

# Analysis of the Spontaneous Mass Generation by Using an Iterative Method

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The standard method to discuss the spontaneous mass generation is to formulate a coupled system of self-consistent equations.

- The non-trivial solution is interpreted as a sum of an infinite number of bubble diagrams.
- Does actually the sum of an INFINITE number of diagrams generate the FINITE (non-vanishing) mass?
- Those equations are no more than the *necessary* condition and it is needed to examine solutions to select correct one by using another mean. (e.g. by referring to the free energy of each solution.)

We adopt the Nambu–Jona-Lasinio(NJL) model and give a new iteration method that directly sums up an infinite number of diagrams in the ladder approximation.

Introduce a bare mass

$$\mathcal{L}_{\text{NJL}} = \bar{\psi} \not{\partial} \psi + \frac{2\pi^2 g}{N} \left[ (\bar{\psi} \psi)^2 + (\bar{\psi} i \gamma_5 \psi)^2 \right] - m_0 \bar{\psi} \psi$$

- NJL model has four-fermion interactions among the massless fermions with the chiral invariance.
- Adding the bare mass  $m_0$  to the Lagrangian to make the standard perturbation theory work well.
- $N$  is the number of fermion flavors. We consider  $1/N$  leading contribution to the mass.

## Introduce “node length”

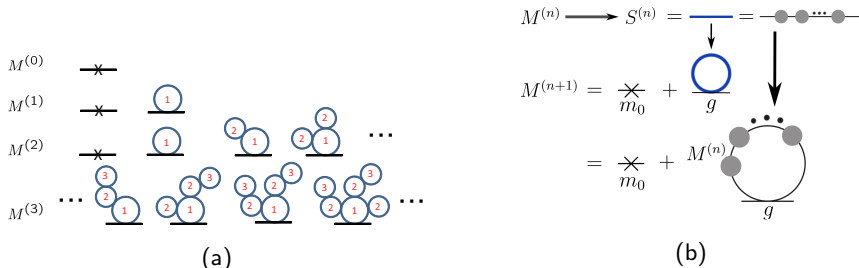


Fig. 1. Node length. (a) Definition. (b) Node length iteration.

- We classify diagrams using the node length of each diagram.
- Node length of a diagram is defined by the maximum number of loops in a continuous route towards the edge loop.
- We define  $M^{(n)}$  is a sum of diagrams whose node length is no greater than  $n$ .

## Transformation function

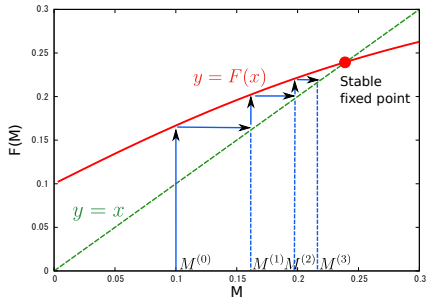
$$M^{(n+1)} = F \left( M^{(n)} \right),$$

$$\text{where } F(M) = m_0 + gM \left( 1 - M^2 \log \left( 1 + M^{-2} \right) \right).$$

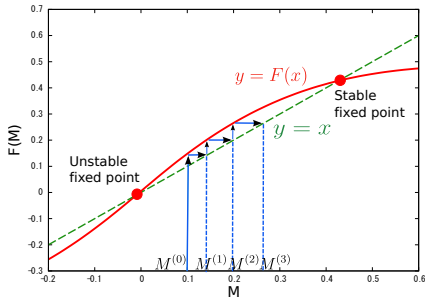
- The transformation function is a one loop integral and we denote it by  $F$ .
- The total sum of the tree diagrams is obtained by  $M^{(\infty)}$ , infinitely many times of transformation of the same  $F$ .

$$\begin{aligned}
 M^{(n)} &\longrightarrow S^{(n)} = \text{---} \text{---} \text{---} \text{---} \text{---} \\
 &\quad \downarrow \quad \quad \quad \downarrow \\
 M^{(n+1)} &= \text{---} \times \text{---} + \frac{\text{---} \circ \text{---}}{g} \\
 &= \text{---} \times \text{---} + \frac{M^{(n)}}{g}
 \end{aligned}$$

# Iterative steps



(a)



(b)

Fig. 2. Iterative steps.

(a) Weak coupling :  $g = 0.7$ . (b) Strong coupling :  $g = 1.5$ .

- In any case the iterative transformation finally reaches a stable fixed point.
- In the weak coupling region, there is only one fixed point.
- In the strong coupling region, there appear pair creation of fixed points.

# Fixed Point Structure

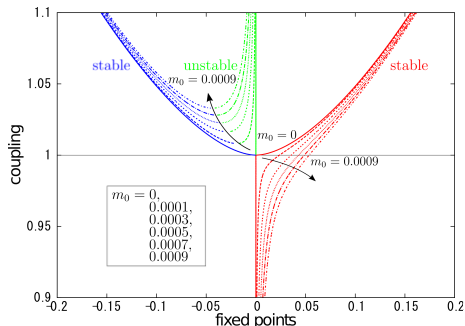
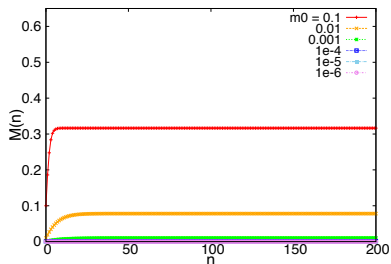


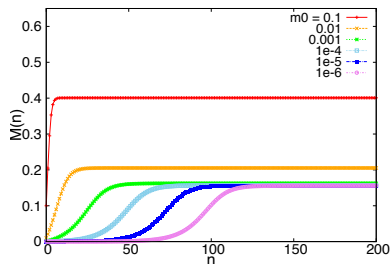
Fig. 3. Fixed point structure with the coupling constant and the bare mass.

- We set a positive value for the bare mass, then the initial point is in the territory of the right-hand side stable fixed point.
- Therefore for all region of the coupling constant, the physical result is controlled by the right most stable fixed point.
- The critical coupling constant for the change of fixed point structure depends on the bare mass and it unifies in the vanishing bare mass limit.

# Mass Generation



(a)



(b)

Fig. 4. Mass generation by the node length iteration.

(a) Weak coupling :  $g = 0.9$ . (b) Strong coupling :  $g = 1.1$ .

- In the weak coupling, the dynamical mass is generated rather quickly at low  $n$  and becomes constant.
- In the strong coupling, the generation of the dynamical mass depends strongly on the bare mass, and it is mainly generated at some narrow range of node length.



- Adding a bare mass  $m_0$  to NJL model, we define the node length iteration.
- Infinite diagrams are summed up explicitly by using iterative method according to the node length.
- The results correctly give the physical mass.



- ▶ Y. Nambu and G. Jona-Lasinio, *Phys. Rev.* **122**, 345 (1961); **124**, 246 (1961).
- ▶ K-I. Aoki, M. Bando, T. Kugo, M.G. Mitchard and H. Nakatani, *Prog. Theor. Phys.* **84**, 683 (1990).
- ▶ K-I. Aoki, S. Onai and D. Sato in preparation.

# Perturbation Theory

We calculate the self-energy  $\Sigma$  in perturbation theory.

$$\boxed{\Sigma} = \begin{array}{c} \text{1-loop} \\ \text{2-loop} \\ \text{3-loop} \\ \dots \end{array}$$

$$\text{Diagram: a circle with momentum } p \text{ inside, sitting on a horizontal line representing an external line.}$$

$$I = \text{Tr} \int d^4 p \frac{1}{\not{p}}$$

The loop part is not affected by the momentum from external line.

$$I = \text{Tr} \Omega_4 \int d^3 p \frac{\not{p}}{p^2}, \quad \text{Tr} \gamma^\mu = 0$$

For the reason of containing at least one  $I$  at the edge of every diagrams, each part becomes zero in the end. It means that all part of diagrams do not receive perturbative contributions.

⇒ Mass cannot be generated in the perturbation theory.

$$\begin{aligned}
 \text{---} \Sigma \text{---} &= \underbrace{\text{---} \bigcirc \text{---} + \text{---} \bigcirc \bigcirc \text{---} + \text{---} \bigcirc \bigcirc \bigcirc \text{---} + \text{---} \bigcirc \bigcirc \bigcirc \text{---} + \dots}_{(1/N)^0} \\
 &+ \underbrace{\text{---} \bigcirc \bigcirc \text{---} + \text{---} \bigcirc \bigcirc \bigcirc \text{---} + \text{---} \bigcirc \bigcirc \bigcirc \text{---} + \dots}_{(1/N)^1} \\
 &+ \dots
 \end{aligned}$$

- A  $1/N$  factor comes out through the interaction at the vertex as the coupling constant is proportional to  $1/N$ .
- On account of the flavor number, a factor  $N$  appears for each loop.