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IR studies of oxygen-related and carbon-related defects in Sn-doped silicon

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The ability of tin (Sn) to capture vacancies in silicon material attracts considerable interest in relation with electronic applications where radiation hardness is necessary. In this sense the full knowledge of the behavior of radiation-induced defects in Sn-doped Si is particularly important. The present report gives a systematic study of the evolution of the vacancy-oxygen (VO_n, $1 \le n \le 6$) defects and the carbon-oxygen (C_iO_i) defect in electron-irradiated Sn-doped Cz-Si containing carbon. The radiation-induced defects and their reactions were investigated by Infrared Spectroscopy (IR) in the course of 20 min isochronal an-

neals from room temperature up to ~750 °C. Interestingly, in Sn-doped Cz-Si samples we were able to monitor and study IR bands correlated with VO₅/VO₆ defects. In particular, we tentatively attributed two bands at 1037, 1051 cm⁻¹ to the VO₅ defect, and two other bands at 967, 1005 cm⁻¹ to the VO₆ defect. Concerning carbon-related defects, our results show that the disappearance of the 862 cm⁻¹ band of the C_iO_i defect is accompanied in the spectra by the emergence of a band at 1048 cm⁻¹ previously correlated with a C_sO_{2i} structure.

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1 Introduction Carbon and oxygen are the main unintentionally added impurities in Si during growth. Carbon is an isovalent impurity in Si at substitutional sites, being electrically inactive. Due to its smaller covalent radius than that of Si ($r_c = 0.77$ Å, $r_{Si} = 1.17$ Å) carbon presence induces tensile strains in the Si lattice. On the other hand, oxygen is located interstitially in the Si lattice being also electrically inactive. Both impurities do not introduce levels in the gap. However, during the various processing stages of the Si material, for instance during irradiation, they participate in the production of electrically active complexes which affect the properties and behavior of Si.

The immediate products of irradiation are vacancies and self-interstitials. Oxygen is the main trap for vacancies leading to the formation of VO defect (V+O_i \rightarrow VO) the socalled A-center [1, 2] which introduces a level at E_c-0.17 eV in the gap [3]. Carbon is the main trap for selfinterstitials leading to the formation of the C_iO_i and the C_iC_s defects [4] in Si, through the Watkins exchange mechanism ($C_s + Si_I \rightarrow C_i$). In Cz-Si the C_iO_i center is the dominant carbon-related defect giving rise to a level at $E_v+0.36$ eV in the gap [4, 5]. Both VO and C_iO_i defects are optically active. An IR band at 830 cm⁻¹ has been assigned to the VO defect [2]. At least six bands, with the strongest one at 860 cm⁻¹, have been assigned to the C_iO_i defect [4].

The main reaction channel that VO defect participates upon annealing is that of the addition of oxygen atoms, leading to the formation of VO_n complexes [2, 6]. VO defect is stable to ~ 300 °C and then becomes mobile to produce VO₂ defect (VO+O_i \rightarrow VO₂). An IR band at 885 cm⁻¹ has been assigned to the VO₂ defect [6, 7]. VO₂ defect is stable to ~ 400°C and then it is captured by a mobile oxygen atom leading to the formation of VO₃ defect (VO₂+O_i \rightarrow VO₃). Three IR bands at 904, 968 and 1000 cm⁻¹ have been assigned to the VO₃ defect [6, 7]. VO₃ de-



fect is stable to $\sim 500~^{\rm o}C$ and then by capturing an oxygen atom forms VO₄ defect (VO₃+O₁ \rightarrow VO₄) which is stable to ~ 550 °C. Two bands at 985 and 1009 cm⁻¹ have been assigned to the VO_4 defect [8]. At higher temperatures, where oxygen atoms are very mobile in the Si lattice, VO_5 and VO₆ defects are formed either from the addition of an oxygen atom via the reactions $VO_4+O_1 \rightarrow VO_5, VO_5+O_1 \rightarrow$ VO₆ correspondingly, or with the participation of oxygen dimers by following the reactions $VO_3+O_{2i} \rightarrow VO_5$, $VO_4+O_{2i} \rightarrow VO_6$ [9, 10]. Other reactions are also possible as for example $V_2O_3 \rightarrow O_i \rightarrow V_2O_4$, $V_2O_3 + 2O_i \rightarrow V_2O_5$, $V_3O_2 + O_i \rightarrow V_3O_3$, etc leading generally to the formation of V_nO_m defects [8]. A number of bands have been generally correlated with the VO₅/VO₆ defects [9, 10] or/and with V_nO_m structures [8]. Any particular assignment of bands to the VO_5 or the VO_6 defect has not been made so far. Furthermore the thermal stability of the VO₅ and VO₆ defects has not been investigated.

The formation of the C_iO_i defects occurs when a mobile C_i atom is trapped by an O_i atom ($C_i + O_i \rightarrow C_iO_i$). It has been reported that the C_iO_i defect, which is stable up to 300 °C, is then destroyed either by dissociation or by reaction with another defect [4, 11].

Doping of Si with isovalent impurities improves the performance of devices for certain applications. Due to its larger covalent radius than that of Si ($r_{Sn} = 1.41$ Å, $r_{Si} = 1.17$ Å) Sn acts as a trap for vacancies and small impurity atoms as for example carbon. Noticeably, the formation of SnV [12] and SnC_i [13] defects has been reported in the literature. On the other hand, C reacts promptly with self-interstitials. Thus, after the irradiation the simultaneous presence of carbon and tin in the Si lattice affects the number of vacancies and self-interstitials available to participate in the various reactions and in particular those involving C and O impurities. Therefore, the formation of oxygen-related and carbon-related defects is expected to be affected.

Given in addition the variety of reactions that carbon interstitials participate and their sensitivity in the annealing behaviour under various conditions [4, 14-16] it would be interesting to investigate the interaction between radiation-induced defects and their behaviour upon thermal anneals in Si material containing carbon and tin. The main objective of this work was to study the effect of Sn doping in the behaviour of oxygen-related and carbon-related defects in Si and in particular those of VO_n and C_iO_i defects.

2 Experimental details The Si material used in this work was Czochralski crystals doped with tin (Sn). The Sn concentration was $[Sn]_o=3x10^{17}$ cm⁻³. The oxygen and carbon concentrations were $[O_i]_o=9.6 \times 10^{17}$ cm⁻³ and $[C_s]_o=4.7\times10^{17}$ cm⁻³, respectively. The samples were irradiated with 2 MeV electrons at 95 °C using the Dynamitron accelerator at Takasaki-JAERI (Japan). The irradiation dose was 1×10^{18} cm⁻². The evolution of the radiation induced defects were studied by 20 min isochronal anneals in steps of 10 °C up to 750 °C. After each annealing step, the

IR spectra were recorded at room temperature by means of a FTIR spectrometer with a resolution of 1cm⁻¹. The two phonon intrinsic absorption was always subtracted by using a float-zone sample of equal thickness.

3 Experimental results and discussion Figure 1 presents sections of the IR spectra of the Sn-doped samples immediately after irradiation and at various selective temperatures in a course of the isochronal anneals. A number of defects were observed. Actually, the well-known defects VO (830 cm⁻¹), VO₂ (888 cm⁻¹), VO₃ (904, 968, 1000 cm⁻¹), VO₄ (985, 1009 cm⁻¹) are present in the spectra together with the less known bands at 967, 1005, 1037 and 1051 cm⁻¹ which appear above 550 °C. Figures 2 and 3 show the



Figure 1 Fragments of the IR spectra at characteristic temperatures for the Sn-doped Si samples.

thermal evolution of all the bands in the samples. Noticeably, at room temperature the two bands at 1037 and 1051 cm⁻¹ have been seen previously and attributed to either V_nO_m structures [8] and/or to VO₅/ VO₆ [9, 6] defects. The two other bands at 967 and 1005 cm⁻¹ have also been detected previously [8] in neutron irradiated Si and correlated with $V_n O_m$ defects. On the other hand, measurements at cryogenic temperatures of electron-irradiated Si have detected additional bands at 992, 1024, and 1108 cm⁻¹ which were correlated [9] with the VO_5/VO_6 defects. In the present detailed work we have observed that the bands at 1037 and 1051 cm⁻¹ appear together in the spectra upon annealing of the VO₄ defects and disappear also together just above 700 °C. Another pair of bands at 967 and 1005 cm⁻¹ also appear together in the spectra upon disappearing of the VO₄ defect and decay out also together about 740 °C. Judging from their thermal stability we tentatively correlated the first pair of bands at 1037, 1051 cm⁻¹ with the VO₅ defect and the other pair at 967 and 1005 cm⁻¹ to the VO₆ defect. In this attribution we have taken into account that all the bands emerge in the spectra upon annealing of the VO_4 defect and therefore it is more reasonable to correlate them with VO₄ defects, instead of the more general attribution to V_nO_m structures. From the general behavior of the thermal stability of the VO_n ($1 \le n \le 4$), where the larger the number



Figure 2 The thermal evolution of VO, VO₂, and VO₃ defects in the Sn-doped Si samples.



Figure 3 The thermal evolution of VO_3 , VO_4 , and VO_5/VO_6 defects in the Sn-doped Si samples.

n of the oxygen atoms the higher the annealing temperature of the structure, we attribute the first pair of bands (1037, 1051 cm⁻¹) to the VO₅ defect annealing out at \sim 710 °C and the second pair of bands at (967, 1005 cm⁻¹) to the VO₆ defect which anneals out at a higher temperature (~750 °C). It is worth noting that the above four weak bands of VO₅ and VO₆ defects appear in the IR spectra of neutron irradiated Cz-Si received at room temperature [8], as well as in the spectra of electron irradiated Si received at He liquid temperatures [9]. Understandingly, in neutron irradiation the concentration of defects is comparably larger than that of electron irradiation, allowing their detection at room temperatures measurements in the former case which may not be detected in the latter case. Also measurements at cryogenic temperatures allow for weak signals to be seen in the IR spectra, although may pass undetected in room temperature measurements. We note that the two bands at 947 and 966 cm⁻¹, arising in the spectra in the course of isochronal anneals at T = 250 °C, are most probably related to the $C_iC_s(Si_1)_2$ defect [17]. Figure 4 shows the evolution of the C_iO_i defect in the Sn-doped Si samples. Interestingly, the disappearance of the C_iO_i band is accompanied in the spectra by the emergence of another band at 1048 cm⁻¹. This band is very weak and has not been reported in early works, see for example Ref. [4]. Recently, this band has been detected [18] and correlated with a C_sO_{2i} structure. Possible formation reaction may be C_iO_i + VO \rightarrow $C_sO_{2i}.$ Both the VO and C_iO_i defects anneal out at around the same temperature ~300 °C. It deserves noting at this point that two bands at 1052 and 1099 cm⁻¹ have been reported in Si material containing carbon either non-irradiated [19] or irradiated at elevated temperatures in the range 400-600 °C [20]. In these studies [19, 20] the measurements were carried out at cryogenic temperatures. The bands were correlated with the C_sO_{2i} defect. The suggested formation reaction was $VO_2 + C_i \rightarrow C_sO_{2i}$ [20]. In our studies, the 1099 cm⁻¹ was not detected. However, its weak signal could have been masked in our RT measurements by the very strong signal of O_i in this region.



Figure 4 The thermal evolution of C_iO_i and C_sO_{2i} in the Sn-doped Si samples.

Based on the above results a possible explanation of the role of the tin impurity on the evolution of the oxygenrelated radiation-induced defects VO_n ($1 \le n \le 6$) and in particular on the formation of the VO_5/VO_6 defects could be envisaged. It has been reported that the diffusion constant of oxygen is considerably enhanced in Si doped with tin [21]. This is possibly due to the elastic strains introduced in the lattice, as a result of the larger covalent radius of Sn than that of the Si atoms. This may leads to a decrease of the potential barrier for the diffusion of the interstitial oxygen atom. Consequently, the trapping rates of oxygen in the reactions $VO_4 + O_i \rightarrow VO_5$, $VO_5 + O_i \rightarrow VO_6$, become larger leading to an enhanced formation of the VO₅ and VO_6 defects. Similarly, the reactions $VO_3+O_{2i} \rightarrow VO_5$ and $VO_4+O_{2i} \rightarrow VO_6$, which involve dimers are affected. Dimers are much faster diffusion species than single oxygen atoms. The formation of dimmers and also their diffusivity is also expected to be enhanced. Thus, larger concentrations of VO₅ and VO₆ defects are produced, making them detectable in the IR spectra. On the other hand, the role of tin on the evolution of carbon-related defects and in par-



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ticular on the formation of the C_sO_{2i} defect cannot be definite extracted from these measurements. The carbon content of the samples was quite high and the detection of the 1048 cm⁻¹ band of the defect may be related simply with this. Of course, Sn is an efficient trap of vacancies, allowing for more self-interstitials to be trapped by substitutional carbon atoms in the course of irradiation. This is expected to lead to the formation of C-related defects in larger concentrations, something which is reflected in the spectra with the enhanced formation of the C_sO_{2i} defect.

4 Conclusions We have studied the evolution of VO_n and C_iO_i defects in Si material containing carbon and tin by FTIR. We found that signals from VO_n structures with n>4 are present in the IR spectra and two pairs of bands at (1037, 1058 cm⁻¹) and (967, 1005 cm⁻¹) have been correlated with the VO_5 and VO_6 defects, respectively. Furthermore the decay out of the 860 cm⁻¹ band of C_iO_i defects was found to be accompanied in the spectra by the emergence of another band at 1048 cm⁻¹ attributed to a C_sO_{2i} structure. An explanation of the role of tin in the evolution of defects was offered.

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