j}{\delta x^\nu} + L_{i\nu}^\ell L_{\ell\mu}^j - L_{i\mu}^\ell L_{\ell\nu}^j - R_{\mu\nu}^\ell C_{i\ell}^j - \tilde{U}_{\mu\nu}^0 \Lambda_{0i}^j,$$

$$(2.5) \quad R_{0\nu\mu}^0 = \frac{\delta L_{0\nu}^0}{\delta x^\mu} - \frac{\delta L_{0\mu}^0}{\delta x^\nu} + L_{0\nu}^0 L_{0\mu}^0 - L_{0\mu}^0 L_{0\nu}^0 - R_{\mu\nu}^i C_{0i}^0 - \tilde{U}_{\mu\nu}^0 \Lambda_{00}^0,$$

$$(2.6) \quad P_{ji\mu}^\ell = \frac{\delta L_{j\mu}^\ell}{\delta y^i} - \frac{\delta C_{ji}^\ell}{\delta x^\mu} + L_{j\mu}^n C_{ni}^\ell - C_{ji}^n L_{n\mu}^\ell + \frac{\delta N_{\mu}^n}{\delta y^i} C_{jn}^\ell + \tilde{Y}_{\mu i}^0 \Lambda_{0j}^\ell,$$

where

$$\tilde{Y}_{\mu i}^0 = Y_{\mu i}^0 + \frac{\delta N_{\mu}^j}{\delta y^i} L_j^0, \quad Y_{\mu i}^0 = \frac{\delta M_{\mu}^0}{\delta y^i} - \frac{\delta L_i^0}{\delta x^\mu}.$$

$$(2.7) \quad P_{0i\mu}^0 = \frac{\delta L_{0\mu}^0}{\delta y^i} - \frac{\delta C_{0i}^0}{\delta x^\mu} + L_{0\mu}^0 C_{0i}^0 - C_{0i}^0 L_{0\mu}^0 + \frac{\delta N_{\mu}^j}{\delta y^i} C_{0j}^0 + \tilde{Y}_{\mu i}^0 \Lambda_{00}^0,$$

$$(2.8) \quad S_{\mu ji}^k = \frac{\delta C_{j\mu}^k}{\delta y^i} - \frac{\delta C_{i\mu}^k}{\delta y^j} + C_{j\mu}^\lambda C_{\lambda i}^k - C_{i\mu}^\lambda C_{\lambda j}^k - R_{ij}^0 \Lambda_{0\mu}^k,$$

$$(2.9) \quad S_{\ell ji}^m = \frac{\delta C_{\ell j}^m}{\delta y^i} - \frac{\delta C_{\ell i}^m}{\delta y^j} + C_{\ell j}^n C_{ni}^m - C_{\ell i}^n C_{nj}^m - R_{ij}^0 \Lambda_{0\ell}^m,$$

$$(2.10) \quad S_{0ji}^0 = \frac{\delta C_{0j}^0}{\delta y^i} - \frac{\delta C_{0i}^0}{\delta y^j} + C_{0j}^0 C_{0i}^0 - C_{0i}^0 C_{0j}^0 - R_{ij}^0 \Lambda_{00}^0,$$

$$(2.11) \quad I_{\nu 0\mu}^k = \frac{\delta \Lambda_{0\nu}^k}{\delta x^\mu} - \frac{\partial L_{\nu\mu}^k}{\partial \lambda} + \Lambda_{0\nu}^\rho L_{\rho\mu}^k - L_{\nu\mu}^\rho \Lambda_{0\rho}^k - \frac{\partial N_{\mu}^i}{\partial \lambda} C_{i\nu}^k - Z_{\mu}^0 \Lambda_{0\nu}^k,$$

where

$$Z_{\mu}^0 = \frac{\partial M_{\mu}^0}{\partial \lambda} + \frac{\partial N_{\mu}^i}{\partial \lambda} L_i^0.$$

$$(2.12) \quad I_{i0\mu}^j = \frac{\delta \Lambda_{0i}^j}{\delta x^\mu} - \frac{\partial L_{i\mu}^j}{\partial \lambda} + \Lambda_{0i}^\ell L_{\ell\mu}^j - L_{i\mu}^\ell \Lambda_{0\ell}^j - \frac{\partial N_{\mu}^\ell}{\partial \lambda} C_{i\ell}^j - Z_{\mu}^0 \Lambda_{0i}^j,$$

$$(2.13) \quad I_{00\mu}^0 = \frac{\delta \Lambda_{00}^0}{\delta x^\mu} - \frac{\partial L_{0\mu}^0}{\partial \lambda} + \Lambda_{00}^0 L_{0\mu}^0 - L_{0\mu}^0 \Lambda_{00}^0 - \frac{\partial N_{\mu}^\ell}{\partial \lambda} C_{0\ell}^0 - Z_{\mu}^0 \Lambda_{00}^0,$$

$$(2.14) \quad J_{\nu 0i}^k = \frac{\delta \Lambda_{0\nu}^k}{\delta y^i} - \frac{\partial C_{i\nu}^k}{\partial \lambda} + \Lambda_{0\nu}^\rho C_{i\rho}^k - C_{i\nu}^\rho \Lambda_{0\rho}^k - \frac{\partial L_i^0}{\partial \lambda} \Lambda_{0\nu}^k,$$

$$(2.15) \quad J_{j0i}^\ell = \frac{\delta \Lambda_{0j}^\ell}{\delta y^i} - \frac{\partial C_{ji}^\ell}{\partial \lambda} + \Lambda_{0j}^n C_{ni}^\ell - C_{ji}^n \Lambda_{0n}^\ell - \frac{\partial L_i^0}{\partial \lambda} \Lambda_{0j}^\ell,$$

$$(2.16) \quad J_{00i}^0 = \frac{\delta \Lambda_{00}^0}{\delta y^i} - \frac{\partial C_{0i}^0}{\partial \lambda} + \Lambda_{00}^0 C_{0i}^0 - C_{0i}^0 \Lambda_{00}^0 - \frac{\partial L_i^0}{\partial \lambda} \Lambda_{00}^0,$$

$$(2.17) \quad K_{\mu 00}^\nu = K_{i00}^j = K_{000}^0 = 0.$$

Thus we can get the following:

Theorem 2.1. In a space $T^{(2)}$ (DF), the relations (2.3)-(2.17) constitute the coefficients of the curvatures of a d -connection D .

§ 3. Conclusions. From a physical point of view, in order to emphasize the physical effects caused by the λ -dependent quantities, we had better focus our attention on such terms as the non-linear connections M_μ^0 , L_i^0 and the connection coefficients $\Lambda_{\mu 0}^\nu$, Λ_{i0}^j . Therefore, we had better examine the λ -torsion tensors $T_{\mu 0}^\lambda$, $T_{\mu 0}^i$, $T_{\mu 0}^0$, T_{i0}^j , T_{i0}^0 and the λ -curvature tensors $I_{\nu 0 \mu}^k$, $I_{i0 \mu}^j$, $I_{00 \mu}^0$, $J_{\nu 0 i}^k$, $J_{j0 i}^\ell$, $J_{00 i}^0$. The physical meaning of them, is of course, dependent on the physical meaning of λ itself. Therefore, for example, with respect to the non-linear connections M_μ^0 , L_i^0 , the λ -deformation field can signify the rate of change of the interaction's values, or the differentiation of the intensity of gravitational field relative to the above mentioned curvature tensors.

Consequently, the existence of λ -field will influence the intensity of tidal forces via of λ -curvature tensors.

The last consideration will be the object of a further our study.

Department of Mathematics
University of Athens
15 784 Panepistimiopolis
Athens, Greece

and

Department of Mechanical
Engineering
Faculty of Science and Technology
Science University of Tokyo
Noda, Chiba 278, Japan

REFERENCES

- [1] M. Anastasiei and H. Kawaguchi : A geometrical theory of time dependent Lagrangians I, II, *Tensor, N.S.*, **48** (1989), 273-282 and 283-293.
- [2] R. Miron, S. Watanabe and S. Ikeda : Some connections on tangent bundle and their applications to the general relativity, *Tensor, N.S.*, **46** (1987), 8-22.
- [3] S. Ikeda : On the theory of gravitational field non-localized by the internal variable IV, *Il Nuovo Cimento*, **108 B** (1993), 397-401.
- [4] S. Ikeda : On the theory of gravitational field in Finsler spaces, *Tensor, N.S.*, **50** (1991), 256-262.
- [5] P.C. Stavrinou : A varying space-time hypersurface with finslerian geometry, 10th Int. Conf. on General Relativity and Gravitation. Contribution, *Padova*, 4-9 July 1983, 638-640.

- [6] R. Miron and M. Anastasiei : The geometry of Lagrange spaces: Theory and applications, *Kluwer*, 1994.
- [7] K. Yano and S. Ishihara : Tangent and cotangent bundles, *Marc. Dekker, Inc.*, *New York*, 1973.

