

Framization of the Temperley-Lieb algebra

DIMOS GOUNDAROULIS, JESÚS JUYUMAYA, ARISTIDES KONTOGEORGIS,
AND SOFIA LAMBROPOULOU

We propose a framization of the Temperley-Lieb algebra. The framization is a procedure that can briefly be described as the adding of framing to a known knot algebra in a way that is both algebraically consistent and topologically meaningful. Our framization of the Temperley-Lieb algebra is defined as a quotient of the Yokonuma-Hecke algebra. The main theorem provides necessary and sufficient conditions for the Markov trace defined on the Yokonuma-Hecke algebra to pass through to the quotient algebra. Using this we construct 1-variable invariants for classical knots and links, which, as we show, are not topologically equivalent to the Jones polynomial.

1	Introduction	300
2	Preliminaries	303
3	Fourier transform and the E-system	309
4	A Framization of the Temperley-Lieb algebra	313
5	Markov traces	323
6	Knot and link invariants	331

2010 Mathematics Subject Classification: 57M25, 20C08, 20F36.

Key words and phrases: framization, Temperley-Lieb algebra, Yokonuma-Hecke algebra, Markov trace, link invariants.

This research has been co-financed by the European Union (European Social Fund - ESF) and Greek national funds through the Operational Program “Education and Lifelong Learning” of the National Strategic Reference Framework (NSRF) - Research Funding Program: THALES: Reinforcement of the interdisciplinary and/or inter-institutional research and innovation.

7 Identifying the invariants from $\text{FTL}_{d,n}(u)$ on classical knots and links	337
Acknowledgements	342
References	342

1. Introduction

Since the original construction of the Jones polynomial the Temperley-Lieb algebra has become a cornerstone of a fruitful interaction between Knot theory and Representation theory. The Temperley-Lieb algebra was introduced by Temperley and Lieb [24] and was rediscovered by Jones [10] as a knot algebra [11].

A knot algebra is an algebra that is used in the construction of invariants of classical links using Jones' method [11]. More precisely, a knot algebra A is a triplet (A, π, τ) , where π is an appropriate representation of the braid group in A and τ is a Markov trace function defined on A . The Temperley-Lieb algebra, the Iwahori-Hecke algebra and the BMW algebra are the most known examples of knot algebras.

The 'framization' of a knot algebra is a mechanism designed by the second and fourth authors, that consists in a generalization of a knot algebra via the addition of framing generators. In this way we obtain a new algebra which is related to framed braids and framed knots. More precisely, the framization procedure can roughly be described as the procedure of adding framing generators to the generating set of a knot algebra, of defining interacting relations between the framing generators and the original generators of the algebra and of applying framing on the original defining relations of the algebra. The resulting framed relations should be topologically consistent. The challenge in this procedure is to apply the framization on the relations of polynomial type.

The basic example of framization is the Yokonuma-Hecke algebra, $Y_{d,n}(u)$, which can be regarded as a framization of the Iwahori-Hecke algebra, $H_n(u)$ [15, 19]. The quadratic relation of $Y_{d,n}(u)$ involves intrinsically the framing generators, while for $d = 1$ the algebra $Y_{1,n}(u)$ coincides with $H_n(u)$. Having in mind this example, the second and fourth authors proposed framizations of several knot algebras [18, 20].

The aim of this paper is to propose a framization of the Temperley-Lieb algebra and to derive from this new algebra knot and link invariants via an

appropriate Markov trace. The Temperley-Lieb algebra can be regarded as a quotient of the Iwahori-Hecke algebra. Therefore, it is natural to search for a quotient of the Yokonuma-Hecke algebra over an appropriate two-sided ideal, that can be considered as a framization of the Temperley-Lieb algebra. Although such an ideal is not unique, it will become clear that our choice for the ideal that leads to the framization of the Temperley-Lieb algebra is the most natural one with respect to the construction of related framed and classical link invariants. Indeed, in Section 4 we first discuss two natural quotients of $Y_{d,n}(u)$ that could possibly lead to a framization of the Temperley-Lieb algebra, the Yokonuma-Temperley-Lieb algebra, $YTL_{d,n}(u)$ (introduced and studied in [9]) and the Complex Reflection Temperley-Lieb algebra, $CTL_{d,n}(u)$. These two quotient algebras, however, are not suitable for our purpose, since: The algebra $YTL_{d,n}(u)$ is too restricted and, as a consequence, the invariants for classical links from the algebra $YTL_{d,n}(u)$ just recover the Jones polynomial [9]. On the other hand, as we shall see, the algebra $CTL_{d,n}(u)$ is too large for our topological purposes. We proceed with introducing a third quotient of $Y_{d,n}(u)$, the Framization of the Temperley-Lieb algebra, $FTL_{d,n}(u)$, which lies between $YTL_{d,n}(u)$ and $CTL_{d,n}(u)$ and which will turn out to be the right one. The connection between all three quotients of $Y_{d,n}(u)$ is then analyzed. We note that for $d = 1$ all three quotients coincide with the Temperley-Lieb algebra $TL_n(u)$. We further provide presentations with non-invertible generators for the quotient algebras $FTL_{d,n}(u)$ and $CTL_{d,n}(u)$. Such a presentation for the quotient algebra $YTL_{d,n}(u)$ was given in [9]. We conclude this section with a result by Chlouveraki and Pouchin [6] regarding the dimensions of the quotient algebras $FTL_{d,n}(u)$ and $CTL_{d,n}(u)$.

Returning to our basic example, the Yokonuma-Hecke algebra, the second author has constructed a unique Markov trace function, tr , on the algebra $Y_{d,n}(u)$ with parameters z, x_1, \dots, x_{d-1} [13]. Consequently, invariants for framed, classical and singular oriented links have been obtained [16, 17, 19] by applying the so-called ‘E-condition’ on the parameters x_1, \dots, x_{d-1} so that tr re-scales according to the negative stabilization move between framed braids [19]. These invariants, in particular those for classical links, was necessary to be compared with other known invariants, especially with the 2-variable Jones or Homflypt polynomial. In [4] it was proved that these polynomial invariants do not coincide with the Homflypt polynomial, except in trivial cases. Yet they could be topologically equivalent to the Homflypt polynomial, in the sense that they might distinguish the same pairs of non-isotopic links. Eventually, in a recent development [3], another presentation for the Yokonuma-Hecke algebra is employed with parameter q in a new quadratic relation, where $q^2 = u$ [2]. Using this presentation, the authors of [3] have

been able to establish that the classical link invariants, Θ_d , obtained from the isomorphic algebra $Y_{d,n}(q)$ coincide with the Homflypt polynomial on *knots*, but they are *not topologically equivalent to the Homflypt polynomial on links* (as it was conjectured in [7]).

The next natural question is to examine under what conditions the trace tr on the algebra $Y_{d,n}(u)$ passes through to the quotient algebras $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$ respectively. We recall that, in the classical case, as Jones showed, the Ocneanu trace on the Iwahori-Hecke algebra [11] passes to the quotient $\text{TL}_n(u)$ if and only if the trace parameter ζ takes certain specific values. Accordingly, in Section 5 we provide the necessary and sufficient conditions for the Markov trace tr [13] on the Yokonuma-Hecke algebra to pass through to the quotient algebras $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$. The corresponding conditions for the algebra $\text{YTL}_{d,n}(u)$ are given in [9]. More precisely, we first find the necessary and sufficient conditions on the trace parameters z, x_1, \dots, x_{d-1} , for the algebra $\text{FTL}_{d,3}(u)$ using tools from harmonic analysis on finite groups (Lemma 8) and then we generalize our result using induction on n (Theorem 6). Using the same methods we prove the analogous theorem for $\text{CTL}_{d,n}(u)$ (Theorem 7). For $d = 1$ the specific values we found for z coincide with those found by Jones for $\text{TL}_n(u)$ [11]. Finally, we discuss the connections between the necessary and sufficient conditions for tr to pass to all three quotient algebras $\text{CTL}_{d,n}(u)$, $\text{FTL}_{d,n}(u)$ and $\text{YTL}_{d,n}(u)$.

Using the above conditions on the trace tr and subjecting the trace parameters x_1, \dots, x_{d-1} to the E-condition, we define in Section 6 invariants for framed and classical links through the quotient algebras $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$. We then show that the invariants from the algebras $\text{CTL}_{d,n}(u)$ coincide either with some of the invariants from $Y_{d,n}(u)$ or with some of the invariants from $\text{FTL}_{d,n}(u)$. Since $\text{CTL}_{d,n}(u)$ is larger than $\text{FTL}_{d,n}(u)$ and since we do not obtain from $\text{CTL}_{d,n}(u)$ any extra invariants, for these reasons $\text{FTL}_{d,n}(u)$ is chosen as the framization of the Temperley-Lieb algebra.

Focusing now on the classical link invariants from the algebra $\text{FTL}_{d,n}(u)$, these need to be compared to the Jones polynomial. Following [3], in Section 7 we give a new presentation for the algebra $\text{FTL}_{d,n}$ with parameter q deriving from the new presentation of the Yokonuma-Hecke algebra $Y_{d,n}(q)$. We then adjust our results so far to the isomorphic algebra $\text{FTL}_{d,n}(q)$ and we apply them to the results of [3]. Namely, by specializing $\Theta_d(q, z)$ to the our specific value for z , we obtain 1-variable invariants for classical knots and links, denoted by $\theta_d(q)$. Finally, adapting the results of [3] to the invariants $\theta_d(q)$ we show that they coincide with the Jones polynomial on *knots* but they are *not topologically equivalent to the Jones polynomial on links*.

The outline of the paper is as follows: Section 2 is dedicated to providing necessary definitions and results, including: the Iwahori-Hecke algebra, the Ocneanu trace and the Yokonuma-Hecke algebra. In Section 3 we recall some basic tools from harmonic analysis of finite groups, such as the convolution product, the product by coordinates and the Fourier transform, necessary for exploring the ‘E-system’. In Section 4 we discuss three quotients of the Yokonuma-Hecke algebra as possible candidates for the framization of the Temperley-Lieb algebra. In Section 5 we provide necessary and sufficient conditions for the tr on the Yokonuma-Hecke algebra to pass through to the quotient algebras $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$. In Section 6 we define 1-variable framed and classical link invariants related to the algebras $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$. Finally, in Section 7 we prove that 1-variable classical link invariants derived from the isomorphic algebra $\text{FTL}_{d,n}(q)$ are not topologically equivalent to the Jones polynomial.

The results of this paper lead to further questions worth investigating, as for example, the possibility of obtaining new 3-manifold invariants related to the invariants θ_d , in analogy to the Witten invariants [26].

2. Preliminaries

2.1. Notation

Throughout the paper by the term algebra we mean an associative unital (with unity 1) algebra over $\mathbb{C}(u)$, where u is an indeterminate. Thus we can regard $\mathbb{C}(u)$ as a subalgebra of the center of the algebra. We will also fix two positive integers, d and n .

As usual we denote by $\mathbb{Z}/d\mathbb{Z}$ the group of integers modulo d . We will also denote the underlying set of the group $\mathbb{Z}/d\mathbb{Z}$ by $\{0, 1, \dots, d-1\}$.

We denote S_n the symmetric group on the set $\{1, 2, \dots, n\}$. Let s_i be the elementary transposition $(i, i+1)$ and let $\langle s_i, s_j \rangle$ denote the subgroup generated by s_i and s_j . We also denote by l the length function on S_n with respect to the s_i 's.

Denote by C the infinite cyclic group and by $C_d = \langle t \mid t^d = 1 \rangle$ the cyclic group of order d . Let $t_i := (1, \dots, 1, t, 1, \dots, 1) \in C_d^n$, where t is in the i -th position. We then have:

$$C_d^n = \langle t_1, \dots, t_n \mid t_i t_j = t_j t_i, t_i^d = 1 \rangle.$$

Define $C_{d,n} := C_d^n \rtimes S_n$, where the action is defined by permutation on the indices of the t_i 's, namely: $s_i t_j = t_{s_i(j)} s_i$. Notice that $C_{d,n}$ is isomorphic

to the *complex reflection group* $G(d, 1, n)$. We also introduce the following notation $C_{\infty, n} := C^n \rtimes S_n$.

Denote by B_n the braid group of type A , that is, the group generated by the elementary braidings $\sigma_1, \dots, \sigma_{n-1}$, subject to the following relations: $\sigma_i \sigma_j \sigma_i = \sigma_j \sigma_i \sigma_j$, for $|i - j| = 1$ and $\sigma_i \sigma_j = \sigma_j \sigma_i$, for $|i - j| > 1$. We will also use the *d-modular framed braid group* $\mathcal{F}_{d, n} := C_d^n \rtimes B_n$, where the action of B_n on C_d^n is defined by the induced permutation on the indices of the t_i 's. We will also refer to the *framed braid group* $\mathcal{F}_n := C^n \rtimes B_n$. Of course, we have isomorphisms: $\mathcal{F}_n \cong \mathbb{Z}^n \rtimes B_n$ and $\mathcal{F}_{d, n} \cong (\mathbb{Z}/d\mathbb{Z})^n \rtimes B_n$. Finally, note that the natural projections $C \rightarrow C_d$ and $B_n \rightarrow S_n$ induce the following commutative diagram:

$$\begin{array}{ccccccc}
 \mathcal{F}_n & \longrightarrow & \mathcal{F}_{d, n} & \longrightarrow & B_n & \longrightarrow & 1 \\
 \downarrow & & \downarrow & & \downarrow & & \\
 C_{\infty, n} & \longrightarrow & C_{d, n} & \longrightarrow & S_n & \longrightarrow & 1 \\
 \downarrow & & \downarrow & & \downarrow & & \\
 1 & & 1 & & 1 & &
 \end{array}$$

From the above diagram one can define the *length function* l' on $C_{d, n}$ as the lift of the ordinary length function l of S_n , that is:

$$(2.1) \quad l'(t^{a_1} s_{i_1} \cdots s_{i_k}) := l(s_{i_1} \cdots s_{i_k}),$$

where $t^a := t_1^{a_1} \cdots t_n^{a_n} \in C_d^n$.

Remark 1. We would like to point out that $C_{d, n}$ and $\mathcal{F}_{d, n}$ appear in the theory of “fields with one element”. This is a theory dreamt by J. Tits in his study of algebraic groups. According to the seminal article of Kapranov and Smirnov [21], $\mathrm{GL}_n(\mathbb{F}_1) = S_n$, $\mathrm{GL}_n(\mathbb{F}_1[t]) = B_n$, $\mathrm{GL}_n(\mathbb{F}_{1^n}) = C_{d, n}$ and $\mathrm{GL}_n(\mathbb{F}_{1^n}[t]) = \mathcal{F}_{d, n}$, where $\mathrm{GL}_n(\mathbb{F}_{1^n})$ (resp. $\mathrm{GL}_n(\mathbb{F}_{1^n}[t])$) is in “some sense” the limit case $q \rightarrow 1$ of $\mathrm{GL}_n(\mathbb{F}_q)$ (resp. $\mathrm{GL}_n(\mathbb{F}_q[t])$).

2.2. Background material

We denote by $H_n(u)$ the *Iwahori-Hecke algebra* associated to S_n , that is, the $\mathbb{C}(u)$ -algebra with linear basis $\{h_w \mid w \in S_n\}$ and the following rules of

multiplication:

$$(2.2) \quad h_{s_i}h_w = \begin{cases} h_{s_iw} & \text{for } l(s_iw) > l(w) \\ uh_{s_iw} + (u - 1)h_w & \text{for } l(s_iw) < l(w). \end{cases}$$

Set $h_i := h_{s_i}$. Then $H_n(u)$ is presented by h_1, \dots, h_{n-1} subject to the following relations:

$$(2.3) \quad h_i h_j = h_j h_i \quad \text{for all } |i - j| > 1$$

$$(2.4) \quad h_i h_j h_i = h_j h_i h_j \quad \text{for all } |i - j| = 1$$

$$(2.5) \quad h_i^2 = u + (u - 1)h_i.$$

Definition 1. For $n \geq 3$, the *Temperley-Lieb algebra* $TL_n(u)$ can be defined as the quotient of the algebra $H_n(u)$ over the two-sided ideal generated by the *Steinberg elements* $h_{i,j}$:

$$(2.6) \quad h_{i,j} := \sum_{w \in \langle s_i, s_j \rangle} h_w, \quad \text{for all } |i - j| = 1.$$

Consequently, the algebra $TL_n(u)$ can be thus presented by h_1, \dots, h_{n-1} subject to relations (2.3)-(2.5) and the following relations:

$$1 + h_i + h_j + h_i h_j + h_j h_i + h_i h_j h_i = 0 \quad \text{for all } |i - j| = 1.$$

The defining ideal of the algebra $TL_n(u)$ is principal and it is generated by the element $h_{1,2}$. Furthermore, using the transformation:

$$(2.7) \quad f_i := \frac{1}{u + 1}(h_i + 1),$$

the algebra $TL_n(u)$ can be presented by the non-invertible generators f_1, \dots, f_{n-1} subject to the following relations:

$$\begin{aligned} f_i^2 &= f_i \\ f_i f_j f_i &= \delta f_i, \quad \text{for all } |i - j| = 1 \\ f_i f_j &= f_j f_i, \quad \text{for all } |i - j| > 1, \end{aligned}$$

where $\delta^{-1} = 2 + u + u^{-1}$ [11].

In [8, 11] Ocneanu constructed a unique Markov trace on $H_n(u)$. More precisely, we have the following theorem.

Theorem 1 (Ocneanu). *Let ζ be an indeterminate. There exists a linear trace τ on $\cup_{n=1}^{\infty} H_n(u)$ uniquely defined by the inductive rules:*

- 1) $\tau(ab) = \tau(ba), \quad a, b \in H_n(u)$
- 2) $\tau(1) = 1$
- 3) $\tau(ah_n) = \zeta \tau(a), \quad a \in H_n(u) \quad (\text{Markov property}).$

The Ocneanu trace τ passes through to $TL_n(u)$ for specific values of ζ . Indeed, as it turned out [11], to factorize τ to the Temperley-Lieb algebra, we only need the fact that τ annihilates the expression of Eq. 2.6. So, in [11] it is proved that τ passes to the Temperley-Lieb algebra if and only if:

$$(2.8) \quad \zeta = -\frac{1}{u+1} \quad \text{or} \quad \zeta = -1.$$

2.3. The Yokonuma-Hecke algebra

The Yokonuma-Hecke algebra of type A , denoted by $Y_{d,n}(u)$ [27], can be defined by generators and relations [13] and can be regarded as a quotient of $\mathbb{C}(u)\mathcal{F}_{d,n}$ over the two-sided ideal that is generated by the elements:

$$\sigma_i^2 - (u-1)e_i - (u-1)e_i\sigma_i - 1,$$

where e_i is the idempotent defined by:

$$(2.9) \quad e_i := \frac{1}{d} \sum_{s=0}^{d-1} t_i^s t_{i+1}^{d-s}, \quad i = 1, \dots, n-1.$$

Equivalently, one can define $Y_{d,n}(u)$ as follows:

Definition 2. The *Yokonuma-Hecke algebra* $Y_{d,n}(u)$ is the algebra presented by generators $g_1, \dots, g_{n-1}, t_1, \dots, t_n$ subject to the following relations:

$$(2.10) \quad g_i g_j = g_j g_i \quad \text{for all } |i-j| > 1$$

$$(2.11) \quad g_{i+1} g_i g_{i+1} = g_i g_{i+1} g_i$$

$$(2.12) \quad t_i t_j = t_j t_i \quad \text{for all } i, j$$

$$(2.13) \quad t_i^d = 1 \quad \text{for all } i$$

$$(2.14) \quad g_i t_i = t_{i+1} g_i$$

$$(2.15) \quad g_i t_{i+1} = t_i g_i$$

$$(2.16) \quad g_i t_j = t_j g_i \quad \text{for } j \neq i, i + 1$$

$$(2.17) \quad g_i^2 = 1 + (u - 1)e_i + (u - 1)e_i g_i.$$

Note that for $d = 1$ the quadratic relation (2.17) becomes:

$$g_i^2 = (u - 1)g_i + u.$$

So, the Yokonuma-Hecke $Y_{1,n}(u)$ coincides with the Iwahori-Hecke algebra.

The algebra $Y_{d,n}(u)$ can also be regarded as a u -deformation of the group algebra $\mathbb{C}C_{d,n}$. Indeed, if $w \in S_n$ is a reduced word in S_n with $w = s_{i_1} \cdots s_{i_k}$ then the expression $g_w = g_{s_{i_1}} \cdots g_{s_{i_r}} \in Y_{d,n}(u)$ is well-defined since the generators $g_i := g_{s_i}$ satisfy the same braiding relations as the generators of S_n [22]. We have the following multiplication rule in $Y_{d,n}(u)$ (see [12, Proposition 2.14]):

$$(2.18) \quad g_{s_i} g_w = \begin{cases} g_{s_i w} & \text{for } l'(s_i w) > l'(w) \\ g_{s_i w} + (u - 1)e_i g_{s_i w} + (u - 1)e_i g_w & \text{for } l'(s_i w) < l'(w). \end{cases}$$

Note also that the generators g_{t_i} correspond to t_i and so, using Eq. 2.1, we have that: $g_{t_i w} = g_{t_i} g_w = t_i g_w$.

The definition of the idempotents e_i can be generalized in the following way. For any indices i, j we define the following elements in $Y_{d,n}(u)$:

$$(2.19) \quad e_{i,j} := \frac{1}{d} \sum_{s=0}^{d-1} t_i^s t_j^{d-s}.$$

We also define, for any $0 \leq m \leq d - 1$, the shift of e_i by m :

$$(2.20) \quad e_i^{(m)} := \frac{1}{d} \sum_{s=0}^{d-1} t_i^{m+s} t_{i+1}^{d-s}.$$

Notice that $e_i = e_{i,i+1} = e_i^{(0)}$. Notice also that $e_i^{(m)} = t_i^m e_i = t_{i+1}^m e_i$. Then one deduces easily that:

$$(2.21) \quad \begin{aligned} e_i^{(m)} e_{i+1} &= e_i e_{i+1}^{(m)} \\ t_1^a t_2^b t_3^c e_1 e_2 &= e_1^{(a+b+c)} e_2, \end{aligned}$$

for all $0 \leq m, a, b, c \leq d - 1$.

The following lemma collects some of the relations among the e_i 's, the t_j 's and the g_i 's. These relations will be used in the paper.

Lemma 1 ([9, Lemma 1]). *For the idempotents e_i and for $1 \leq i, j \leq n - 1$ the following relations hold:*

$$\begin{aligned} t_j e_i &= e_i t_j \\ e_{i+1} g_i &= g_i e_{i+2} \\ e_i g_j &= g_j e_i, \quad \text{for } j \neq i - 1, i + 1 \\ e_j g_i g_j &= g_i g_j e_i \quad \text{for } |i - j| = 1 \\ e_i e_{i+1} &= e_i e_{i+2} \\ e_i e_{i+1} &= e_{i,i+2} e_{i+1}. \end{aligned}$$

A word in the defining generators of the algebra will be called a *monomial*. Notice that using relations (2.14) and (2.15) one can write any monomial \mathbf{m} in $Y_{d,n}(u)$ in the following form:

$$\mathbf{m} = t_1^{a_1} \cdots t_n^{a_n} \mathbf{m}',$$

where $\mathbf{m}' = g_{i_1} \cdots g_{i_n}$. We then say that every monomial in $Y_{d,n}(u)$ has the *splitting property*, which is in fact inherited from the framed braid group \mathcal{F}_n . That is, one can separate the *framing part* of \mathbf{m} (which is the subword in the framing generators t_j) from the *braiding part* (which is the subword in the braiding generators g_i).

2.4. A Markov trace on $Y_{d,n}(u)$

Using the multiplication formulas (2.18), the second author proved in [13] that $Y_{d,n}(u)$ has the following standard linear basis:

$$(2.22) \quad \{t_1^{a_1} \cdots t_n^{a_n} g_w \mid a_i \in \mathbb{Z}/d\mathbb{Z}, w \in S_n\}.$$

This above linear basis led naturally to the following inductive basis for the Yokonuma-Hecke algebra, which we will use in the proof of the main theorem (Theorem 6).

Proposition 1 ([13, Proposition 8]). *Every element in $Y_{d,n+1}(u)$ is a unique linear combination of words, each of one of the following types:*

$$\mathbf{m}_n g_n g_{n-1} \cdots g_i t_i^k \quad \text{or} \quad \mathbf{m}_n t_{n+1}^k,$$

where $0 \leq k \leq d - 1$ and \mathbf{m}_n is a word in the inductive basis of $Y_{d,n}(u)$.

Employing the above inductive basis, the second author proved that $Y_{d,n}(u)$ supports a unique Markov trace. We have the following theorem:

Theorem 2 ([13, Theorem 12]). *For indeterminates z, x_1, \dots, x_{d-1} there exists a unique linear Markov trace tr :*

$$\text{tr} : \cup_{n=1}^{\infty} Y_{d,n}(u) \longrightarrow \mathbb{C}(u)[z, x_1, \dots, x_{d-1}],$$

defined inductively on n by the following rules:

$$\begin{aligned} \text{tr}(ab) &= \text{tr}(ba) \\ \text{tr}(1) &= 1 \\ \text{tr}(ag_n) &= z \text{tr}(a) && \text{(Markov property)} \\ \text{tr}(at_{n+1}^s) &= x_s \text{tr}(a) && (s = 1, \dots, d - 1), \end{aligned}$$

where $a, b \in Y_{d,n}(u)$.

Using the trace rules of Theorem 2 and setting $x_0 := 1$, we deduce that $\text{tr}(e_i)$ takes the same value for all i , and this value is denoted by E :

$$E := \text{tr}(e_i) = \frac{1}{d} \sum_{s=0}^{d-1} x_s x_{d-s}.$$

Moreover, we also define *the shift by m of E* , where $0 \leq m \leq d - 1$, by:

$$E^{(m)} := \text{tr}(e_i^{(m)}) = \frac{1}{d} \sum_{s=0}^{d-1} x_{m+s} x_{d-s}.$$

Notice that $E = E^{(0)}$.

3. Fourier transform and the E-system

An important tool in the proof of the main theorem are some classical identities of harmonic analysis on the group of integers modulo d . More precisely,

we will use identities linking the convolution product and the product by coordinates through the Fourier transform. These tools were also used in solving the so-called E-system, see [19, Appendix]. Thus, in this section we shall give some notations and recall some well-known and useful facts of the Fourier transform along with some facts for the E-system.

3.1. Computations in $\mathbb{C}C_d$

Recall that C_d is the cyclic group of order d , generated by t . The *product by coordinates* in $\mathbb{C}C_d$ is defined by the formula:

$$\left(\sum_{r=0}^{d-1} a_r t^r \right) \cdot \left(\sum_{s=0}^{d-1} b_s t^s \right) = \sum_{i=0}^{d-1} a_i b_i t^i$$

and the *convolution product* is defined by the formula:

$$(3.1) \quad \left(\sum_{r=0}^{d-1} a_r t^r \right) * \left(\sum_{s=0}^{d-1} b_s t^s \right) = \sum_{r=0}^{d-1} \left(\sum_{s=0}^{d-1} a_s b_{r-s} \right) t^r.$$

In order to define the Fourier transform on C_d we need to introduce the following elements:

$$\mathbf{i}_a := \sum_{s=0}^{d-1} \chi_a(t^s) t^s \quad (a \in \mathbb{Z}/d\mathbb{Z}),$$

where the χ_k 's denote the characters of the group C_d , namely:

$$(3.2) \quad \chi_k(t^m) = \cos \frac{2\pi km}{d} + i \sin \frac{2\pi km}{d} \quad (k, m \in \mathbb{Z}/d\mathbb{Z}).$$

One can verify that:

$$\mathbf{i}_a * \mathbf{i}_b = \begin{cases} d \mathbf{i}_a & \text{if } a = b \\ 0 & \text{if } a \neq b. \end{cases}$$

On the other hand, we shall denote by δ_a the element t^a of the canonical linear basis of $\mathbb{C}C_d$. It is clear that:

$$\delta_a \cdot \delta_b = \begin{cases} \delta_a & \text{if } a = b \\ 0 & \text{if } a \neq b. \end{cases}$$

The *Fourier transform* is the linear automorphism on $\mathbb{C}C_d$, defined by:

$$(3.3) \quad y := \sum_{r=0}^{d-1} a_r t^r \mapsto \widehat{y} := \sum_{s=0}^{d-1} (y * \mathbf{i}_s)(0) t^s,$$

where $(y * \mathbf{i}_s)(0)$ denote the coefficient of δ_0 in the convolution $y * \mathbf{i}_s$.

The next proposition gathers the most important properties of the Fourier transform used in the paper.

Proposition 2 ([25, Chapter 2]). *For any y and y' in $\mathbb{C}C_d$, we have:*

- 1) $\widehat{y * y'} = \widehat{y} \cdot \widehat{y'}$
- 2) $\widehat{y \cdot y'} = d^{-1} \widehat{y} * \widehat{y'}$
- 3) $\widehat{\delta}_a = \mathbf{i}_{-a}$
- 4) $\widehat{\mathbf{i}}_a = d \delta_a$
- 5) *If $y = \sum_{r=0}^{d-1} a_r t^r$, then $\widehat{y} = d \sum_{r=0}^{d-1} a_{-r} t^r$.*

Finally, we note that the elements in the group algebra $\mathbb{C}C_d$ can also be identified to the set of functions $f : C_d \rightarrow \mathbb{C}$, where the identification is as follows:

$$(3.4) \quad (f : C_d \rightarrow \mathbb{C}) \longleftrightarrow \sum_{k=0}^{d-1} f(t^k) t^k \in \mathbb{C}C_d.$$

Some times we shall use this identification, since it makes some computations easier.

3.2. The E-system and its solutions

The E-system is a non-linear system of equations that was introduced in order to find the necessary and sufficient conditions that need to be applied on the parameters x_i of tr so that the definition of link invariants from the Yokonuma-Hecke algebra would be possible [19].

Definition 3 ([19, Definition 11]). We say that the $(d - 1)$ -tuple of complex numbers (x_1, \dots, x_{d-1}) satisfies the *E-condition* if x_1, \dots, x_{d-1} satisfy the following system of non-linear equations in \mathbb{C} , the *E-system*:

$$(3.5) \quad E^{(m)} = x_m E \quad (1 \leq m \leq d - 1).$$

In [19, Appendix] the full set of solutions of the E-system is given by Gérardin using some tools of harmonic analysis on finite groups. More precisely, he interpreted the solution (x_1, \dots, x_d) of the E-system, as a certain complex function $x_D : C_d \rightarrow \mathbb{C}$. The solution is parametrized by a non-empty subset D of C_d^* , where C_d^* denotes the dual group of C_d , i.e. the space of characters of C_d . Since $C_d \cong C_d^* \cong \mathbb{Z}/d\mathbb{Z}$, by small abuse of notation, we will consider D as a subset of $\mathbb{Z}/d\mathbb{Z}$. Recall that the characters χ_k of C_d are given by $t^a \mapsto \chi_k(t^a)$, where k runs over $\mathbb{Z}/d\mathbb{Z}$, see Eq. 3.2.

The dependence of x_D on D is given by the following equation of functions:

$$(3.6) \quad x_D = \frac{1}{|D|} \sum_{k \in D} \chi_k.$$

Notice that the function x_D can be also seen as an element in $\mathbb{C}C_d$, namely:

$$(3.7) \quad x_D = \sum_{j=0}^{d-1} x_j t^j,$$

where $x_j = x_D(t^j) = \frac{1}{|D|} \sum_{k \in D} \chi_k(t^j)$.

A simple computation shows that the convolution products, where x is an element in the group algebra $\mathbb{C}C_d$, are given by:

$$(3.8) \quad \begin{aligned} x * x &= d \sum_{k=0}^{d-1} \text{tr}(e_i^{(k)}) t^k = d \sum_{k=0}^{d-1} E^{(k)} t^k, \\ x * x * x &= d^2 \sum_{k=0}^{d-1} \text{tr}(e_1^{(k)} e_2) t^k, \end{aligned}$$

see also [9, Lemma 2]

Remark 2. It is worth noting that the formula for the solutions of the E-system can be interpreted as a generalization of the Ramanujan sum. Indeed, by taking the subset P of C_d consisting of the numbers coprime to d , then the solution parametrized by P is, up to the factor $|P|$, the Ramanujan sum $c_d(k)$ (see [23]).

We finish this section with a theorem which yields the main connection among the solutions of the E-system and the trace tr .

Theorem 3 ([19, Theorem 7]). *If the trace parameters (x_1, \dots, x_{d-1}) satisfy the E-condition, then*

$$\text{tr}(\alpha e_n) = \text{tr}(\alpha)\text{tr}(e_n) \quad (a \in Y_{d,n}(u)).$$

4. A Framization of the Temperley-Lieb algebra

In this section we explore three quotients of the Yokonuma-Hecke algebra, the algebras $YTL_{d,n}(u)$, $FLL_{d,n}(u)$ and $CTL_{d,n}(u)$, as potential candidates for the framization of the Temperley-Lieb algebra and we select one of them, namely $FLL_{d,n}(u)$, as the most appropriate in view of our topological aims.

4.1. The three potential candidates

As discussed during the Introduction, the Yokonuma-Hecke algebra can be interpreted as the framization of the Iwahori-Hecke algebra, which is a knot algebra. Thus a natural question arises, the definition of a framization for the knot algebra Temperley-Lieb. Considering the fact that the Temperley-Lieb algebra can be defined as a quotient of the Iwahori-Hecke algebra, it is natural to try and define a framization of the Temperley-Lieb algebra as a quotient of the Yokonuma-Hecke algebra. Recall now that the defining ideal of the Temperley-Lieb algebra (Definition 1) is generated by the Steinberg elements which are related to the subgroups $\langle s_i, s_{i+1} \rangle$ of S_n , for all i . These subgroups can be also regarded as subgroups of $C_{d,n}$. Therefore, using the multiplication rule of Eq. 2.18 we are able to define the analogous Steinberg elements $g_{i,i+1}$ in $Y_{d,n}(u)$,

$$(4.1) \quad g_{i,i+1} := \sum_{w \in \langle s_i, s_{i+1} \rangle} g_w \quad \text{for all } i.$$

In [9, Definition 2] we defined a potential candidate for the framization of the Temperley-Lieb algebra, the *Yokonuma-Temperley-Lieb algebra*, denoted by $YTL_n(u)$ which is defined as the quotient of $Y_{d,n}(u)$ over the two-sided ideal generated by the $g_{i,i+1}$'s for all i . It is not difficult to show that this ideal is in fact principal and it is generated by the element $g_{1,2}$. Moreover, the necessary and sufficient conditions for the trace tr to pass through to $YTL_n(u)$ were studied [9, Theorem 6]. Unfortunately, these conditions turn out to be too strong. Namely, the trace parameters x_i must be d^{th} roots of unity, giving rise to obvious, special solutions of the E-system, which imply topologically loss of the framing information. Moreover, if we restrict to the case of classical

links, by representing the Artin braid group B_n in $Y_{d,n}(u)$, considering the t_i 's as formal generators, and then taking the quotient over the ideal that is generated by the $g_{i,i+1}$'s [9, Section 5], and using the results of [4], the derived classical link invariants for the algebras $YTL_{d,n}(u)$ coincide with the classical Jones polynomial. For the above reasons, $YTL_{d,n}(u)$ is discarded as framization of $TL_n(u)$. Finally, we note that the representation theory of this algebra has been studied extensively in [5].

Given the fact that $Y_{d,n}(u)$ can be considered as a u -deformation of $CC_{d,n}$ (recall the discussion in Section 2.3), it is natural to consider subgroups of $C_{d,n}$ that involve in their generating set the framing generators of the i -th and j -th strands along with $\langle s_i, s_j \rangle$. As a first attempt, we consider the following subgroups of $C_{d,n}$:

$$C_{d,n}^i := \langle t_i, t_{i+1}, t_{i+2} \rangle \rtimes \langle s_i, s_{i+1} \rangle \quad \text{for all } i.$$

Notice that these subgroups are isomorphic to the group $C_{d,3}$, in analogy to the classical case of $TL_n(u)$. We define now the elements $c_{i,i+1}$ in $Y_{d,n}(u)$ as follows:

$$(4.2) \quad c_{i,i+1} = \sum_{c \in C_{d,n}^i} g_c.$$

We then have the following definition:

Definition 4. For $n \geq 3$, we define the algebra $CTL_{d,n}(u)$ as the quotient of the algebra $Y_{d,n}(u)$ by the two-sided ideal generated by the $c_{i,i+1}$'s, for all i . We shall call $CTL_{d,n}(u)$ the *Complex Reflection Temperley-Lieb algebra*.

Remark 3. The denomination Complex Reflection Temperley-Lieb algebra has to do with the fact that the underlying group of $CTL_{d,n}(u)$ is isomorphic to the complex reflection group $G(d, 1, 3)$.

As it will be shown in Theorem 7, the necessary and sufficient conditions such that tr passes to $CTL_{d,n}(u)$ are, contrary to the case of $YTL_{d,n}(u)$, too relaxed, especially on the trace parameters x_i . So, in order to define link invariants from the algebras $CTL_{d,n}(u)$, the E-condition must be imposed on the x_i 's.

This indicates that the desired framization of the Temperley-Lieb algebra for our topological purposes could be an intermediate algebra between these two. We achieve this, by using for the defining ideal an intermediate subgroup

of C_d^n that lies between $\langle s_i, s_{i+1} \rangle$ and $C_{d,n}^i$. Indeed, we consider the following subgroups of $C_{d,n}$,

$$H_{d,n}^i := \langle t_i t_{i+1}^{-1}, t_{i+1} t_{i+2}^{-1} \rangle \rtimes \langle s_i, s_{i+1} \rangle \quad \text{for all } i.$$

We now introduce the following elements:

$$r_{i,i+1} := \sum_{x \in H_{d,n}^i} g_x \quad \text{for all } i.$$

Definition 5. For $n \geq 3$, the *Framization of the Temperley-Lieb algebra*, denoted $\text{FTL}_{d,n}(u)$, is defined as the quotient $Y_{d,n}(u)$ over the two-sided ideal generated by the elements $r_{i,i+1}$, for all i .

Remark 4. Notice that when $d = 1$, the Yokonuma-Hecke algebra coincides with the Iwahori-Hecke algebra, hence it follows that $\text{YTTL}_{1,n}(u)$ also coincides with $\text{TL}_n(u)$. Moreover, in this case the subgroups $H_{d,n}^i$ and $C_{d,n}^i$ also collapse to $\langle s_i, s_{i+1} \rangle$, which is isomorphic to S_3 . Hence, $\text{FTL}_{1,n}(u)$ and $\text{CTL}_{1,n}(u)$ coincide with $\text{TL}_n(u)$ too.

4.2. Relating the three quotient algebras

We shall now show how the algebras defined above are related. Notice that the defining ideal for each quotient algebra mentioned above is generated by sums of elements g_x , where x belongs to the underlying group of each ideal. More precisely, the underlying group of the defining ideal of $\text{YTTL}_{d,n}(u)$ is S_3 of $\text{FTL}_{d,n}(u)$ is $H_{d,n}^i$ and of $\text{CTL}_{d,n}(u)$ is $C_{d,n}^i$. We have the following inclusion of groups : $S_3 \leq H_{d,n}^i \leq C_{d,n}^i$. We will show that this implies the following inclusions of ideals:

$$(4.3) \quad \langle c_{i,i+1} \rangle \triangleleft \langle r_{i,i+1} \rangle \triangleleft \langle g_{i,i+1} \rangle.$$

The second inclusion of the ideals, $\langle r_{i,i+1} \rangle \triangleleft \langle g_{i,i+1} \rangle$, is clear. Indeed, every x in $H_{d,n}^i$ can be written in the form:

$$x = t_i^a t_{i+1}^{-a} t_{i+1}^b t_{i+2}^{-b} w = t_i^a t_{i+1}^{b-a} t_{i+2}^{-b} w, \quad \text{where } w \in S_3.$$

Therefore, from the multiplication rule of Eq. 2.18, we have that $g_x = t_i^a t_{i+1}^{b-a} t_{i+2}^{-b} g_w$. Thus we can rewrite the elements $r_{i,i+1}$ in the following form:

$$r_{i,i+1} = \sum_{\substack{a,b=0 \\ w \in S_3}}^{d-1} t_i^a t_{i+1}^{b-a} t_{i+2}^{-b} g_w = \left(\sum_{a,b=0}^{d-1} t_i^a t_{i+1}^{b-a} t_{i+2}^{-b} \right) \left(\sum_{w \in S_3} g_w \right),$$

hence

$$(4.4) \quad r_{i,i+1} = d^2 e_i e_{i+1} g_{i,i+1}.$$

We shall proceed now with the proof of the first inclusion of ideals. We observe that:

$$(4.5) \quad C_{d,n}^i = H_{d,n}^i \rtimes C_d.$$

Indeed, let $x = t_i^a t_{i+1}^b t_{i+2}^c w$ an element in $C_{d,n}^i$, where $w \in S_3$, and let ϕ be the following homomorphism:

$$\begin{aligned} \phi : C_{d,n}^i &\rightarrow \langle t_i \rangle \cong C_d \\ x &\mapsto t_i^{a+b+c}. \end{aligned}$$

Observe that $\ker \phi = H_{d,n}^i$, so $\phi|_{H_{d,n}^i} = \text{id}_{C_d}$, which implies Eq. 4.5. Therefore, for the element $x \in C_{d,n}^i$ we have a unique decomposition $x = t_i^k y$, where $0 \leq k \leq d - 1$ and $y \in H_{d,n}^i$. This decomposition of the elements of $C_{d,n}^i$ together with the multiplication rule in Eq. 2.18, implies $g_x = t_i^k g_y$. This allows us to write the elements $c_{i,i+1}$ of Eq. 4.2 in the following equivalent form:

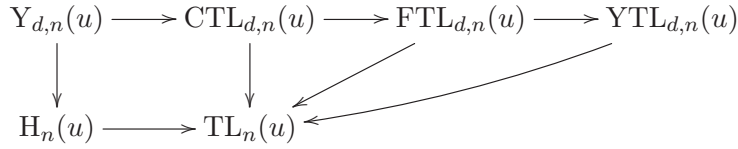
$$c_{i,i+1} = \sum_{\substack{0 \leq k \leq d-1 \\ y \in H_{d,n}^i}} t_i^k g_y,$$

hence:

$$(4.6) \quad c_{i,i+1} = \left(\sum_{k=0}^{d-1} t_i^k \right) r_{i,i+1}.$$

Equation 4.6 implies that $\text{CTL}_{d,n}(u)$ projects onto $\text{FTL}_{d,n}(u)$ while Equation 4.4 implies that $\text{FTL}_{d,n}(u)$ projects onto $\text{YTL}_{d,n}(u)$. We have thus proved the following:

Proposition 3. *The inclusions of ideals of Eq. 4.3 yield the following natural commutative diagram of epimorphisms:*



where the non-horizontal arrows are defined by mapping the framing generators to 1.

4.3. Principality of the ideals

It is known that the defining ideal of the Temperley-Lieb algebra is principal [11]. We are going now to prove that the defining ideals of $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$ respectively are principal ideals too. The method used in the proof is standard [11] but for self-containedness of the paper we will sketch the basic ideas. We start with a technical lemma.

Lemma 2. *The following hold in $Y_{d,n}(u)$ for all $i = 1, \dots, n-2$, $j = 1, \dots, n$ and $0 \leq a, b, c \leq d-1$:*

- (1) $t_j = (g_1 \cdots g_{n-1})^{j-1} t_1 (g_1 \cdots g_{n-1})^{-(j-1)}$
- (2) $g_i = (g_1 \cdots g_{n-1})^{i-1} g_1 (g_1 \cdots g_{n-1})^{-(i-1)}$
- (3) $t_i^a t_{i+1}^b t_{i+2}^c = (g_1 \cdots g_{n-1})^{i-1} t_1^a t_2^b t_3^c (g_1 \cdots g_{n-1})^{-(i-1)}$
- (4) $t_i^a t_{i+1}^b t_{i+2}^c g_i = (g_1 \cdots g_{n-1})^{i-1} t_1^a t_2^b t_3^c g_1 (g_1 \cdots g_{n-1})^{-(i-1)}$
- (5) $t_i^a t_{i+1}^b t_{i+2}^c g_{i+1} = (g_1 \cdots g_{n-1})^{i-1} t_1^a t_2^b t_3^c g_2 (g_1 \cdots g_{n-1})^{-(i-1)}$
- (6) $t_i^a t_{i+1}^b t_{i+2}^c g_i g_{i+1} = (g_1 \cdots g_{n-1})^{i-1} t_1^a t_2^b t_3^c g_1 g_2 (g_1 \cdots g_{n-1})^{-(i-1)}$
- (7) $t_i^a t_{i+1}^b t_{i+2}^c g_{i+1} g_i = (g_1 \cdots g_{n-1})^{i-1} t_1^a t_2^b t_3^c g_2 g_1 (g_1 \cdots g_{n-1})^{-(i-1)}$
- (8) $t_i^a t_{i+1}^b t_{i+2}^c g_i g_{i+1} g_i = (g_1 \cdots g_{n-1})^{i-1} t_1^a t_2^b t_3^c g_1 g_2 g_1 (g_1 \cdots g_{n-1})^{-(i-1)}$.

Proof. Statement (1) is proved by application of Eqs. 2.14–2.16 and induction on j . The proof of the second statement is standard in the literature; it follows from the braid relations (2.10) and (2.11) and induction on i . The other statements of the Lemma are proved by repeated applications of statements (1) and (2). □

Lemma 3. *The defining ideal of $\text{FTL}_{d,n}(u)$ is generated by any single element $r_{i,i+1}$.*

Proof. It is enough to prove that $r_{i,i+1} = (g_1 \cdots g_{n-1})^{(i-1)} r_{1,2} (g_1 \cdots g_{n-1})^{-(i-1)}$. Indeed, expanding $r_{1,2}$ in the right-hand side of the equality, we have:

$$\begin{aligned} & (g_1 \cdots g_{n-1})^{i-1} r_{1,2} (g_1 \cdots g_{n-1})^{-(i-1)} \\ &= \sum_{\substack{a,b=0 \\ w \in S_3}}^{d-1} (g_1 \cdots g_{n-1})^{i-1} t_1^a t_2^{b-a} t_3^{-b} g_w (g_1 \cdots g_{n-1})^{-(i-1)} \\ &= \sum_{a,b=0}^{d-1} (g_1 \cdots g_{n-1})^{i-1} t_1^a t_2^{b-a} t_3^{-b} \left(\sum_{w \in S_3} g_w \right) (g_1 \cdots g_{n-1})^{-(i-1)} = r_{i,i+1}, \end{aligned}$$

Therefore the proof is concluded. □

The following is an immediate corollary of Lemma 3.

Corollary 1. *FTL_{d,n}(u) is the algebra generated by $t_1, \dots, t_n, g_1, \dots, g_{n-1}$ which are subject to the defining relations of $Y_{d,n}(u)$ and the relation:*

$$(4.7) \quad r_{1,2} = 0.$$

Further, an analogous result (with analogous proofs) holds for the quotient algebra CTL_{d,n}(u). So we have the following:

Corollary 2. *The defining ideal of CTL_{d,n}(u) is generated by any single element $c_{i,i+1}$. Hence CTL_{d,n}(u) can be presented by $t_1, \dots, t_n, g_1, \dots, g_{n-1}$ together with the defining relations of $Y_{d,n}(u)$ and the relation:*

$$(4.8) \quad c_{1,2} = 0.$$

4.4. Presentations with non-invertible generators

By using the analogous transformation to Eq. 2.7, we obtain presentations for FTL_{d,n}(u) and CTL_{d,n}(u) through non-invertible generators. More precisely, set

$$(4.9) \quad \ell_i := \frac{1}{u+1} (g_i + 1).$$

Proposition 4. *The algebra $\text{FTL}_{d,n}(u)$ can be presented with generators $\ell_1, \dots, \ell_{n-1}, t_1, \dots, t_n$, subject to the following relations:*

$$(4.10) \quad \ell_i \ell_j = \ell_j \ell_i, \quad \text{for } |i - j| > 1$$

$$(4.11) \quad \ell_i \ell_{i+1} \ell_i - \frac{(u-1)e_i + 1}{(u+1)^2} \ell_i = \ell_{i+1} \ell_i \ell_{i+1} - \frac{(u-1)e_{i+1} + 1}{(u+1)^2} \ell_{i+1}$$

$$(4.12) \quad t_i^d = 1, \quad t_i t_j = t_j t_i$$

$$(4.13) \quad \ell_i t_i = t_{i+1} \ell_i + \frac{1}{u+1} (t_i - t_{i+1})$$

$$(4.14) \quad \ell_i t_{i+1} = t_i \ell_i + \frac{1}{u+1} (t_{i+1} - t_i)$$

$$(4.15) \quad \ell_i t_j = t_j \ell_i, \quad \text{for } |i - j| > 1$$

$$(4.16) \quad \ell_i^2 = \frac{(u-1)e_i + 2}{u+1} \ell_i,$$

$$(4.17) \quad e_i e_{i+1} \ell_i \ell_{i+1} \ell_i = \frac{u}{(u+1)^2} e_i e_{i+1} \ell_i.$$

Proof. It is a straightforward computation to see that relations (2.10)–(2.17) are transformed via Eq. 4.9 into the relations (4.11)–(4.17). We will prove here some indicative cases. The rest are proved in an analogous way. First we will prove the quadratic relation (4.16). From Eq. 4.9 we have that:

$$g_i^2 = ((u+1)\ell_i - 1)^2,$$

using Eq. 2.17, this is equivalent to:

$$1 + (u-1)e_i + (u-1)e_i g_i = (u+1)^2 \ell_i^2 - 2(u+1)\ell_i + 1,$$

or, via Eq. 4.9, equivalently:

$$(u-1)(u+1)e_i \ell_i = (u+1)^2 \ell_i^2 - 2(u+1)\ell_i,$$

which leads to :

$$\ell_i^2 = \frac{(u-1)e_i + 2}{u+1} \ell_i.$$

Next we will prove Eq. 4.11. From Eq. 4.9 we obtain:

$$(4.18) \quad \begin{aligned} g_i g_{i+1} g_i &= (u+1)l_i l_{i+1} l_i - (u+1)^2 l_i^2 - (u+1)^2 l_{i+1} l_i \\ &\quad + (u+1)l_i - (u+1)^2 l_i l_{i+1} + (u+1)l_i \\ &\quad + (u+1)l_{i+1} - 1. \end{aligned}$$

$$(4.19) \quad \begin{aligned} g_{i+1} g_i g_{i+1} &= (u+1)l_{i+1} l_i l_{i+1} - (u+1)^2 l_{i+1}^2 - (u+1)^2 l_i l_{i+1} \\ &\quad + (u+1)l_{i+1} - (u+1)^2 l_{i+1} l_i + (u+1)l_{i+1} \\ &\quad + (u+1)l_i - 1. \end{aligned}$$

Equations 2.11, 4.18, 4.19 and 4.16 lead us to the desired result:

$$\begin{aligned} l_i l_{i+1} l_i - \frac{(u-1)e_i + 1}{(u+1)^2} l_i &= l_{i+1} l_i l_{i+1} - \frac{(u-1)e_{i+1} + 1}{(u+1)^2} l_{i+1}, \\ 1 \leq i \leq n-2. \end{aligned}$$

Finally, from relations $e_i e_{i+1} g_{i,i+1} = 0$ using Eqs. 4.1 and 4.9 we have for $1 \leq i \leq n-2$ that:

$$\begin{aligned} 0 &= e_i e_{i+1} g_{i,i+1} = e_i e_{i+1} (g_i g_{i+1} g_i + g_{i+1} g_i + g_i g_{i+1} + g_{i+1} + g_i + 1) \\ &= e_i e_{i+1} ((u+1)^3 l_i l_{i+1} l_i - (u+1)^2 l_i^2 + (u+1)l_i). \end{aligned}$$

From Eq. 4.16 we have that:

$$e_i e_{i+1} ((u+1)^2 l_i l_{i+1} l_i) = e_i e_{i+1} ((u-1)e_i + 1) l_i,$$

or equivalently:

$$e_i e_{i+1} l_i l_{i+1} l_i = \frac{u}{(u+1)^2} e_i e_{i+1} l_i,$$

which is Eq. 4.17. □

Proposition 5. *The algebra $\text{CTL}_{d,n}(u)$ can be presented with generators $\ell_1, \dots, \ell_{n-1}, t_1, \dots, t_n$, subject to the following relations:*

$$\begin{aligned} \ell_i \ell_j &= \ell_j \ell_i, \quad \text{for } |i - j| > 1 \\ \ell_i \ell_{i+1} \ell_i - \frac{(u-1)e_i + 1}{(u+1)^2} \ell_i &= \ell_{i+1} \ell_i \ell_{i+1} - \frac{(u-1)e_{i+1} + 1}{(u+1)^2} \ell_{i+1} \\ t_i^d &= 1, \quad t_i t_j = t_j t_i \\ \ell_i t_i &= t_{i+1} \ell_i + \frac{1}{u+1} (t_i - t_{i+1}) \\ \ell_i t_{i+1} &= t_i \ell_i + \frac{1}{u+1} (t_{i+1} - t_i) \\ \ell_i t_j &= t_j \ell_i, \quad \text{for } |i - j| > 1 \\ \ell_i^2 &= \frac{(u-1)e_i + 2}{u+1} \ell_i \\ \sum_{k=0}^{d-1} e_i^{(k)} e_{i+1} \ell_i \ell_{i+1} \ell_i &= \sum_{k=0}^{d-1} e_i^{(k)} e_{i+1} \frac{u}{(u+1)^2} \ell_i. \end{aligned}$$

Proof. The proof is a straightforward computation and totally analogous to the proof of Proposition 4. □

Remark 5. We know that a linear basis of the Temperley-Lieb algebra can be constructed from the interpretation of the generators ℓ_i as diagrams. In virtue of Remark 4, then it is desirable to construct a basis of $\text{FTL}_{d,n}(u)$ from the presentation given in Proposition 4. Unfortunately, we do not have a diagrammatic interpretation for the generators ℓ_i yet. In a recent result [5] Chlouveraki and Pouchin studied extensively the representation theories of the algebras $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$. Further, they provided linear bases for both and they also computed their dimensions. We will present here the dimensions of both of the algebras $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$. For this purpose, let $\text{Comp}_d(n) := \{ \mu = (\mu_1, \mu_2, \dots, \mu_d) \in \mathbb{N}^d \mid \mu_1 + \mu_2 + \dots + \mu_d = n \}$ and let also $c_k := \frac{1}{k+1} \binom{2k}{k}$ be the k -th Catalan number. We then have:

Theorem 4 ([5, Theorems 3.10, 5.5 and Remark 5.6]). *The dimension of the quotient algebra $\text{FTL}_{d,n}(u)$ is:*

$$(4.20) \quad \dim_{\mathbb{C}(u)} \text{FTL}_{d,n}(u) = \sum_{\mu \in \text{Comp}_d(n)} \left(\frac{n!}{\mu_1! \mu_2! \dots \mu_d!} \right)^2 c_{\mu_1} c_{\mu_2} \dots c_{\mu_d}.$$

The dimension of the quotient algebra $\text{CTL}_{d,n}(u)$ is:

$$\dim_{\mathbb{C}(u)} \text{CTL}_{d,n}(u) = \sum_{\mu \in \text{Comp}_d(n)} \left(\frac{n!}{\mu_1! \mu_2! \cdots \mu_d!} \right)^2 c_{\mu_1} \mu_2! \cdots \mu_d!.$$

4.5. Technical lemmas

We finish this section with two technical lemmas concerning the interaction with the braiding generators g_1, g_2 of the generators $g_{1,2}, r_{1,2}, c_{1,2}$ of the three ideals discussed above. Also, these lemmas will be used in the proof of Theorems 6 and 7.

Lemma 4. *For the element $g_{1,2}$ we have in $Y_{d,n}(u)$ the following:*

- (1) $g_1 g_{1,2} = [1 + (u - 1)e_1]g_{1,2}$
- (2) $g_2 g_{1,2} = [1 + (u - 1)e_2]g_{1,2}$
- (3) $g_1 g_2 g_{1,2} = [1 + (u - 1)e_1 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2]g_{1,2}$
- (4) $g_2 g_1 g_{1,2} = [1 + (u - 1)e_2 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2]g_{1,2}$
- (5) $g_1 g_2 g_1 g_{1,2} = [1 + (u - 1)(e_1 + e_2 + e_{1,3}) + (u - 1)^2 (u + 2) e_1 e_2]g_{1,2}.$

Proof. See [9, Lemma 5]. Cf. [14, Lemma 7.5]. □

Lemma 5. *For the element $r_{1,2}$ we have in $Y_{d,n}(u)$:*

- (1) $g_1 r_{1,2} = [1 + (u - 1)e_1]r_{1,2}$
- (2) $g_2 r_{1,2} = [1 + (u - 1)e_2]r_{1,2}$
- (3) $g_1 g_2 r_{1,2} = [1 + (u - 1)e_1 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2]r_{1,2}$
- (4) $g_2 g_1 r_{1,2} = [1 + (u - 1)e_2 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2]r_{1,2}$
- (5) $g_1 g_2 g_1 r_{1,2} = [1 + (u - 1)(e_1 + e_2 + e_{1,3}) + (u - 1)^2 (u + 2) e_1 e_2]r_{1,2}.$

Proof. For proving this lemma we will make extensive use of Lemmas 4 and 1. For statement (1) we have:

$$\begin{aligned} g_1 r_{1,2} &= g_1 e_1 e_2 g_{1,2} = e_1 e_{1,3} g_1 g_{1,2} \\ &= e_1 e_2 [1 + (u - 1)e_1] g_{1,2} \\ &= [1 + (u - 1)e_1] e_1 e_2 g_{1,2} \\ &= [1 + (u - 1)e_1] r_{1,2}. \end{aligned}$$

In an analogous way we prove statement (2). For statement (3) we have that:

$$\begin{aligned} g_1 g_2 r_{1,2} &= g_1 g_2 e_1 e_2 g_{1,2} = e_2 e_{1,3} g_1 g_2 g_{1,2} \\ &= e_1 e_2 [1 + (u - 1)e_1 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2] g_{1,2} \\ &= [1 + (u - 1)e_1 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2] e_1 e_2 g_{1,2} \\ &= [1 + (u - 1)e_1 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2] r_{1,2}. \end{aligned}$$

In an analogous way we prove statement (4). Finally, we have for statement (5):

$$\begin{aligned} g_1 g_2 g_1 r_{1,2} &= g_1 g_2 g_1 e_1 e_2 g_{1,2} \\ &= e_1 e_2 g_1 g_2 g_1 g_{1,2} \\ &= e_1 e_2 [1 + (u - 1)(e_1 + e_2 + e_{1,3}) + (u - 1)^2 (u + 2) e_1 e_2] g_{1,2} \\ &= [1 + (u - 1)(e_1 + e_2 + e_{1,3}) + (u - 1)^2 (u + 2) e_1 e_2] e_1 e_2 g_{1,2} \\ &= [1 + (u - 1)(e_1 + e_2 + e_{1,3}) + (u - 1)^2 (u + 2) e_1 e_2] r_{1,2}. \end{aligned}$$

□

Lemma 6. *For the element $c_{1,2}$ we have in $Y_{d,n}(u)$:*

- (1) $g_1 c_{1,2} = [1 + (u - 1)e_1] c_{1,2}$
- (2) $g_2 c_{1,2} = [1 + (u - 1)e_2] c_{1,2}$
- (3) $g_1 g_2 c_{1,2} = [1 + (u - 1)e_1 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2] c_{1,2}$
- (4) $g_2 g_1 c_{1,2} = [1 + (u - 1)e_2 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2] c_{1,2}$
- (5) $g_1 g_2 g_1 c_{1,2} = [1 + (u - 1)(e_1 + e_2 + e_{1,3}) + (u - 1)^2 (u + 2) e_1 e_2] c_{1,2}$.

Proof. The proof is completely analogous to the proof of Lemma 5. □

5. Markov traces

The main purpose of this section is to find the necessary and sufficient conditions in order that the trace tr defined on $Y_{d,n}(u)$ [13] passes to the quotient algebras $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$. Since the defining ideal of $\text{FTL}_{d,n}(u)$ (respectively of $\text{CTL}_{d,n}(u)$) is principal, by the linearity of tr , we have that tr passes to $\text{FTL}_{d,n}(u)$ (respectively to $\text{CTL}_{d,n}(u)$) if and only if we have:

$$(5.1) \quad \text{tr}(\mathfrak{m} r_{1,2}) = 0 \quad (\text{respectively} \quad \text{tr}(\mathfrak{m} c_{1,2}) = 0),$$

for all monomials \mathbf{m} in the inductive basis of $Y_{d,n}(u)$. So, we seek necessary and sufficient conditions for Eq. 5.1 to hold. The strategy is to find such conditions first for $n = 3$ and then to generalize using induction.

5.1. Computations on tr

Recall that elements in the inductive basis of $Y_{d,3}(u)$ are of the following forms:

$$(5.2) \quad t_1^a t_2^b t_3^c, \quad t_1^a g_1 t_1^b t_3^c, \quad t_1^a t_2^b g_2 g_1 t_1^c, \quad t_1^a t_2^b g_2 t_2^c, \quad t_1^a g_1 t_1^b g_2 t_2^c, \quad t_1^a g_1 t_1^b g_2 g_1 t_1^c,$$

where $0 \leq a, b, c \leq d - 1$ (see Proposition 1). We need now to compute the trace of the elements $\mathbf{m} r_{1,2}$, where \mathbf{m} runs the monomials listed in (5.2). To do these computations we will use the following lemma and proposition.

Lemma 7. *For all $0 \leq m \leq d - 1$, we have:*

$$\text{tr} \left(e_1^{(m)} e_2 g_{1,2} \right) = (u + 1)z^2 x_m + (u + 2)z E^{(m)} + \text{tr}(e_1^{(m)} e_2).$$

Proof. By direct computation we have:

$$\begin{aligned} & \text{tr} \left(e_1^{(m)} e_2 g_{1,2} \right) \\ &= \text{tr} \left(e_1^{(m)} e_2 g_1 \right) + \text{tr} \left(e_1^{(m)} e_2 g_2 \right) + \text{tr} \left(e_1^{(m)} e_2 g_1 g_2 \right) \\ & \quad + \text{tr} \left(e_1^{(m)} e_2 g_2 g_1 \right) + \text{tr} \left(e_1^{(m)} e_2 g_1 g_2 g_1 \right) + \text{tr} \left(e_1^{(m)} e_2 \right) \\ &= \frac{1}{d^2} \sum_{s=0}^{d-1} \sum_{k=0}^{d-1} \text{tr}(t_1^{m+s} t_2^{-s+k} t_3^{-k} g_1) + \frac{1}{d^2} \sum_{s=0}^{d-1} \sum_{k=0}^{d-1} \text{tr}(t_1^{m+s} t_2^{-s+k} t_3^{-k} g_2) \\ & \quad + \frac{1}{d^2} \sum_{s=0}^{d-1} \sum_{k=0}^{d-1} \text{tr}(t_1^{m+s} t_2^{-s+k} t_3^{-k} g_1 g_2) + \frac{1}{d^2} \sum_{s=0}^{d-1} \sum_{k=0}^{d-1} \text{tr}(t_1^{m+s} t_2^{-s+k} t_3^{-k} g_2 g_1) \\ & \quad + \frac{1}{d^2} \sum_{s=0}^{d-1} \sum_{k=0}^{d-1} \text{tr}(t_1^{m+s} t_2^{-s+k} t_3^{-k} g_1 g_2 g_1) + \text{tr} \left(e_1^{(m)} e_2 \right) \\ &= 2z E^{(m)} + 2z^2 x_m + z E^{(m)} + (u - 1)z E^{(m)} + (u - 1)z^2 x_m \\ &= (u + 1)z^2 x_m + (u + 2)z E^{(m)} + \text{tr} \left(e_1^{(m)} e_2 \right). \end{aligned}$$

□

Proposition 6. *For all $0 \leq a, b, c \leq d - 1$, we have:*

1) If $\mathbf{m} = t_1^a t_2^b t_3^c$,

$$\text{tr}(\mathbf{m}r_{1,2}) = (u + 1)z^2 x_{a+b+c} + (u + 2)E^{(a+b+c)}z + \text{tr}(e_1^{(a+b+c)} e_2)$$

2) If $\mathbf{m} = t_1^a g_1 t_1^b t_3^c$ and $\mathbf{m} = t_1^a t_2^b g_2 t_2^c$,

$$\text{tr}(\mathbf{m}r_{1,2}) = u \left[(u + 1)z^2 x_{a+b+c} + (u + 2)E^{(a+b+c)}z + \text{tr}(e_1^{(a+b+c)} e_2) \right]$$

3) If $\mathbf{m} = t_1^a t_2^b g_2 g_1 t_1^c$ and $\mathbf{m} = t_1^a g_1 t_1^b g_2 t_2^c$,

$$\text{tr}(\mathbf{m}r_{1,2}) = u^2 \left[(u + 1)z^2 x_{a+b+c} + (u + 2)E^{(a+b+c)}z + \text{tr}(e_1^{(a+b+c)} e_2) \right]$$

4) If $\mathbf{m} = t_1^a g_1 t_1^b g_2 g_1 t_1^c$,

$$\text{tr}(\mathbf{m}r_{1,2}) = u^3 \left[(u + 1)z^2 x_{a+b+c} + (u + 2)E^{(a+b+c)}z + \text{tr}(e_1^{(a+b+c)} e_2) \right].$$

Proof. We will prove claim (1). According to Eq. 4.4 we have: $\mathbf{m}r_{1,2} = t_1^a t_2^b t_3^c r_{1,2} = t_1^a t_2^b t_3^c e_1 e_2 g_{1,2}$. But $t_1^a t_2^b t_3^c e_1 e_2 = e_1^{(a+b+c)} e_2$, hence:

$$\mathbf{m}r_{1,2} = e_1^{(a+b+c)} e_2 g_{1,2}.$$

Thus, claim (1) follows by applying Lemma 7.

For proving the rest of the claims we use Lemmas 5 and 7 and we follow the same argument, so we finish the proof of the proposition by proving only one representative case. We shall prove claim (3) for $\mathbf{m} = t_1^a g_1 t_1^b g_2 t_2^c$. This monomial can be rewritten as $t_1^a t_2^b t_3^c g_1 g_2$. Now, by using Lemma 5 on $g_1 g_2 r_{1,2}$, we obtain:

$$\begin{aligned} \mathbf{m}r_{1,2} &= t_1^a t_2^b t_3^c g_1 g_2 r_{1,2} \\ &= t_1^a t_2^b t_3^c [1 + (u - 1)e_1 + (u - 1)e_{1,3} + (u - 1)^2 e_1 e_2] r_{1,2}, \end{aligned}$$

then using now Eq. 4.4 and the fact the e_i 's are idempotents, it follows that:

$$\begin{aligned} \mathbf{m}r_{1,2} &= t_1^a t_2^b t_3^c [e_1 e_2 + (u - 1)e_1 e_2 + (u - 1)e_1 e_2 + (u - 1)^2 e_1 e_2] g_{1,2} \\ &= u^2 t_1^a t_2^b t_3^c e_1 e_2 g_{1,2}. \end{aligned}$$

Then, applying Eq. 2.21 we have:

$$\mathbf{m}r_{1,2} = u^2 t_1^a t_2^b t_3^c e_1 e_2 g_{1,2} = u^2 e_1^{(a+b+c)} e_2 g_{1,2}.$$

Therefore, by using Lemma 7, we obtain the desired expression for $\text{tr}(\mathbf{m}r_{1,2})$. □

5.2. Passing tr to the algebra $\text{YTL}_{d,n}(u)$

In [9] we found the necessary and sufficient conditions so that tr passes to $\text{YTL}_{d,n}(u)$. Indeed, we have the following:

Theorem 5 ([9, Theorem 6]). *The trace tr passes to the quotient algebra $\text{YTL}_{d,n}(u)$ if and only if the x_i 's are solutions of the E-system and one of the two cases holds:*

- (i) the x_ℓ 's are d^{th} roots of unity and $z = -\frac{1}{u+1}$ or $z = -1$,
- (ii) the x_ℓ 's are the solutions of the E-system that are parametrized by the set $D = \{m_1, m_2 \mid 0 \leq m_1, m_2 \leq d - 1 \text{ and } m_1 \neq m_2\}$ and they are expressed as:

$$x_\ell = \frac{1}{2} \left(\chi_{m_1}(t^\ell) + \chi_{m_2}(t^\ell) \right), \quad 0 \leq \ell \leq d - 1.$$

In this case we have that $z = -\frac{1}{2}$.

5.3. Passing tr to the algebra $\text{FTL}_{d,n}(u)$

The following lemma is key to proving one of our main results (Theorem 6). Recall that the *support* of a function $x : C_d \rightarrow \mathbb{C}$ (or equivalently of an element $\sum_{k=0}^{d-1} x(t^k)t^k \in \mathbb{C}C_d$) is the subset of C_d where the values of x are non-zero.

Lemma 8. *The trace tr passes to $\text{FTL}_{d,3}(u)$ if and only if the parameters of the trace tr satisfy:*

$$x_k = -z \left(\sum_{m \in \text{Sup}_1} \chi_m(t^k) + (u + 1) \sum_{m \in \text{Sup}_2} \chi_m(t^k) \right)$$

and
$$z = -\frac{1}{|\text{Sup}_1| + (u + 1)|\text{Sup}_2|},$$

where $\text{Sup}_1 \cup \text{Sup}_2$ (disjoint union) is the support of the Fourier transform of x , and x is the complex function on C_d , that maps 0 to 1 and k to the trace parameter x_k (cf. Section 3.2).

Proof. Recall that the trace tr passes to $\text{FTL}_{d,3}$ if and only if the Eqs. 5.1 hold, for all \mathbf{m} in the inductive basis of $Y_{d,3}$. By using Proposition 6 follows

that the trace tr passes to the quotient algebra $\text{FTL}_{d,3}(u)$ if and only if the trace parameters z, x_1, \dots, x_{d-1} satisfy the following system of equations:

$$\mathbb{E}_0 = \mathbb{E}_1 = \dots = \mathbb{E}_{d-1} = 0,$$

where

$$\mathbb{E}_m := (u + 1)z^2x_m + (u + 2)E^{(m)}z + \text{tr}(e_1^{(m)}e_2) = 0, \quad 0 \leq m \leq d - 1.$$

We note now that this system of equations above is equivalent to the system:

$$(5.3) \quad \begin{aligned} \mathbb{E}_0 &= 0 \\ \mathbb{E}_m - x_m\mathbb{E}_0 &= 0 \quad \text{where } 1 \leq m \leq d - 1. \end{aligned}$$

We will solve this system of equations, obtaining thus the proof of the lemma.

Recall that $x_0 := 1, E^{(0)} = E$ and $e_i^{(0)} = e_i$, hence $\mathbb{E}_0 = (u + 1)z^2 + (u + 2)Ez + \text{tr}(e_1e_2)$. Then the $(d - 1)$ equations $\mathbb{E}_m - x_m\mathbb{E}_0 = 0$ of Eq. 5.3 become:

$$(5.4) \quad \begin{aligned} z(u + 2)(E^{(m)} - x_mE) &= - \left(\text{tr}(e_1^{(m)}e_2) - x_m \text{tr}(e_1e_2) \right), \\ 1 \leq m \leq d - 1. \end{aligned}$$

Interpreting now the above equation in the functional notation of Section 3 and having in mind Eq. 3.8, it follows that Eq. 5.4 can be rewritten as:

$$(u + 2)z \left(\frac{1}{d}x * x - Ex \right) = - \left(\frac{1}{d^2}x * x * x - \text{tr}(e_1e_2)x \right).$$

Applying now the Fourier transform on the above functional equality and using Proposition 2, we obtain:

$$(5.5) \quad (u + 2)z \left(\frac{\widehat{x}^2}{d} - E\widehat{x} \right) = - \left(\frac{\widehat{x}^3}{d^2} - \text{tr}(e_1e_2)\widehat{x} \right).$$

Let now $\widehat{x} = \sum_{m=0}^{d-1} y_m t^m$. Then Eq. 5.5 becomes:

$$(u + 2)z \left(\frac{y_m^2}{d} - Ey_m \right) = - \left(\frac{y_m^3}{d^2} - \text{tr}(e_1e_2)y_m \right).$$

Hence

$$(5.6) \quad y_m \left(\frac{y_m^2}{d^2} + (u + 2)z \frac{y_m}{d} - (u + 2)zE - \text{tr}(e_1e_2) \right) = 0.$$

Now, from equation $\widehat{\mathbb{E}}_0 = 0$, we have that $-(u+2)zE = (u+1)z^2 + \text{tr}(e_1e_2)$. Replacing this expression of $-(u+2)zE$ in Eq. 5.6 we have that:

$$y_m \left(\frac{y_m^2}{d^2} + (u+2)z \frac{y_m}{d} + (u+1)z^2 \right) = 0,$$

or equivalently (notice that the equivalence still holds even if we specialize $u = -1$, where the above equation is not quadratic):

$$(5.7) \quad y_m (y_m + dz) (y_m + dz(u+1)) = 0.$$

Denote $\text{Sup}_1 \cup \text{Sup}_2$ the support of \widehat{x} , where

$$\begin{aligned} \text{Sup}_1 &:= \{m \in C_d; y_m = -dz\} \\ \text{and } \text{Sup}_2 &:= \{m \in C_d; y_m = -dz(u+1)\}, \end{aligned}$$

hence

$$\widehat{x} = \sum_{m \in \text{Sup}_1} -dzt^m + \sum_{m \in \text{Sup}_2} -dz(u+1)t^m.$$

Notice again that if specialize $u = -1$, then the support of \widehat{x} is just Sup_1 . Then

$$\widehat{\widehat{x}} = -dz \sum_{m \in \text{Sup}_1} \widehat{\delta}_m - dz(u+1) \sum_{m \in \text{Sup}_2} \widehat{\delta}_m,$$

thus from argument (4) of Proposition 2 we have:

$$\widehat{\widehat{\widehat{x}}} = -z \left(\sum_{m \in \text{Sup}_1} \mathbf{i}_{-m} + (u+1) \sum_{m \in \text{Sup}_2} \mathbf{i}_{-m} \right).$$

Therefore, having in mind now (5) of Proposition 2, we deduce that:

$$(5.8) \quad x_k = -z \left(\sum_{m \in \text{Sup}_1} \chi_m(t^k) + (u+1) \sum_{m \in \text{Sup}_2} \chi_m(t^k) \right).$$

Having in mind that $x_0 = 1$, one can determine the values of z . Indeed, from Eq. 5.8, we have that:

$$(5.9) \quad 1 = x_0 = -z(|\text{Sup}_1| + (u+1)|\text{Sup}_2|),$$

or equivalently (keep in mind that the assumption $x_0 = 1$ forces the denominator to be non-zero and hence the support of \hat{x} is not empty):

$$(5.10) \quad z = -\frac{1}{|\text{Sup}_1| + (u + 1)|\text{Sup}_2|}.$$

By the same reasoning z is also non-zero. □

Keeping the same notation with the above lemma, we have:

Theorem 6. *The trace tr defined on $Y_{d,n}(u)$ passes to the quotient algebra $\text{FTL}_{d,n}(u)$ if and only if the trace parameters z, x_1, \dots, x_{d-1} satisfy the conditions of Lemma 8, namely Eqs. 5.8 and 5.10.*

Proof. The proof is by induction on n . The case $n = 3$ is the lemma above. Assume now that the statement holds for all $\text{FTL}_{d,k}(u)$, where $k \leq n$, that is:

$$\text{tr}(a_k r_{1,2}) = 0,$$

for all $a_k \in Y_{d,k}(u)$, $k \leq n$. We will show the statement for $k = n + 1$. It suffices to prove that the trace vanishes on any element of the form $a_{n+1} r_{1,2}$, where a_{n+1} belongs to the inductive basis of $Y_{d,n+1}(u)$ (recall Eq. 1), given the conditions of the theorem. Namely:

$$\text{tr}(a_{n+1} r_{1,2}) = 0.$$

Since a_{n+1} is in the inductive basis of $Y_{d,n+1}(u)$, it is of one of the following forms:

$$a_{n+1} = a_n g_n \cdots g_i t_i^k \quad \text{or} \quad a_{n+1} = a_n t_{n+1}^k,$$

where a_n is in the inductive basis of $Y_{d,n}(u)$. For the first case we have:

$$\text{tr}(a_{n+1} r_{1,2}) = \text{tr}(a_n g_n \cdots g_i t_i^k r_{1,2}) = z \text{tr}(a_n g_{n-1} \cdots g_i t_i^k r_{1,2}) = z \text{tr}(w r_{1,2}),$$

where $w := a_n g_{n-1} \cdots g_i t_i^k$. Notice now that w is a word in $Y_{d,n}(u)$ and so, by the linearity of the trace, we have that $\text{tr}(w r_{1,2})$ is a linear combination of traces of the form $\text{tr}(a_n r_{1,2})$, where a_n is in the inductive basis of $Y_{d,n}(u)$. Therefore, by the induction hypothesis, we deduce that:

$$\text{tr}(w r_{1,2}) = 0,$$

if and only if the conditions of the Theorem are satisfied. Therefore the statement is proved. The second case is proved similarly. Hence, the proof is concluded. □

Corollary 3. *In the case where one of the sets Sup_1 or Sup_2 is the empty set, the values of the x_k 's are solutions of the E-system. More precisely, if Sup_1 is the empty set, the x_k 's are the solutions of the E-system parametrized by Sup_2 and $z = -1/(u + 1)|\text{Sup}_2|$. If Sup_2 is the empty set, then x_k 's are the solutions of the E-system parametrized by Sup_1 and $z = -1/|\text{Sup}_1|$.*

Proof. The proof follows from Eq. 3.6 and the expression given in theorem above for the x_k 's. □

5.4. Passing tr to the algebra $\text{CTL}_{d,n}(u)$

The method for finding the necessary and sufficient conditions for tr to pass to the quotient algebra $\text{CTL}_{d,n}(u)$ is completely analogous to that of the previous subsection. So, we will need the following analogue of Proposition 6.

Proposition 7. *Define \mathbb{G} , as follows:*

$$\mathbb{G} = (u + 1)z^2 \sum_{k=0}^{d-1} x_k + (u + 2)z \sum_{k=0}^{d-1} E^{(k)} + \sum_{k=0}^{d-1} \text{tr}(e_1^{(k)} e_2).$$

Then for all $0 \leq a, b, c \leq d - 1$, we have:

- (1) $\text{tr}(\mathbf{m}c_{1,2}) = \mathbb{G}$ for $\mathbf{m} = t_1^a t_2^b t_3^c$
- (2) $\text{tr}(\mathbf{m}c_{1,2}) = u\mathbb{G}$ for $\mathbf{m} = t_1^a g_1 t_1^b t_3^c$ and $\mathbf{m} = t_1^a t_2^b g_2 t_2^c$
- (3) $\text{tr}(\mathbf{m}c_{1,2}) = u^2\mathbb{G}$ for $\mathbf{m} = t_1^a t_2^b g_2 g_1 t_1^c$ and $\mathbf{m} = t_1^a g_1 t_1^b g_2 t_2^c$
- (4) $\text{tr}(\mathbf{m}c_{1,2}) = u^3\mathbb{G}$ for $\mathbf{m} = t_1^a g_1 t_1^b g_2 g_1 t_1^c$.

Following now the analogous reasoning that was used to prove Theorem 6 and having in mind Eq. 5.1, Corollary 2, Lemma 6 and Proposition 7, we obtain the following theorem.

Theorem 7. *The trace tr passes to the quotient algebra $\text{CTL}_{d,n}(u)$ if and only if the parameter z and the x_i 's are related through the equation:*

$$(5.11) \quad (u + 1)z^2 \sum_{k \in \mathbb{Z}/d\mathbb{Z}} x_k + (u + 2)z \sum_{k \in \mathbb{Z}/d\mathbb{Z}} E^{(k)} + \sum_{k \in \mathbb{Z}/d\mathbb{Z}} \text{tr}(e_1^{(k)} e_2) = 0.$$

5.5. Comparison of the three trace conditions

In this section we will compare the conditions that need to be applied to the trace paramaters z and $x_i, i = 1, \dots, d - 1$ so that tr passes to each one of the quotient algebras $\text{YTL}_{d,n}(u)$, $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$.

By comparing Theorem 5 and Theorem 6, we observe that the conditions such that tr passes to $\text{YTL}_{d,n}(u)$ are contained in the conditions such that tr passes to $\text{FTL}_{d,n}(u)$.

Moreover, Theorem 6 can be rephrased in the following way:

Theorem 8. *The trace tr passes to the quotient algebra $\text{FTL}_{d,n}(u)$ if and only if the parameter z and the x_i 's are related through the equation:*

$$(u + 1)z^2x_k + (u + 2)zE^{(k)} + \text{tr}(e_1^{(k)}e_2) = 0, \quad k \in \mathbb{Z}/d\mathbb{Z}.$$

This implies that the conditions such that the trace passes to the quotient algebra $\text{FTL}_{d,n}(u)$ are contained in those of Theorem 7.

All of the above can be summarized in the following table:

	$\text{Y}_{d,n}(u)$	\rightarrow	$\text{CTL}_{d,n}(u)$	\rightarrow	$\text{FTL}_{d,n}(u)$	\rightarrow	$\text{YTL}_{d,n}(u)$
z	free						
x_i	free	\leftrightarrow	Theorem 7	\leftrightarrow	Theorem 8	\leftrightarrow	Theorem 5

Table 1: Relations of the algebras and the trace conditions.

The first row includes the projections between the algebras while the second shows the inclusions of the trace conditions for each case.

Remark 6. By Theorems 6 and 8 and by Corollary 3, the necessary and sufficient conditions for the trace tr to pass to $\text{FTL}_{d,n}(u)$ include the solutions of the E-system, leading directly to link invariants derived from this algebra (see Section 6). On the other hand, the conditions on the x_i 's for the algebra $\text{CTL}_{d,n}(u)$ are too loose as indicated by Theorem 7. Moreover, as we shall see in Section 6 the resulting invariants from $\text{CTL}_{d,n}(u)$ coincide either with invariants from $\text{Y}_{d,n}(u)$ or with invariants from $\text{FTL}_{d,n}(u)$. For these reasons, the algebra $\text{CTL}_{d,n}(u)$ will be discarded as a possible framization of the Temperley-Lieb algebra.

6. Knot and link invariants

In this section we define framed and classical link invariants related to the algebras $\text{FTL}_{d,n}(u)$ and $\text{CTL}_{d,n}(u)$, using the results of the previous sections. The general scheme for defining these invariants follows Jones' method [8, 11]. More precisely, one uses the (framed) braid equivalence corresponding to

(framed) link isotopy, the mapping of the (framed) braid group to the knot algebra in question and the Markov trace on this algebra, which, upon re-scaling and normalization according to the braid equivalence, yields isotopy invariants of (framed) links.

6.1. The Homflypt and the Jones polynomials

It is known that by re-scaling and normalizing the Ocneanu trace τ on $H_n(u)$, one can define the 2-variable Jones or Homflypt polynomial, $P(\lambda_H, u)$ [11]. Namely, we have:

$$P(\lambda_H, u)(\hat{\alpha}) = \left(-\frac{1 - \lambda_H u}{\sqrt{\lambda_H}(1 - u)} \right)^{n-1} \left(\sqrt{\lambda_H} \right)^{\varepsilon(\alpha)} \tau(\pi(\alpha)),$$

where: $\alpha \in \cup_{\infty} B_n$, $\lambda_H = \frac{1-u+\zeta}{u\zeta}$ is the “re-scaling factor”, π is the natural epimorphism of $\mathbb{C}(u)B_n$ on $H_n(u)$ that sends the braid generator σ_i to h_i and $\varepsilon(\alpha)$ is the algebraic sum of the exponents of the σ_i ’s in α . Further, by specializing ζ to $-\frac{1}{u+1}$, the non-trivial value for which the Ocneanu trace τ passes to the quotient algebra $TL_n(u)$, the Jones polynomial, $V(u)$, can be defined through the Homflypt polynomial [11], as follows:

$$V(u)(\hat{\alpha}) = \left(-\frac{1 + u}{\sqrt{u}} \right)^{n-1} (\sqrt{u})^{\varepsilon(\alpha)} \tau(\pi(\alpha)) = P(u, u)(\hat{\alpha}).$$

6.2. Invariants from $Y_{d,n}(u)$

In [19] it is proved that the trace tr defined on $Y_{d,n}(u)$ can be re-scaled according to the braid equivalence corresponding to isotopic framed links if and only if the framing parameters x_i ’s of tr furnish a solution of the E-system (recall discussion in Section 3). Let $X_D = (x_1, \dots, x_{d-1})$ be a solution of the E-system parametrized by the non-empty set D of $\mathbb{Z}/d\mathbb{Z}$. We have the following definition:

Definition 6 ([4, Definition 3]). The trace map tr_D defined as the trace tr with the parameters x_i specialized to the values x_i , shall be called the *specialized trace* with parameter z .

Note that for $d = 1$ the traces tr and tr_D coincide with the Ocneanu trace. By normalizing tr_D , an invariant for *framed links* is obtained [19]:

$$(6.1) \quad \Gamma_{d,D}(w, u)(\hat{\alpha}) = \left(-\frac{(1-wu)|D|}{\sqrt{w}(1-u)} \right)^{n-1} (\sqrt{w})^{\varepsilon(\alpha)} \text{tr}_D(\gamma(\alpha)),$$

where: $w = \frac{z+(1-u)E}{uz}$ is the re-scaling factor, $E = \frac{1}{|D|}$ [16, 19], γ is the natural epimorphism of the framed braid group algebra $\mathbb{C}(u)\mathcal{F}_n$ on the algebra $Y_{d,n}(u)$, and $\alpha \in \cup_{\infty}\mathcal{F}_n$.

Further, by restricting the invariants $\Gamma_{d,D}(w, u)$ to *classical links*, seen as framed links with all framings zero, in [16] invariants of classical oriented links $\Delta_{d,D}(w, u)$ are obtained.

In [4] it was proved that for generic values of the parameters u, z the invariants $\Delta_{d,D}(w, u)$ do not coincide with the Homflypt polynomial except in the trivial cases $u = 1$ and $E = 1$. More details are given in Section 7.

6.3. Invariants from $\text{YTL}_{d,n}(u)$

In [9] the invariants that are defined through the Yokonuma-Temperley-Lieb were studied. More precisely, it was shown that in order that the trace tr passes to the quotient algebra $\text{YTL}_{d,n}(u)$ it is necessary that the x_i 's are d^{th} roots of unity. These furnish a (trivial) solution of the E-system and in this case $E = 1$. By [19, Remark 5] the framed link invariants $\mathcal{V}_D(u)$ that are derived from $\text{YTL}_{d,n}(u)$ are not very interesting, since basic pairs of framed links are not distinguished. On the other hand, by [4], the classical link invariants $V_D(u)$ that we obtain from $\text{YTL}_{d,n}(u)$ coincide with the Jones polynomial. This is the main reason that the algebra $\text{YTL}_{d,n}(u)$ does not qualify for being the framization of the Temperley-Lieb algebra.

6.4. Invariants from $\text{FTL}_{d,n}(u)$

As it has already been stated, the trace parameters x_i should be solutions of the E-system so that a link invariant through tr is well-defined. Recall that, the conditions of Theorem 6 include these solutions for the x_i 's. So, in order to define link invariants on the level of the quotient algebra $\text{FTL}_{d,n}(u)$, we discard any values of the x_i 's that do not furnish a solution of the E-system. For a solution of the E-system parametrized by $D \subset \mathbb{Z}/d\mathbb{Z}$, using Corollary 3, we have the following:

Proposition 8. *The specialized trace tr_D passes to the quotient algebra $\text{FTL}_{d,n}(u)$ if and only if:*

$$z = -\frac{1}{(u+1)|D|} \quad \text{or} \quad z = -\frac{1}{|D|}.$$

Remark 7. As argued also in [11] for the classical case of $\text{TL}_n(u)$, we do not take into consideration the case where $z = -\frac{1}{|D|}$, since important topological information is lost. For example, the trace tr_D gives the same value for all even (resp. odd) powers of the g_i 's, for $m \in \mathbb{Z}^{>0}$ [19]:

$$(6.2) \quad \text{tr}_D(g_i^m) = \left(\frac{u^m - 1}{u + 1}\right) z + \left(\frac{u^m - 1}{u + 1}\right) \frac{1}{|D|} + 1 \quad \text{if } m \text{ is even}$$

and

$$(6.3) \quad \text{tr}_D(g_i^m) = \left(\frac{u^m - 1}{u + 1}\right) z + \left(\frac{u^m - 1}{u + 1}\right) \frac{1}{|D|} - \frac{1}{|D|} \quad \text{if } m \text{ is odd,}$$

so the corresponding knots and links are not distinguished.

From the remaining case, where the x_i 's are solutions of the E-system and $z = -\frac{1}{(u+1)|D|}$, we deduce for the rescaling factor w that appears in Eq. 6.1 that $w = u$. We then have the following definition:

Definition 7. Let X_D be a solution of the E-system, parametrized by the non-empty subset D of $\mathbb{Z}/d\mathbb{Z}$ and let $z = -\frac{1}{(u+1)|D|}$. We obtain from $\Gamma_{d,D}(w, u)$ the following 1-variable framed link invariant:

$$\Gamma_{d,D}(u, u)(\hat{\alpha}) := \left(-\frac{(1+u)|D|}{\sqrt{u}}\right)^{n-1} (\sqrt{u})^{\varepsilon(\alpha)} \text{tr}_D(\gamma(\alpha)),$$

for any $\alpha \in \cup_{\infty} \mathcal{F}_n$. Further, in analogy to the invariants of $\Gamma_{d,D}(w, u)$, if we restrict to framed links with all framings zero, we obtain from $\Gamma_{d,D}(u, u)$ an 1-variable invariant of classical links $\Delta_{d,D}(u, u)$.

6.5. Invariants from $\text{CTL}_{d,n}(u)$

The conditions of Theorem 7 do not involve the solutions of the E-system at all, so in order to obtain a well-defined link invariant on the level of $\text{CTL}_{d,n}(u)$ we must impose E-condition on the x_i 's. Recall that the solutions

of the E-system can be expressed in the form:

$$x_D = \frac{1}{|D|} \sum_{k \in D} \mathbf{i}_k \in \mathbb{C}C_d,$$

where $\mathbf{i}_k = \sum_{j=0}^{d-1} \chi_k(t^j)t^j$, χ_k is the character that sends $t^m \mapsto \cos \frac{2\pi km}{d} + i \sin \frac{2\pi km}{d}$ and D is the subset of $\mathbb{Z}/d\mathbb{Z}$ that parametrizes this solution of the E-system. Let now ε be the augmentation function of the group algebra $\mathbb{C}C_d$, sending $\sum_{j=0}^{d-1} x_j t^j$ to $\sum_{j=0}^{d-1} x_j$. We have that:

$$(6.4) \quad \varepsilon(x_D) = \frac{1}{|D|} \sum_{k \in D} \varepsilon(\mathbf{i}_k) = \frac{1}{|D|} \sum_{j=0}^{d-1} \sum_{k \in D} \chi_k(t^j) = \begin{cases} \frac{d}{|D|}, & \text{if } 0 \in D \\ 0, & \text{if } 0 \notin D \end{cases}.$$

From this we deduce that:

$$(6.5) \quad \begin{aligned} \sum_{j=0}^{d-1} E^{(j)} &= \varepsilon\left(\frac{x * x}{d}\right) = \frac{1}{d|D|^2} \sum_{k \in D} \varepsilon(\mathbf{i}_k * \mathbf{i}_k) \\ &= \frac{1}{|D|^2} \sum_{k \in D} \varepsilon(\mathbf{i}_k) = \begin{cases} \frac{d}{|D|^2}, & \text{if } 0 \in D \\ 0, & \text{if } 0 \notin D \end{cases} \end{aligned}$$

and also that:

$$(6.6) \quad \begin{aligned} \sum_{j=0}^{d-1} \text{tr}(e_1^{(j)} e_2) &= \varepsilon\left(\frac{x * x * x}{d^2}\right) = \frac{1}{d^2|D|^3} \sum_{k \in D} \varepsilon(\mathbf{i}_k * \mathbf{i}_k * \mathbf{i}_k) \\ &= \frac{1}{|D|^3} \sum_{k \in D} \varepsilon(\mathbf{i}_k) = \begin{cases} \frac{d}{|D|^3}, & \text{if } 0 \in D \\ 0, & \text{if } 0 \notin D \end{cases} \end{aligned}$$

Using now Eqs. 6.4–6.6, we have that Eq. 5.11, for the case where $0 \in D$, becomes :

$$(6.7) \quad \frac{d}{|D|} \left((u+1)z^2 + \frac{(u+2)}{|D|}z + \frac{1}{|D|^2} \right) = 0.$$

Notice also that for the case where $0 \notin D$, Eq. 5.11 vanishes. We thus have the following:

Proposition 9. *Assume that the x_i 's are restricted to solutions of the E-system. Then, the specialized trace tr_D passes to the quotient algebra $\text{CTL}_{d,n}(u)$ if and only if one of the following cases hold:*

(i) When $0 \in D$, the trace parameter z takes the values:

$$z = -\frac{1}{(u+1)|D|} \quad \text{or} \quad z = -\frac{1}{|D|}.$$

(ii) When $0 \notin D$, the trace parameter z is free.

We will discuss now the invariants that are derived from the algebras $\text{CTL}_{d,n}(u)$.

Case (i) $0 \in D$. In this case the values for z in case (i) of Proposition 9 coincide with the values for z in Proposition 8. Further, much like the case of $\text{FTL}_{d,n}(u)$, the value $z = -\frac{1}{|D|}$ is not taken into consideration. Thus, the invariants that are obtained from tr_D on the level of the quotient algebra $\text{CTL}_{d,n}(u)$, for subsets D containing zero, coincide with the corresponding invariants $\Gamma_{d,D}(u, u)$ and $\Delta_{d,D}(u, u)$ derived from $\text{FTL}_{d,n}(u)$, since the conditions that are applied to the trace parameters are the same for both quotient algebras.

Case (ii) $0 \notin D$. In this case z remains an indeterminate. Thus, the only condition that is required so that the trace tr_D passes to the quotient algebra, is that the x_i 's comprise a solution of the E-system. This means that the invariants that are derived from the quotient algebra $\text{CTL}_{d,n}(u)$, for subsets D not containing zero, coincide with the corresponding invariants $\Gamma_{d,D}(w, u)$ and $\Delta_{d,D}(w, u)$ that are derived from $Y_{d,n}(u)$. We have thus proved the following:

Proposition 10. *Let X_D be a solution of the E-system parametrized by $D \subset \mathbb{Z}/d\mathbb{Z}$. The invariants derived from the algebra $\text{CTL}_{d,n}(u)$:*

(i) *if $0 \in D$, they coincide with the invariants derived from the algebra $\text{FTL}_{d,n}(u)$ and*

(ii) *if $0 \notin D$, they coincide with the invariants derived from the algebra $Y_{d,n}(u)$.*

Remark 8. As we see from the above, the type of invariants we obtain from the algebra $\text{CTL}_{d,n}(u)$ depends on whether zero belongs or not to the parametrizing set D of the specific solution of the E-system. The intrinsic reason for this peculiar condition on the set D is the fact that when summing up all n -roots of unity we get zero, unless $n = 1$, see Eq. 6.4.

Remark 9. The results of Propositions 9 and 10 seem to be in accordance with the recent results of Chlouveraki and Pouchin [6], where they prove

that the algebra is isomorphic to a direct sum of matrix algebras over tensor products of Temperley-Lieb and Iwahori-Hecke algebras.

To summarize, the solutions of the E-system (which are the necessary and sufficient conditions so that topological invariants for framed links can be defined) are included in the conditions of Theorem 6, while for the case of $\text{CTL}_{d,n}(u)$ we still have to impose them. Even by doing so, this algebra does not deliver any new invariants for (framed) links. This is the main reason that led us to consider the quotient algebra $\text{FTL}_{d,n}(u)$ as the most natural non-trivial analogue of the Temperley-Lieb algebra in the context of framed links.

We conclude this section with presenting the following tables that give an overview of the invariants for each quotient algebra:

$d, D > 1$	$Y_{d,n}(u)$	$\text{CTL}_{d,n}(u)$		$\text{FTL}_{d,n}(u)$	$\text{YTL}_{d,n}(u)$
		$0 \notin D$	$0 \in D$		
$\mathcal{F}_{d,n}$	$\Gamma_{d,D}(w, u)$	$\Gamma_{d,D}(w, u)$	$\Gamma_{d,D}(u, u)$	$\Gamma_{d,D}(u, u)$	–
B_n	$\Delta_{d,D}(w, u)$	$\Delta_{d,D}(w, u)$	$\Delta_{d,D}(u, u)$	$\Delta_{d,D}(u, u)$	–

Table 2: Overview of the invariants for $|D| > 1$.

$d, D = 1$	$Y_{d,n}(u)$	$\text{CTL}_{d,n}(u)$		$\text{FTL}_{d,n}(u)$	$\text{YTL}_{d,n}(u)$
		$0 \notin D$	$0 \in D$		
$\mathcal{F}_{d,n}$	$\Gamma_{d,D}(w, u)$	$\Gamma_{d,D}(w, u)$	$\mathcal{V}_D(u)$	$\mathcal{V}_D(u)$	$\mathcal{V}_D(u)$
B_n	$P(\lambda, u)$	$P(\lambda, u)$	$V_D(u)$	$V_D(u)$	$V_D(u)$

Table 3: Overview of the invariants for $|D| = 1$.

7. Identifying the invariants from $\text{FTL}_{d,n}(u)$ on classical knots and links

It has been a long standing problem how the classical link invariants derived from the Yokonuma-Hecke algebras compare with other known invariants, especially with the Homflypt polynomial. Finally, in a recent development [3] it is proved that these invariants are topologically equivalent to the Homflypt polynomial on *knots* but not on *links*. For proving these results, a different presentation for the algebra $Y_{d,n}(u)$ was employed, leading to classical link

invariants denoted by $\Theta_{d,D}$. As proved in [3] the invariants $\Theta_{d,D}$ do not depend on the sets D , so the notation was simplified. More precisely, by results in [3], the specialized trace tr_D on classical knots and links depends only on $|D|$ and not on the solution X_D of the E-system. Further, for d, d' positive integers with $d \leq d'$, we have $\Theta_{d,D} = \Theta_{d',D'}$ as long as $|D| = |D'|$. We deduce that, if $|D'| = d$, then $\Theta_{d',D'} = \Theta_{d,\mathbb{Z}/d\mathbb{Z}}$. Therefore, the invariants $\Theta_{d,D}$ can be parametrized by the natural numbers, setting $\Theta_d := \Theta_{d,\mathbb{Z}/d\mathbb{Z}}$ for all $d \in \mathbb{Z}_{>0}$ [3, Proposition 4.6]. However, to avoid confusion, we will keep the notation tr_D for the specialized trace, keeping in mind that $D = \mathbb{Z}/d\mathbb{Z}$.

In order to compare the classical link invariants from the Framization of the Temperley-Lieb algebra with the Jones polynomial we will consider in this section a new presentation for this algebra (according to [3]) and we will adapt our results so far.

7.1. A different presentation for $\mathbf{H}_n(\mathbf{u})$ and $\mathbf{Y}_{d,n}(\mathbf{u})$

The Iwahori-Hecke algebra is generated by the elements h'_1, \dots, h'_{n-1} satisfying the relations $h'_i h'_j = h'_j h'_i$, for $|i - j| > 1$ and $h'_i h'_{i+1} h'_i = h'_{i+1} h'_i h'_{i+1}$, together with the quadratic relations: $(h'_i)^2 = 1 + (q - q^{-1})h'_i$. The transformation from the presentation that was given in Section 2 to this one can be achieved by taking $u = q^2$ and $h_i = qh'_i$ [3]. Consequently, the defining ideal (2.6) for the algebra $\text{TL}_n(q)$ becomes:

$$1 + q(h'_1 + h'_2) + q^2(h'_1 h'_2 + h'_2 h'_1) + q^3 h'_1 h'_2 h'_1.$$

Further, the Ocneanu trace τ passes to the quotient algebra for the following values of the trace parameter ζ' :

$$\zeta' = -\frac{q^{-1}}{q^2 + 1} \quad \text{or} \quad \zeta' = -q^{-1}.$$

On the other hand, the algebra $\mathbf{Y}_{d,n}(q)$ is generated by the elements $g'_1, \dots, g'_{n-1}, t_1, \dots, t_n$, satisfying the relations (2.10)–(2.16) and the quadratic relations [3]:

$$(7.1) \quad (g'_i)^2 = 1 + (q - q^{-1})e_i g'_i.$$

One can obtain this presentation from the one given in Definition 2 by taking $u = q^2$ and

$$(7.2) \quad g_i = g'_i + (q - 1)e_i g'_i \quad (\text{or, equivalently, } g'_i = g_i + (q^{-1} - 1)e_i g_i).$$

Further, on the algebra $Y_{d,n}(q)$ a unique Markov trace is defined, analogous to tr , satisfying the same rules [3], for which we retain here the same notation. Note also that, the E-system remains the same for $Y_{d,n}(q)$ so we can talk about the specialized trace tr_D [3]. Consequently, in [3], invariants for framed links were derived which restrict to invariants of classical links:

$$(7.3) \quad \Theta_d(\lambda_d, q)(\widehat{\alpha}) = \left(\frac{1 - \lambda_d}{\sqrt{\lambda_d}(q - q^{-1})E} \right)^{n-1} \sqrt{\lambda_d}^{\varepsilon(\alpha)} \text{tr}_D(\delta(\alpha)),$$

where $\alpha \in \cup_{\infty} B_n$, $E = 1/d$, $\varepsilon(a)$ is as in Definition 7, δ is the natural homomorphism $\mathbb{C}(q)B_n \rightarrow Y_{d,n}(q)$ and $\lambda_d = \frac{z' - (q - q^{-1})E}{z'}$ is the re-scaling factor for the trace tr .

7.2. A different presentation for $\text{FTL}_{d,n}(u)$

Applying now Equations 7.1 and 7.2 to the defining relation (4.4) of the Framization of the Temperley-Lieb algebra, we obtain:

$$(7.4) \quad e_1 e_2 \left(1 + q(g'_1 + g'_2) + q^2(g'_1 g'_2 + g'_2 g'_1) + q^3 g'_1 g'_2 g'_1 \right) = 0.$$

This gives rise to a new presentation, with parameter q , for the Framization of the Temperley-Lieb algebra, as the quotient of $Y_{d,n}(q)$ over the ideal that is generated by the relations (7.4). We shall denote this isomorphic algebra by $\text{FTL}_{d,n}(q)$.

Given this new presentation for $\text{FTL}_{d,n}(q)$, the necessary and sufficient conditions of Theorem 6 such that the trace tr on $Y_{d,n}(q)$ passes to the quotient become:

$$(7.5) \quad x'_k = -qz' \left(\sum_{m \in \text{Sup}_1} \chi_m(t^k) + (q^2 + 1) \sum_{m \in \text{Sup}_2} \chi_m(t^k) \right),$$

$$(7.6) \quad z' = -\frac{1}{q|\text{Sup}_1| + q(q^2 + 1)|\text{Sup}_2|}.$$

Here we used the symbols x'_k and z' for the trace parameters in order to distinguish them from those of $\text{FTL}_{d,n}(u)$. If we choose the x'_k 's of Eq. 7.5 to be solutions of the E-system (by letting either Sup_1 or Sup_2 to be the empty

set), we obtain (respectively) two values for z' , the following:

$$(7.7) \quad z' = -\frac{q^{-1}E}{q^2 + 1} \quad \text{or} \quad z' = -q^{-1}E.$$

The parameters z and z' are related through the equation: $z = qz'$. Thus, by specializing in $\Theta_d(\lambda_d, q)$ the trace parameter $z' = -\frac{q^{-1}E}{q^2+1}$, we obtain the invariants for classical knots and links:

$$(7.8) \quad \theta_d(q)(\hat{\alpha}) := \left(-\frac{1 + q^2}{qE}\right)^{n-1} q^{2\varepsilon(\alpha)} \text{tr}_D(\delta(a)) = \Theta_d(q, q^4)(\hat{\alpha}),$$

where $\alpha \in \cup_{\infty} B_n$, d and E , λ_d , $\varepsilon(a)$ and δ are as in Eq. 7.3 By choosing the values mentioned above for the trace parameters z' and x'_k , $0 \leq k \leq d - 1$, we obtain $\lambda_d = q^4$.

Remark 10. The value $z' = -q^{-1}E$ of Eq. 7.7 corresponds to the value $z = -\frac{1}{|D|}$ of Proposition 8 for the algebra $\text{FTL}_{d,n}(u)$. As in Remark 7, we also discard here the value $z' = -q^{-1}E$ for topological reasons. Indeed, taking the corresponding classical Temperley-Lieb algebra $\text{TL}_n(q)$ and setting $\zeta' = -q^{-1}$, we notice that the resulting invariant fails to distinguish the unknot from the Hopf link.

7.3. Identification on knots

For the case of braids in $\cup_n B_n$, whose closure is a *knot*, the results of [3] adapt to the following:

Proposition 11. *The invariants θ_d are topologically equivalent to the Jones polynomial on knots.*

Proof. Let $\alpha \in B_n$ such that its closure $\hat{\alpha}$ is a knot. By [3, Theorem 5.17] we have that:

$$(7.9) \quad \Theta_d(q)(\hat{\alpha}) = \Theta_1(q, \lambda_d^{z'/E})(\hat{\alpha}) = P(q, \lambda_H^{z'/E})(\hat{\alpha}),$$

where $\lambda_d^{z'/E}$ (resp. $\lambda_H^{z'/E}$) stands for the re-scaling factor λ_d (resp. λ_H) with the trace parameter z' (resp. ζ') specialized to z'/E .

Notice now that: $\frac{z'}{E} = -\frac{q^{-1}}{q^2+1}$, which is the value of ζ' for which the Ocneanu trace τ passes to the algebra $\text{TL}_n(q)$. This implies that: $\lambda_d^{z'/E} =$

$q^4 = \lambda_H^{z'/E}$. Thus, Eq. 7.9 becomes:

$$\theta_d(q)(\hat{\alpha}) = \theta_1(q, \lambda_d^{z'/E})(\hat{\alpha}) = P(q, \lambda_d^{z'/E})(\hat{\alpha}) = P(q, q^4)(\hat{\alpha}) = V(q)(\hat{\alpha}).$$

□

7.4. Identification on links

For the case of *classical links*, we work as follows. In [3], using data from [1], it was observed that, out of 89 pairs of non-isotopic links, which have the same Homflypt polynomial, there are 6 pairs that are distinguished by the invariants $\Theta_d(\lambda_d, q)$. More precisely, the differences of the polynomials for each pair of links were computed and were found to be non-zero. Indeed:

$$\begin{aligned} & \Theta_d(L11n358\{0, 1\}) - \Theta_d(L11n418\{0, 0\}) \\ &= \frac{(E - 1)(\lambda_d - 1)(q - 1)^2(q + 1)^2 (q^2 - \lambda_d) (\lambda_d q^2 - 1)}{E\lambda_d^4 q^4}, \\ & \Theta_d(L11a467\{0, 1\}) - \Theta_d(L11a527\{0, 0\}) \\ &= \frac{(E - 1)(\lambda_d - 1)(q - 1)^2(q + 1)^2 (q^2 - \lambda_d) (\lambda_d q^2 - 1)}{E\lambda_d^4 q^4}, \\ & \Theta_d(L11n325\{1, 1\}) - \Theta_d(L11n424\{0, 0\}) \\ &= -\frac{(E - 1)(\lambda_d - 1)(q - 1)^2(q + 1)^2 (q^2 - \lambda_d) (\lambda_d q^2 - 1)}{E\lambda_d^3 q^4}, \\ & \Theta_d(L10n79\{1, 1\}) - \Theta_d(L10n95\{1, 0\}) \\ &= \frac{(E - 1)(\lambda_d - 1)(q - 1)^2(q + 1)^2 (\lambda_d + \lambda_d q^4 + \lambda_d q^2 - q^2)}{E\lambda_d^4 q^4}, \\ & \Theta_d(L11a404\{1, 1\}) - \Theta_d(L11a428\{0, 1\}) \\ &= \frac{(E - 1)(\lambda_d - 1)(\lambda_d + 1)(q - 1)^2(q + 1)^2 (q^4 - \lambda_d q^2 + 1)}{Eq^4}, \\ & \Theta_d(L10n76\{1, 1\}) - \Theta_d(L11n425\{1, 0\}) \\ &= \frac{(E - 1)(\lambda_d - 1)(\lambda_d + 1)(q - 1)^2(q + 1)^2}{E\lambda_d^3 q^2}. \end{aligned}$$

Note that the factor $(E - 1)$, that is common in all six pairs, confirms that the pairs have the same Homflypt polynomial, since for $E = 1$ the difference collapses to zero. Further, in [3] the values of Θ_d were computed

theoretically for one of the six pairs, using a *special skein relation* satisfied by Θ_d . Namely, as it is shown in [3], the invariants Θ_d satisfy the Homflypt skein relation, but only for crossings between different components.

For the invariants θ_d , we specialize in the above computations $z' = -\frac{q^{-1}E}{q^2+1}$ (which implies $\lambda_d = q^A$). Clearly, for $E \neq 1$ the six pairs of links above are also distinguished by the invariants θ_d . Moreover, the special skein relation of Θ_d is also valid for the invariants θ_d , specializing to the following:

$$(7.10) \quad q^{-2} \theta_d(L_+) - q^2 \theta_d(L_-) = (q - q^{-1}) \theta_d(L_0),$$

where the oriented links L_+ , L_- , L_0 comprise a Conway triple involving a crossing between different components. From the above, we have thus proved the following:

Theorem 9. *For $d \in \mathbb{Z}_{>1}$, the invariants $\theta_d(q)$ for classical links are not topologically equivalent to the Jones polynomial. Further, the invariants $\theta_d(q)$ satisfy the special skein relation of Eq. 7.10 only for crossings between different components.*

7.5. Concluding notes

The link invariants from the algebras $\text{FTL}_{d,n}(u)$ still remain under investigation. In this paper the invariants from $\text{FTL}_{d,n}(q)$ have been compared to the Jones polynomial and have been proved to be topologically non-equivalent. So the related framed link invariants might lead to new 3-manifold invariants analogous to the Witten invariants. Note that in the case of the algebras $\text{YTL}_{d,n}(u)$ the Witten invariants only can be recovered, since the related link invariants recover the Jones polynomial [9].

Acknowledgements

We would like to thank the Referee for the careful reading and for his/her valuable remarks.

References

- [1] J. C. Cha and C. Livingston, *LinkInfo: table of knot invariants*, <http://www.indiana.edu/~linkinfo>, (2015).
- [2] M. Chlouveraki and L. P. D'Andecy, *Representation theory of the Yokonuma-Hecke*, *Advances in Mathematics* **259** (2014), 134–172.

- [3] M. Chlouveraki, J. Juyumaya, K. Karvounis, and S. Lambropoulou, *Identifying the invariants for classical knots and links from the Yokonuma-Hecke algebras*, submitted for publication (2015).
- [4] M. Chlouveraki and S. Lambropoulou, *The Yokonuma-Hecke algebras and the Homflypt polynomial*, *J. Knot Theory and Its Ramifications* **22** (2013).
- [5] M. Chlouveraki and G. Pouchin, *Determination of the representations and a basis for the Yokonuma-Temperley-Lieb algebra*, *Algebras and Representation Theory* **18** (2015).
- [6] M. Chlouveraki and G. Pouchin, *Representation theory and an isomorphism theorem for the Framisation of the Temperley-Lieb algebra*, *Math. Z.* **285** (2017), no. 3, 1357–1380.
- [7] S. Chmutov, S. Jablan, K. Karvounis, and S. Lambropoulou, *On the knot invariants from the Yokonuma-Hecke algebras*, *J. Knot Theory Ramific.* **25** (2016), no. 9 (special issue dedicated to the memory of Slavik Jablan), 1641004.
- [8] P. Freyd, D. Yetter, J. Hoste, W. Lickorish, K. Millett, and A. Ocneanu, *A new polynomial invariant of knots and links*, *Bull. AMS* **12** (1985), 239–246.
- [9] D. Goundaroulis, J. Juyumaya, A. Kontogeorgis, and S. Lambropoulou, *The Yokonuma-Temperley-Lieb Algebra*, *Banach Center Pub.* **103** (2014), 73–95.
- [10] V. Jones, *Index for subfactors*, *Inventiones Mathematicae* **72** (1983), 1–25.
- [11] V. Jones, *Hecke algebra representations of braid groups and link polynomials*, *Annals of Mathematics* **126** (1987), 335–388.
- [12] J. Juyumaya, *Sur les nouveaux générateurs de l’algèbre de Hecke $\mathcal{H}(G, U, 1)$* , *J. Algebra* **204** (1998), 40–68.
- [13] J. Juyumaya, *Markov trace on the Yokonuma-Hecke algebra*, *J. Knot Theory and Its Ramifications* **13** (2004), 25–39.
- [14] J. Juyumaya, *A partition Temperley-Lieb algebra*, preprint (2013), [arXiv:1304.5158 \[math.QA\]](https://arxiv.org/abs/1304.5158).
- [15] J. Juyumaya and S. Lambropoulou, *p -adic framed braids*, *Topology and its Applications* **154** (2007), 1804–1826.

- [16] J. Juyumaya and S. Lambropoulou, *An adelic extension of the Jones polynomial*, in: M. Banagl and D. Vogel, editors, *The mathematics of knots*, Contributions in the Mathematical and Computational Sciences, Vol. 1, 825–840, Springer (2009).
- [17] J. Juyumaya and S. Lambropoulou, *An invariant for singular knots*, *J. Knot Theory and Its Ramifications* **18** (2009), no. 6, 825–840.
- [18] J. Juyumaya and S. Lambropoulou, *Modular framization of the BMW algebra* (2013). [arXiv:1007.0092v1](https://arxiv.org/abs/1007.0092v1) [[math.GT](https://arxiv.org/abs/1007.0092v1)].
- [19] J. Juyumaya and S. Lambropoulou, *p-adic framed braids II*, *Advances in Mathematics* **234** (2013), 149–191.
- [20] J. Juyumaya and S. Lambropoulou, *On the framization of knot algebras*, in: L. Kauffman and V. Manturov, editors, *New Ideas in Low-dimensional Topology*, Series on Knots and everything, World Scientific (2014).
- [21] M. Kapranov and A. Smirnov, *Cohomology determinants and reciprocity laws: Number field case* (1996), unpublished preprint.
- [22] H. Matsumoto, *Générateurs et relations des groupes de Weyl généralisés*, *C. R. Acad. Sci. Paris* **258** (1964), 3419–3422.
- [23] S. Ramanujan, *On Certain Trigonometric Sums and their Applications in the Theory of Numbers*, *Transactions of the Cambridge Philosophical Society* **22** (1918), no. 13.
- [24] H. Temperley and E. H. Lieb, *Relations between the ‘percolation’ and ‘coloring’ problem and other graph-theoretical problem associated with regular planar lattice: some exact results for the ‘percolations problems’*, *Proc. Roy. Soc. London Ser. A* **322** (1971), 251–280.
- [25] A. Terras, *Fourier Analysis of Finite Groups and Applications*, London Math. Soc. student text (1999).
- [26] E. Witten, *Quantum field theory and the Jones polynomial*, *Communications in Mathematical Physics* **121** (1989), no. 3, 351–399.
- [27] T. Yokonuma, *Sur la structure des anneaux de Hecke d’un group de Chevalley fin*, *C. R. Acad. Sc. Paris* **264** (1967), 344–347.

DEPARTMENT OF MATHEMATICS
NATIONAL TECHNICAL UNIVERSITY OF ATHENS
ZOGRAFOU CAMPUS, GR-157 80 ATHENS, GREECE
E-mail address: dgound@mail.ntua.gr

INSTITUTO DE MATEMÁTICAS, UNIVERSIDAD DE VALPARAÍSO
GRAN BRETAÑA 1091, VALPARAÍSO, CHILE
E-mail address: juyumaya@uvach.cl

DEPARTMENT OF MATHEMATICS,
NATIONAL AND KAPODISTRIAN UNIVERSITY OF ATHENS
PANEPISTIMIOUPOLIS, 15784 ATHENS, GREECE
E-mail address: kontogar@math.uoa.gr

DEPARTMENT OF MATHEMATICS
NATIONAL TECHNICAL UNIVERSITY OF ATHENS
ZOGRAFOU CAMPUS, GR-157 80 ATHENS, GREECE
E-mail address: sofia@math.ntua.gr

RECEIVED JANUARY 28, 2015

