

# Isochronal annealing studies of carbon-related defects in irradiated Si

C.A. Londos\*, M.S. Potsidi, G.D. Antonaras, A. Andrianakis

*University of Athens, Solid State Physics Section, Panepistimiopolis, Zografos, Athens 157 84, Greece*

## Abstract

We report infrared spectroscopy studies of defects in neutron-irradiated, carbon-doped, Cz-grown silicon. At room temperature irradiations, among the main defects formed are the  $C_iC_s$  and  $C_iO_i$  complexes. A peak in the spectra at  $544\text{ cm}^{-1}$  was found to be the contribution of two bands at  $543.5$  and  $545.5\text{ cm}^{-1}$ . From the corresponding annealing behavior of these bands, the  $543.5\text{ cm}^{-1}$  band was correlated with the  $C_iC_s$  defect although the  $544.5\text{ cm}^{-1}$  band with the  $C_iO_i$  defect. At high-irradiation doses, complexes as the  $C_i(Si_I)$  ( $953, 960\text{ cm}^{-1}$ ),  $C_iO_i(Si_I)$  ( $934, 1018\text{ cm}^{-1}$ ),  $C_iC_s(Si_I)$  ( $987, 993\text{ cm}^{-1}$ ) and  $C_sC_s$  ( $527\text{ cm}^{-1}$ ) form. Isochronal anneals performed in order to study the thermal evolution of these centers, showed that the  $C_i(Si_I)$  and  $C_iO_i(Si_I)$  begin to decay in the spectra around  $150^\circ\text{C}$ . Their disappearance is not accompanied by the emergence of any signal. The  $C_iC_s$  and the  $C_iC_s(Si_I)$  centers begin to decay around  $\sim 250^\circ\text{C}$ . Their disappearance is accompanied by the emergence of two pairs of bands at ( $918, 1006\text{ cm}^{-1}$ ) and ( $945, 964\text{ cm}^{-1}$ ), respectively. The origin of the centers, giving rise to these bands is discussed.

© 2005 Elsevier B.V. All rights reserved.

PACS: 61.72.Ji; 61.80.Hg; 61.72.Tt

Keywords: Silicon; Carbon impurity; Irradiation

## 1. Introduction

Carbon is, besides oxygen, the most important impurity in Si. It has been studied intensively during the last 50 years. However, in spite of the very large amount of work that has been done so far, some aspects of its behavior, especially those concerning the reaction processes that carbon participates in, are not known in detail. Upon irradiation of carbon-rich Cz-grown Si, carbon-substitutional atoms ( $C_s$ ) are ejected [1] at interstitial sites ( $C_i$ ), according to the Watkins-replacement mechanism ( $C_s + Si_I \rightarrow C_i$ ).  $C_i$  is very mobile at room temperature and it readily reacts [2] with  $C_s$  and  $O_i$  defects forming the  $C_iC_s$  and  $C_iO_i$  pairs, respectively. Most of the primary defects produced by the irradiation tend to annihilate between themselves ( $V + Si_I \rightarrow \emptyset$ ), although some of them are captured by impurities present in the material ( $V + O_i \rightarrow VO, Si_I + C_s \rightarrow C_i$ , etc.) and some other pair with each other ( $V + V \rightarrow V_2, Si_I + Si_I \rightarrow (Si_I - Si_I)$ ). De-

fect reaction modeling foresees and experiments verify [3–5] that a percentage of the  $Si_I$ 's is also captured by centers formed during the irradiation. Thus, the  $C_i$ , the  $C_iC_s$  and  $C_iO_i$  defects act as nucleation sites for the  $Si_I$ 's and at high radiation doses, complexes as the  $C_i(Si_I)$ , the  $C_iO_i(Si_I)$  and the  $C_iC_s(Si_I)$  form, as well. A pair of local vibrational mode (LVM) bands at ( $953, 966\text{ cm}^{-1}$ ) has been correlated [2] with the  $C_i(Si_I)$ , another pair at ( $940, 1024\text{ cm}^{-1}$ ) has been correlated [2] with the  $C_iO_i(Si_I)$  complex and another one at ( $987, 993\text{ cm}^{-1}$ ) has been correlated [6] with the  $C_iC_s(Si_I)$  complex. On the other hand, vacancies are also trapped by the  $C_iC_s$  complexes [7] leading to the formation of the  $C_sC_s$  defect ( $527.4, 748.7\text{ cm}^{-1}$ ), through the reaction  $C_iC_s + V \rightarrow C_sC_s$ .

In neutron-irradiated Si, due to the spatial separation [8] of the produced vacancies and self-interstitials, complexes related to these defects form in larger concentrations. Therefore, weak signals in the spectra related to such complexes are expected to be detected more easily, facilitating their study. The purpose of this work is to study the production and evolution with temperature of the  $C_i(Si_I)$ ,  $C_iC_s(Si_I)$ ,  $C_iO_i(Si_I)$  and  $C_iC_s$  defects. It is a

\*Corresponding author. Tel./fax: +30 210 7276726.

E-mail address: [hlontos@phys.uoa.gr](mailto:hlontos@phys.uoa.gr) (C.A. Londos).

continuation of a previous work [6], aiming in particular at throwing new light on carbon-related defects and processes, where aggregations of primary defects are involved.

## 2. Experimental details

Prepolished Si samples of 2 mm thickness with initial oxygen and carbon concentrations  $[O_i]_0 = 7.2 \times 10^{17} \text{ cm}^{-3}$  and  $[C_s]_0 = 1.5 \times 10^{17} \text{ cm}^{-3}$ , respectively, were irradiated with 5 MeV fast neutrons at a fluence of  $\sim 1 \times 10^{17} \text{ n/cm}^2$ . The samples were wrapped in Cd foils to eliminate thermal neutrons and put inside sealed quartz boxes to avoid water contamination. After the irradiations, 20 min isochronal anneals of  $\sim 10^\circ\text{C}$  steps, in open furnaces were performed. After each annealing stage, infrared spectroscopy measurements were carried out, at room temperature, with a Jasco-IR 700 double beam dispersive spectrometer.

## 3. Experimental results and discussion

Fig. 1 shows a section of the absorption spectrum of our Si samples immediately after the irradiation. By using Lorentzian profiles (Fig. 2a) it is found that the band at  $544 \text{ cm}^{-1}$  is the overlap of two contributing bands at  $543.5$  and  $545.5 \text{ cm}^{-1}$ . Figs. 2(b) and (c) exhibit the thermal evolution of the later bands. Previously IR studies [9], in electron-irradiated Si, have detected the formation of the  $C_1C_s$  center and six bands at  $540.4$ ,  $543.3$ ,  $579.8$ ,  $640.6$ ,  $730.4$  and  $842.4 \text{ cm}^{-1}$  have been attributed to the B configuration of the defect, corresponding to its neutral charge state, in which the center is expected to be in irradiated material. These bands are seen only at liquid He temperatures. They are very weak, and this fact along with strong unharmonic effects [9,10] prevents their detection at room temperature. The bands are stable up to  $\sim 250^\circ\text{C}$ . It has also been reported [11,12] that a band at  $\sim 540 \text{ cm}^{-1}$ , detected in electron-irradiated material at liquid He

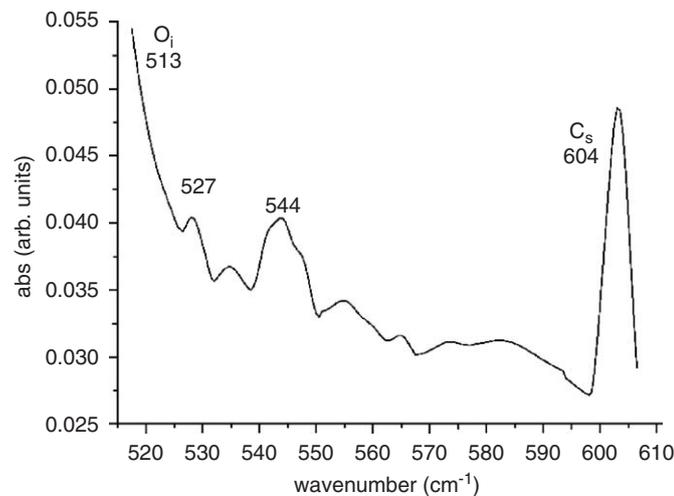


Fig. 1. Section of absorption spectrum of neutron-irradiated, carbon-doped, Cz-grown Si.

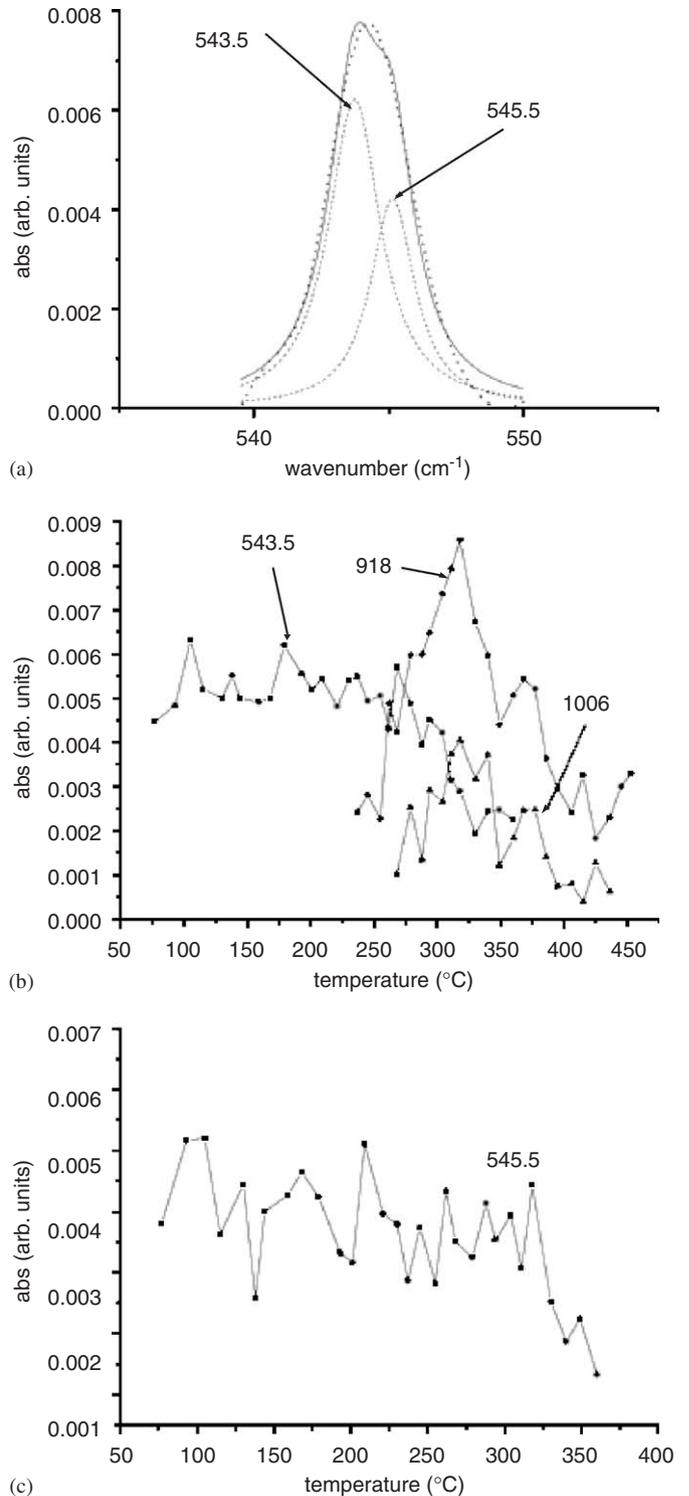


Fig. 2. (a) Lorentzian profiles of the  $544 \text{ cm}^{-1}$  IR band, (b) the thermal evolution of the  $543.5$  and the  $(918, 1006 \text{ cm}^{-1})$  pair of bands, (c) the thermal evolution of the  $545.5 \text{ cm}^{-1}$  band.

temperatures, being stable up to  $\sim 300^\circ\text{C}$  is related to the  $C_1O_i$  center. Thus, in our case, where both the  $C_1C_s$  and  $C_1O_i$  defects are expected to form and be stronger due to the neutron irradiation, we can reasonably assume that they contribute to the same broad peak making their signal

detectable at room temperature. In this sense, based on their annealing behavior (Figs. 2(b) and (c)) we correlate the bands at  $543.5$  and  $545.5\text{ cm}^{-1}$  with the  $\text{C}_i\text{C}_s$  and  $\text{C}_i\text{O}_i$  defects, correspondingly. We note at this point that the decay of the  $543.5\text{ cm}^{-1}$  band is accompanied in the spectra (Fig. 2(b)) by the emergence of two bands at  $918$  and  $1006\text{ cm}^{-1}$  (Fig. 3).

As we mentioned in the Introduction, the  $\text{Si}_i$ 's tend to aggregate on carbon-related complexes. Fig. 4(a) shows a section of the absorption spectrum received immediately after the irradiation. The pairs of bands at  $(953, 960\text{ cm}^{-1})$ ,  $(934, 1018\text{ cm}^{-1})$  and  $(987, 993\text{ cm}^{-1})$  have been previously attributed [2,6] to the  $\text{C}_i(\text{Si}_i)$ ,  $\text{C}_i\text{O}_i(\text{Si}_i)$  and  $\text{C}_i\text{C}_s(\text{Si}_i)$  complexes, respectively. Figs. 4(b) and (c) show the evolution with temperature of the  $\text{C}_i(\text{Si}_i)$  and  $\text{C}_i\text{O}_i(\text{Si}_i)$  complexes, respectively.

The decay of the corresponding peaks is not accompanied by the emergence of other peaks in the spectra. Further on, Fig. 5(a) shows the thermal evolution of the  $(987, 993\text{ cm}^{-1})$  pair of bands the decay of which is accompanied by the emergence in the spectra of two bands at  $945$  and  $964\text{ cm}^{-1}$  (Fig. 5(b)). Photoluminescence studies [13,14] have reported that the destruction of the  $969\text{ meV}$  PL line of the  $\text{C}_i\text{C}_s$  defect (G-line) is partly connected with the growth of lines at  $951, 953, 954$  and  $957\text{ meV}$  which have been attributed to G-centers perturbed by nearby self-interstitials. Additional PL lines have been reported [16] to arise upon the destruction of the G-line of the  $\text{C}_i\text{C}_s$  defect. Moreover, a line at  $949.9\text{ meV}$ , the so-called F-line, has been suggested [17] to arise from a  $\text{C}_3\text{O}$  defect formed when a mobile  $\text{C}_i\text{C}_s$  defect is captured by a  $\text{C}_i\text{O}_i$  defect. We suggest that our pairs of LVM bands at  $(918, 1006\text{ cm}^{-1})$  (Fig. 2(b)) and  $(945, 964\text{ cm}^{-1})$  (Fig. 5(a)), may arise from the same centers that give rise to the above PL lines. Our suggestion is based mainly on the similarities in the emergence of the bands from the IR and PL studies and at this stage, any positive identification by assigning certain IR bands to certain structure cannot be made.

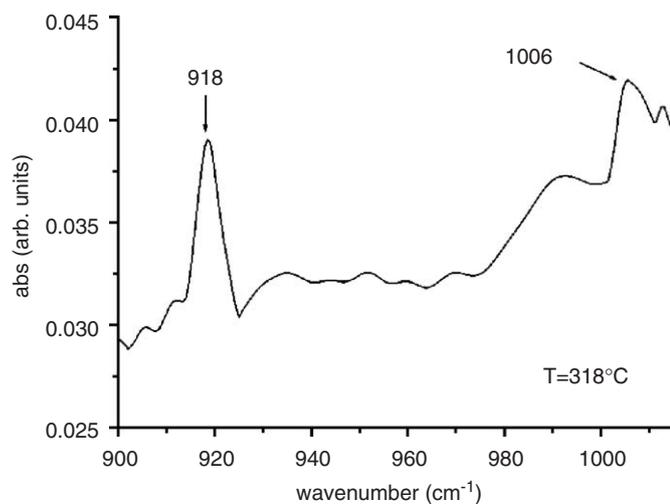


Fig. 3. Section of the absorption spectrum showing the  $918$  and  $1006\text{ cm}^{-1}$  bands.

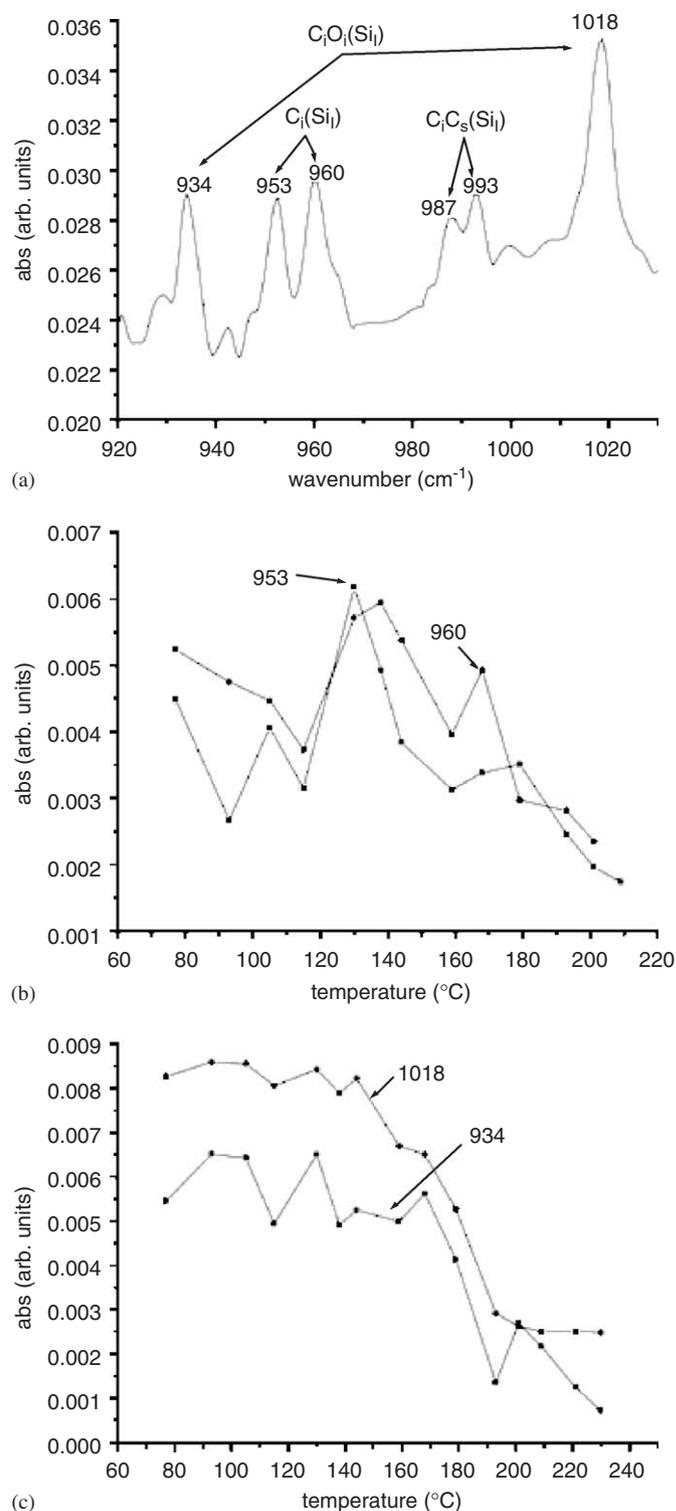


Fig. 4. (a) Section of absorption spectrum of neutron-irradiated, carbon-doped, Cz-grown Si taken immediately after room temperature irradiation, (b) the thermal evolution of the  $(953, 960\text{ cm}^{-1})$  bands of the  $\text{C}_i(\text{Si}_i)$  defect, (c) the thermal evolution of the  $(934, 1018\text{ cm}^{-1})$  bands of the  $\text{C}_i\text{O}_i(\text{Si}_i)$  defect.

Vacancies also aggregate on carbon-related defects produced by the irradiation, leading to the formation of larger complexes. Such a complex is the  $\text{C}_s\text{C}_s$  structure

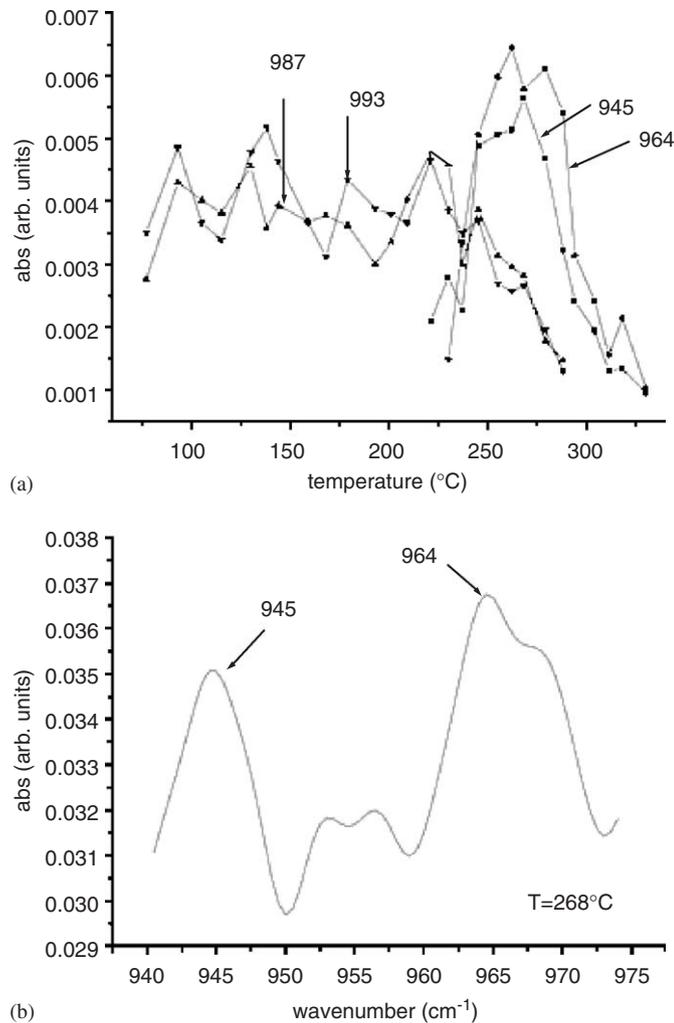


Fig. 5. (a) The thermal evolution of the (945 and 964  $\text{cm}^{-1}$ ) pair of bands grown in the spectra when the (987 and 993  $\text{cm}^{-1}$ ) pair of bands of the  $\text{C}_i\text{C}_s(\text{Si}_i)$  defect anneals out, (b) section of the absorption spectrum showing the (945 and 964  $\text{cm}^{-1}$ ) pair of bands.

[7,15] which forms through the reaction  $\text{C}_i\text{C}_s + V \rightarrow \text{C}_s\text{C}_s$ . In its neutral charge state the defect gives rise to an IR band at  $527.4\text{cm}^{-1}$ . In our neutron-irradiated samples, a band at  $527\text{cm}^{-1}$ , detected at room temperature (Fig. 1), is tentatively attributed to the  $\text{C}_s\text{C}_s$  defect. Its evolution with temperature is given in Fig. 6.

#### 4. Summary

We have investigated carbon-related complexes in neutron-irradiated Si, by means of IR spectroscopy. A broad band at  $544\text{cm}^{-1}$  was found to be the convolution of two bands at  $543.5$  and  $545.5\text{cm}^{-1}$ . Based on their thermal stability, we attributed these bands to the  $\text{C}_i\text{C}_s$  and the  $\text{C}_i\text{O}_i$  defects, correspondingly. The evolutions with temperature of the  $\text{C}_i(\text{Si}_i)$  ( $953, 960\text{cm}^{-1}$ ),  $\text{C}_i\text{O}_i(\text{Si}_i)$  ( $934, 1018\text{cm}^{-1}$ ) and  $\text{C}_i\text{C}_s(\text{Si}_i)$  ( $987, 993\text{cm}^{-1}$ ) complexes were stu-

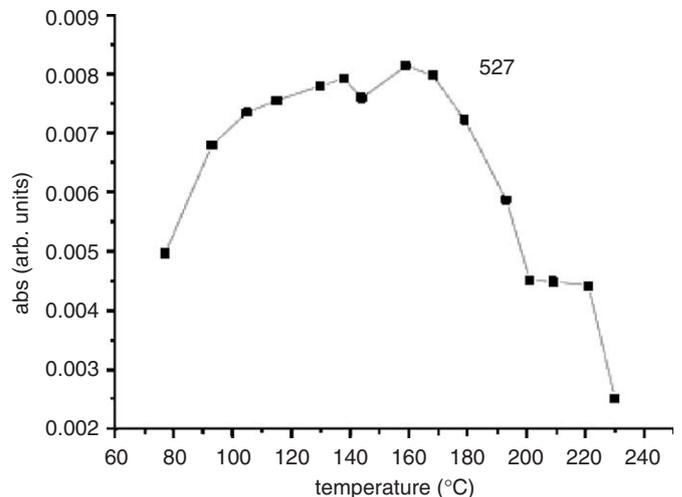


Fig. 6. The thermal evolution of the  $527\text{cm}^{-1}$  band.

died. Further on, a band at  $527\text{cm}^{-1}$  was tentatively attributed to the  $\text{C}_s\text{C}_s$  center.

#### Acknowledgments

This work was supported by the special account for research grants of the National and Kapodistrian University of Athens (Grant no. 70/4/3307).

#### References

- [1] R.C. Newman, A.R. Bean, *Radiat. Eff.* 8 (1970) 189.
- [2] G. Davies, R.C. Newman, in: S. Mahajan (Ed.), *Handbook in Semiconductors*, vol. 3, Elsevier, Amsterdam, 1994, pp. 1557–1635.
- [3] G. Davies, E.C. Lighthowers, R.C. Newman, A.S. Oates, *Semicond. Sci. Technol.* 2 (1987) 524.
- [4] S.P. Chappell, R.C. Newman, *Semicond. Sci. Technol.* 2 (1987) 691.
- [5] B.C. MacEvoy, S.J. Watts, *Solid State Phenomena* 57–58 (1997) 221.
- [6] C.A. Londos, M.S. Potsidi, E. Stakakis, *Physica B* 340–342 (2003) 551.
- [7] E.V. Lavrov, B.B. Nielsen, J.R. Byberg, B. Hourasine, R. Jones, S. Öberg, P.R. Briddon, *Phys. Rev. B* 62 (2000) 158.
- [8] Fukata, T. Ohori, M. Suezawa, H. Takahashi, *J. Appl. Phys.* 91 (2002) 5831.
- [9] E.V. Lavrov, L. Hoffman, B. Bech Nielsen, *Phys. Rev. B* 60 (1999) 8081.
- [10] P. Leary, R. Jones, S. Öberg, V.J.B. Torres, *Phys. Rev. B* 55 (1997) 2188.
- [11] L.I. Murin, V.P. Markevich, J.L. Lindström, M. Kleverman, J. Hermansson, T. Hallberg, B.G. Svensson, *Solid State Phenomena* 82–84 (2002) 57.
- [12] J. Coutinho, R. Jones, P.R. Briddon, S. Öberg, L.I. Murin, V.P. Markevich, J.L. Lindström, *Phys. Rev. B* 65 (2001) 014109.
- [13] G. Davies, *Mater. Sci. Forum* 38–41 (1989) 151.
- [14] G. Davies, E.C. Lighthowers, M.C. doCarmo, J.G. Wilkes, G.R. Wolstenholme, *Solid State Commun.* 59 (1984) 1059.
- [15] J.R. Byberg, B. Bech Nielsen, M. Fanciulli, S.K. Estreicher, P.A. Fedders, *Phys. Rev. B* 61 (2000) 1239.
- [16] T.K. Kwok, *Phys. Rev. B* 51 (1995) 17188.
- [17] E.C. Lightowers, A.N. Safonov, *Mater. Sci. Forum* 258–263 (1997) 617.