

Exact results in gauge theories

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ΗΜΕΡΙΔΑ ΠΡΟΣ ΤΙΜΗΝ ΤΟΥ ΦΩΚΙΩΝΑ ΧΑΤΖΗΩΑΝΝΟΥ

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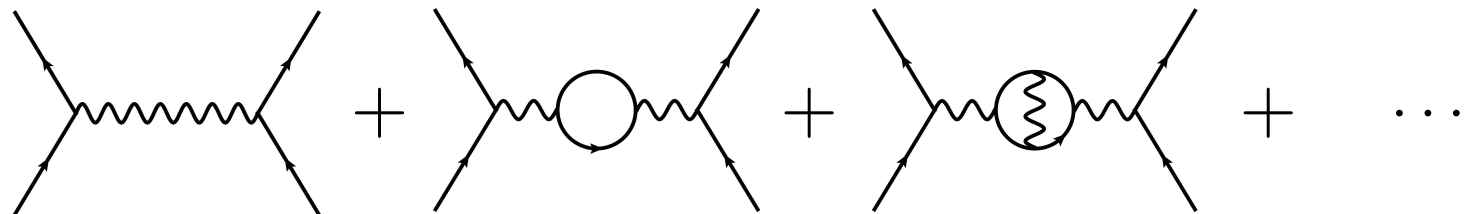
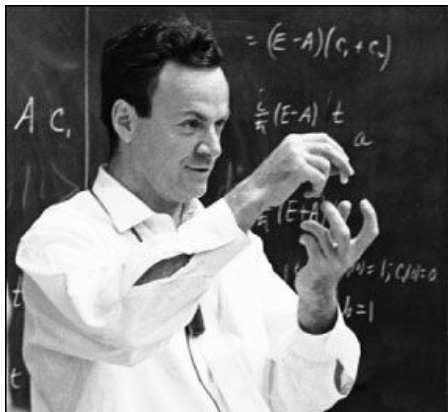
Gauge theories

QUARKS	u up	c charm	t top	γ photon
	d down	s strange	b bottom	g gluon
LEPTONS	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	Z^0 Z boson
	e electron	μ muon	τ tau	W^\pm W boson
	FORCE CARRIERS			

U(N) or SU(N) gauge group
gauge bosons force carriers: A_μ

Matter fields (quarks and leptons): ψ

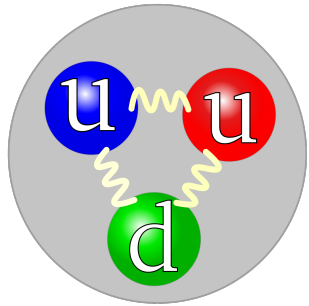
Feynman: how to calculate a process **perturbatively**:



$$= c_1 \lambda + c_2 \lambda^2 + c_3 \lambda^3 + \dots$$

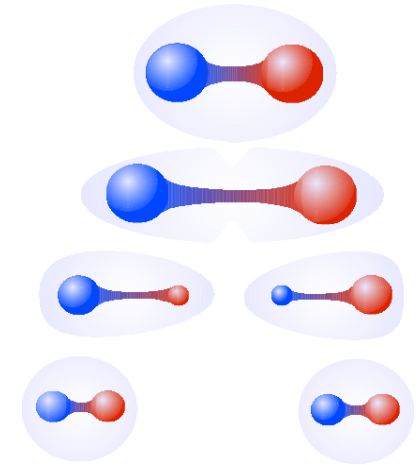
$$\lambda \ll 1$$

A big open problem



Confinement:

Quarks cannot be isolated in Nature, and cannot be directly observed.



Find an analytic proof that quantum chromodynamics (QCD) should be confining.

Cannot be done perturbatively

Emmy Noether



Symmetry → *Conservation law*

Use the Symmetry to solve the problem.

***The more symmetry the easier it
is to solve the problem.***

Gauge theories are very hard to understand:
Let's add **Supersymmetry**.

boson ↔ **fermion**

Supersymmetric Gauge theories in 4D

$\mathcal{N} = 4$ Super-Yang-Mills (SYM)

$$A_\mu, \lambda^A, \phi^{AB} \quad A, B = 1, \dots, 4$$

No matter fields are allowed! Only “gluons”!
Conformal: $\lambda \neq \lambda(E)$!

*More susy easier
to calculate*

$$\mathcal{N} = 2 \quad \left\{ \begin{array}{ll} \text{Vector multiplet} & A_\mu, \lambda^{\mathcal{I}}, \phi \\ \text{Matter multiplet} & Q_{\mathcal{I}}, \psi, \tilde{\psi} \end{array} \right. \quad \mathcal{I} = 1, 2$$

*Less susy more
realistic*

$$\mathcal{N} = 1 \quad \left\{ \begin{array}{ll} \text{Vector} & A_\mu, \lambda \\ \text{Matter} & q, \psi \end{array} \right.$$

$$\mathcal{N} = 0 \quad \left\{ \begin{array}{ll} \text{Vector} & A_\mu \\ \text{Matter} & \psi \end{array} \right.$$

The real world

The success story

Exact results for many observables

$$c_1 \lambda + c_2 \lambda^2 + c_3 \lambda^3 + \dots = F(\lambda)$$

due to:

- * *AdS/CFT correspondence*
- * *Integrability*
- * *Localization*
- * *4D/2D relations*

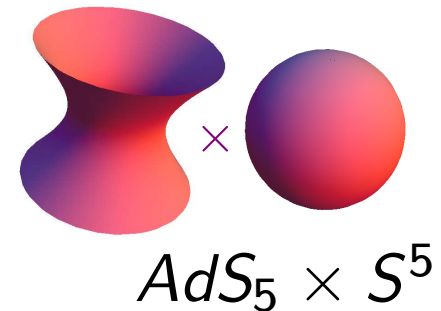
AdS/CFT correspondence

[’t Hooft 1993 Susskind 1995]

A relation: Gravity theories in $d+1$ and gauge theories d -dimensions.

[Maldacena 1998]

The 4D $\mathcal{N} = 4$ SYM: *10D critical string theory on*



A duality:

Use weakly coupled gravity: strongly coupled gauge theory.

Holographic models: strong coupling regime

Geometric description of:

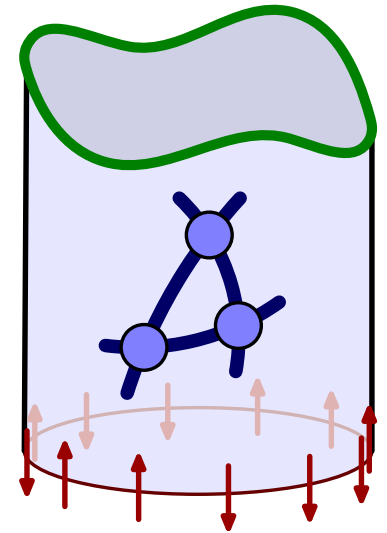
- Confinement
- Chiral symmetry breaking
- Ultrahot QuarkGluonPlasma/ Ultracold atoms
- Applications of Condensed matter systems

Integrability

$\mathcal{N} = 4$ SYM is integrable in the planar limit for **any coupling**

- *Perturbation theory*: mapped to an integrable spin chain
- *Strong coupling*: integrable 2D theory on the string world-sheet

Powerful integrability toolkit



- The spectral problem is solved **exactly**: for **any coupling**

Integrability now is applied to **other observables**.

Localization

[Pestun 2007]

$$Z_{S^4} = \int [D\Phi] e^{-S[\Phi]} = \int da |\mathcal{Z}(a)|^2$$

The **path integral** localizes to an **ordinary integral**
(*Cancelations due to supersymmetry*)

We can do an ordinary integral.

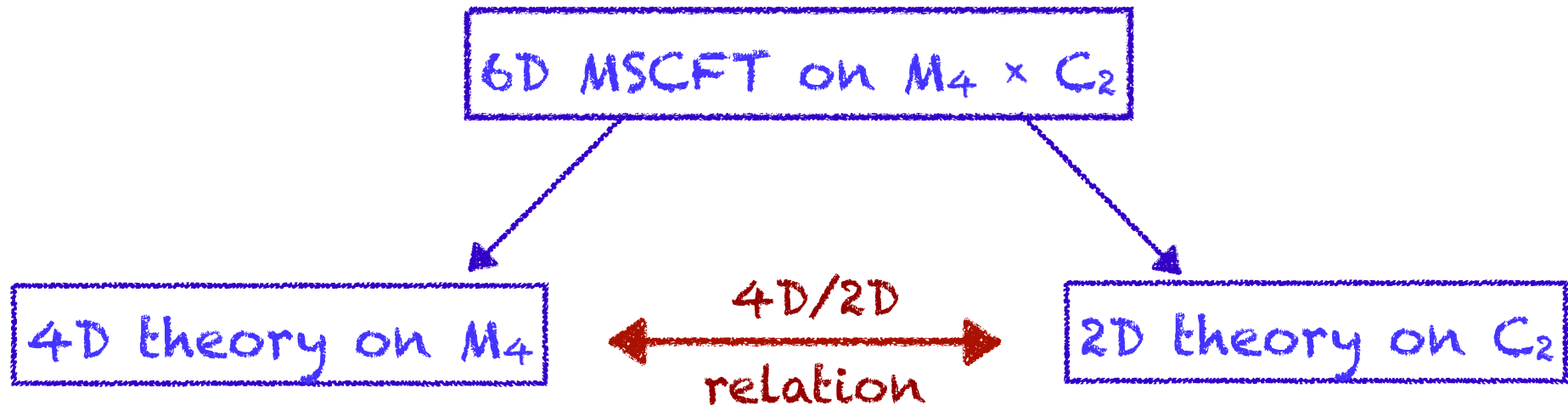
Compute the path integral exactly.

For **any value of the coupling constant**.

Example of exact observable:

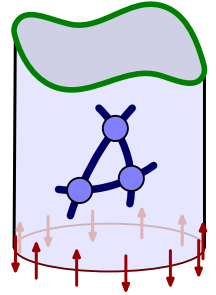
$$W(\lambda) = 2 \frac{I_1(\sqrt{\lambda})}{\sqrt{\lambda}} = \begin{cases} 1 + \frac{\lambda}{8} + \frac{\lambda^2}{192} + \frac{\lambda^3}{9216} + \dots, & \lambda \ll 1 \\ \sqrt{\frac{2}{\pi}} \lambda^{-\frac{3}{4}} e^{\sqrt{\lambda}} + \dots, & \lambda \gg 1 \end{cases}$$

4D/2D relations



- * 4D $N=4$ SYM: 6D MSCFT on T^2 (EM duality = modular transformations)
- * 4D $N=2$ class: 6D MSCFT on Riemann surface $C_{g,n}$ [Gaiotto 2009]
- * 4D partition functions = 2D CFT correlators [Alday, Gaiotto, Tachikawa 2009]
- * 4D SC Index = 2D TFT correlators [Gadde, EP, Rastelli, Razamat 2009]

AdS/CFT integrability beyond $N=4$ SYM



Any $\mathcal{N} = 2$ superconformal gauge theory has a closed **purely gluonic sector** that is **integrable** in the planar limit.

[EP 2013]

[Mitev, EP 2014 + 2015]

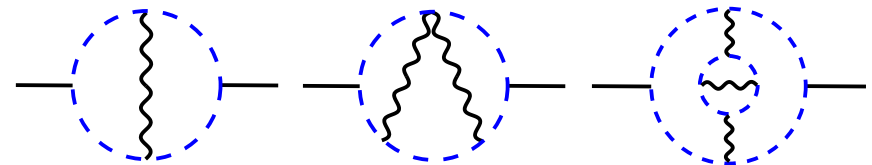
Observable in this sector is obtained from its $\mathcal{N} = 4$ counterpart by replacing: $\lambda \rightarrow f(\lambda)$

Relative finite renormalization: Can calculate it with Feynman diagrams

$$f(\lambda) = \lambda + \lambda(Z_{\mathcal{N}=2} - Z_{\mathcal{N}=4})$$

Compute it exactly using localization.

And check with AdS/CFT.



Lesson:

Gauge symmetry is important: *not* supersymmetry!

Similar sector for $N=1$ theories

[Carstensen, EP work in progress]

Beyond Localization

We can do even better than usual localization:

The partition function for a theory with no Lagrangian description.
Using powerful **string theory** tools.

[Bao, Mitev, EP, Taki, Yagi 2013]

Solving a **long standing open problem in Mathematical physics:**

3-point functions of 2D Toda CFT

[Mitev, EP 2014]

[Isachenkov, Mitev, EP 2014]

[Coman, EP, Teschner to appear]

Carving out a 2D CFT

Can we have 4D/2D relations for $N=1$ theories?

$$\begin{array}{c} \text{4D partition functions} \\ = \\ \text{2D CFT correlators} \end{array}$$

Developing a method for finding such relations.

A large class of $N=1$ theories

[Coman,EP,Taki,Yagi 2015]

* *The 2D symmetry algebra and representations*

[Mitev,EP 2017]

* *Exact results for $N=1$ instantons*

[Bourton, EP 2017]

* *Partition functions for $N=1$ SCFTs*

[work in progress]



Summary (past and future)

- * *AdS/CFT and integrability beyond $N=4$ SYM*
- * *4D/2D relations*
- * *Enlarging the list of Exact $N=2$ observables*
- * *Exact results for $N=1$ theories!*
- * *$N=0$??? (QCD conformal window)*

$N=4$ SYM



*Thanks for your
attention!*

Real world QCD